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by

Larmor Methods for Neutron Scattering

Neutron Scattering

• What fraction of neutrons were scattered through an angle 2θ and with a change in velocity of ΔV ?



- Neutrons are **particles**
- Relate neutron scattering intensity distribution to sample properties -- Neutrons are **waves** $\lambda = 2d \sin\theta$
- But I slipped something past you.....

Larmor labeling

 Neutrons are quantum particles with a spin that has two possible states – "up" and "down". It's state can be written as a linear combination of "up" and "down" which evolve separately in a magnetic field ^{"In just 38 months, you can earn big PROFIT\$"} as a fully trained QUANTUM MECHANIC!"

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- Neutrons carry a pointed stick called the "classical spin" that rotates when the neutron is in a magnetic field
- The rotation is called Larmor precession and occurs at a rate of 1.83 x 10⁸ x B(T) per second
- This precession can be used as a clock to measure the neutron speed Plug & chug





Larmor Methods use the Neutron Spin to improve neutron instrument performance

- An example is Neutron Spin Echo (NSE)
- Traditional methods
 - define *both* incident & scattered wavevectors in order to define E and Q accurately
 - use collimators, monochromators, choppers etc to define both \mathbf{k}_i and \mathbf{k}_f

• NSE

- measure as a function of the *difference* between appropriate components of \mathbf{k}_i and \mathbf{k}_f (original use: measure $\mathbf{k}_i \mathbf{k}_f$ i.e. energy change)
- use the neutron's spin polarization to encode the difference between components of \mathbf{k}_i and \mathbf{k}_f
- can use large beam divergence &/or poor monochromatization to increase signal intensity, while maintaining very good resolution

The Principles of NSE are Very Simple

- If a spin rotates anticlockwise & then clockwise by the same amount it comes back to the same orientation
 - Need to reverse the direction of the applied field to reverse rotation
 - The effect is independent of neutron speed provided the speed is constant
- The same effect can be obtained by reversing the precession angle at the mid-point and continuing the precession in the same sense
 - Use a π flip
- If the neutron's velocity, v, is changed by the sample, its spin will not come back to the same orientation
 - The difference will be a measure of the *change* in the neutron's speed or energy

For Quasi-elastic Scattering, the Echo Polarization depends on Energy Transfer

• If the neutron changes energy when it scatters, the precession phases before & after scattering, $\phi_1 \& \phi_2$, will be different:

using
$$\hbar\omega = \frac{1}{2}m(v_1^2 - v_2^2) \approx mv\delta v$$

 $\phi_1 - \phi_2 = \gamma Bd\left(\frac{1}{v_1} - \frac{1}{v_2}\right) \approx \frac{\gamma Bd}{v^2}\delta v \approx \frac{\gamma Bd\hbar\omega}{mv^3} = \frac{\gamma Bdm^2\lambda^3\omega}{2\pi\hbar^2}$

- To lowest order, the difference between φ₁ & φ₂ depends only on ω (I.e. δv) & <u>not</u> on v₁ & v₂ separately
- The measured polarization, <P>, is the average of $cos(\phi_1 \phi_2)$ over all transmitted neutrons I.e.

$$\left\langle P \right\rangle = \frac{\iint I(\lambda)S(\vec{Q},\omega)\cos(\phi_1 - \phi_2)d\lambda d\omega}{\iint I(\lambda)S(\vec{Q},\omega)d\lambda d\omega}$$

Neutron Polarization at the Echo Point is a Measure of the Intermediate Scattering Function

$\left\langle P \right\rangle = \frac{\iint I(\lambda)S(\vec{Q},\omega)\cos(\phi_1 - \phi_2)d\lambda d\omega}{\iint I(\lambda)S(\vec{Q},\omega)d\lambda d\omega} \approx \left\langle \int S(\vec{Q},\omega)\cos(\omega\tau)d\omega \right\rangle = I(\vec{Q},\tau)$						
$\int J^{-}(U)^{-}(\mathcal{L}, U)^{-}(U)^{-}(\mathcal{L}, U)^{-}(U$	Bd	λ(nm)	τ(ns)			
where the "spin echo time" $\tau = \gamma B d \frac{m^2}{m^2} \lambda^3$	(T.m)					
$2\pi h^2$	1	0.4	12			
	1	0.6	40			
	1	1.0	186			

- I(Q,t) is called the intermediate scattering function
 - Time Fourier transform of $S(\vec{Q},\omega)$ or the \vec{Q} Fourier transform of $G(\vec{r},t)$, the two particle correlation function
- NSE probes the sample dynamics as a function of time rather than as a function of $\boldsymbol{\omega}$
- The spin echo time, τ , is the "correlation time"

What does a NSE Spectrometer Look Like? IN11 at ILL was the First



The Principle of Neutron Resonant Spin Echo

- Within a coil, the neutron is subjected to a steady, strong field, B_0 , and a weak rf field $B_1 cos(\omega t)$ with a frequency $\omega = \omega_0 = \gamma B_0$
 - Typically, $B_0 \sim 100 \text{ G}$ and $B_1 \sim 1 \text{ G}$



- In a frame rotating with frequency $\omega_{\rm 0},$ the neutron spin sees a constant field of magnitude $\rm B_1$
- The length of the coil region is chosen so that the neutron spin precesses around B_1 thru an angle π .
- The neutron precession phase is:

 $\varphi_{neutron}^{exit} = \varphi_{RF}^{exit} + (\varphi_{RF}^{entry} - \varphi_{neutron}^{entry})$

$$= 2\varphi_{RF}^{entry} - \varphi_{neutron}^{entry} + \omega_0 d / v$$





Table 1. Spin orientation

	Time t	Phase field B_r	neutron Spin phase ${\cal S}$
Α	t_A	ωt_A	0
A'	$t_{A'} = t_A + \frac{d}{v}$	$\omega t_{A'}$	$2\omega t_A + \omega rac{d}{v}$
В	$t_B = t_A + \frac{l_{AB} + d}{v}$	ωt_B	$2\omega t_A + \omega rac{d}{v}$
в'	$t_{B'} = t_A + \frac{l_{AB} + 2d}{v}$	$\omega t_{B'}$	$2\omega \frac{l_{AB}+d}{v}$
С	t_C	$-\omega t_C$	$2\omega \frac{l_{AB}+d}{v}$
C'	$t_{C'} = t_C + rac{d}{v}$	$-\omega t_{C'}$	$-\omega \frac{d}{v'} - 2\omega t_C - 2\omega \frac{l_{AB}+d}{v}$
D	$t_D = t_C + rac{l_{CD}+d}{v'}$	$-\omega t_D$	$-\omega \frac{d}{v'} - 2\omega t_C - 2\omega \frac{l_{AB}+d}{v}$
D'	$t_{D'} = t_C + \frac{l_{CD} + 2d}{v'}$	$-\omega t_{D'}$	$2\omega(\frac{l_{AB}+d}{v}-\frac{l_{CD}+d}{v'})$

Echo occurs for elastic scattering when $l_{AB} + d = l_{CD} + d$

* Courtesy of S. Longeville



The Measured Polarization for NRSE is given by an Expression Similar to that for Classical NSE

• Assume that $v' = v + \delta v$ with δv small and expand to lowest order, giving:

$$\left\langle P \right\rangle = \frac{\iint I(\lambda)S(\vec{Q},\omega)\cos(\omega\tau_{NRSE})d\lambda d\omega}{\iint I(\lambda)S(\vec{Q},\omega)d\lambda d\omega}$$

where the "spin echo time"
$$\tau_{NRSE} = 2\gamma B_0 (l+d) \frac{m^2}{2\pi h^2} \lambda^3$$

- Note the additional factor of 2 in the echo time compared with classical NSE (a factor of 4 is obtained with "bootstrap" rf coils)
- The echo is obtained by varying the distance, *l*, between rf coils
- In NRSE, we measure neutron velocity using fixed "clocks" (the rf coils) whereas in NSE each neutron "carries its own clock" whose (Larmor) rate is set by the local magnetic field

The Classical Picture of Spin Precession



Spin Echo Scattering Angle Measurement (SESAME) No Sample in Beam





Spin Echo Scattering Angle Measurement (SESAME) Scattering of a Divergent Beam



Spin-echo angular encoding: the experiment



The neutron is a spin-1/2 particle with a magnetic moment



- In an applied magnetic field, B_z, the neutron has two spin eigenstates, "up" and "down" denoted I0> & I1>
 - Magnetic moment either along z or –z (magnetic moment is antiparallel to the spin)
- Quantum mechanics tells us that any neutron spin state is a linear combination of the two eigenstates

$$\chi = a |0\rangle + b |1\rangle \quad \text{where } |a|^2 + |b|^2 = 1\chi$$
$$\chi = \cos(\theta/2) |0\rangle + e^{i\phi} \sin(\theta/2) |1\rangle$$

 The general state can be viewed as a point on the Bloch sphere



Differential Interference Contrast Microscopy



Two polarization states of light "visit" neighboring parts of a sample and interfere to produce contrast that depends on the phase difference between the paths.

The Quantum View



Increasing the Spin Echo Length



Cutting the slab into prisms allows the spatial separation of red and blue rays to be controlled by the separation of the prisms



Adding prisms with the opposite dispersion increases the separation of the red and blue rays

Spin echo small angle neutron scatteringa polarized neutron interferometer

• Quantum view



 ξ : spin echo length-max accessible length scale, which depends on, The separation of the two devices, *S* Inclination angle, θ Field intensity

The final neutron polarization is a measure of the degree of correlation between scattering from points separated by a distance ξ .

Conclusion:

NSE Provides a Way to Separate Resolution from Monochromatization & Beam Divergence

- The method currently provides the best energy resolution for inelastic neutron scattering (~ neV)
 - Both classical NSE and NRSE achieve similar energy resolution
 - NRSE is more easily adapted to "phonon focusing" I.e. measuring the energy line-widths of phonon excitations
- The method can also be used to improve (Q) resolution for elastic scattering
 - Extend size range for SANS (SESANS)
 - May allow 100 1000 x gain in measurement speed for some SANS exps
 - Separate specular and diffuse scattering in reflectometry
 - Measure in-plane ordering in thin films (SERGIS)?

Summary of NSE

- Field boundaries perpendicular to the neutron beam code neutron's speed (not it's trajectory)
 - Correlation time (up to ~ 300 ns on IN15) scales as λ^3
- Field boundaries inclined to neutron beam code neutron's trajectory
 - Correlation length (also called the spin echo length) scales as λ^2
 - Up to 25 microns correlation length has been achieved (Delft)
- Neither method requires highly collimated or very monochromatic neutron beams
 - NSE is suitable for long correlation times or large length scales typical of soft materials such as complex fluid films

THE END

Questions?

In a uniform magnetic field the neutron spin precesses around the the field

$$\frac{d\mathbf{S}}{dt} = \gamma_L \mathbf{S} \times \mathbf{B}$$
Bloch equation
$$\Phi = \frac{\gamma_L m\lambda}{h} \int_{path} Sign(B) |B| dl \propto \lambda \cdot FI$$





In a 10G field, 3cm of travel will rotate the spin of a 4Å neutron by a complete turn.

The "up" and "down" states have different Zeeman energies in a magnetic field

$$E_{Zeeman} = -\vec{\mu}_N \cdot \vec{E}$$

In NSE*, Neutron Spins Precess Before and After Scattering & a Polarization Echo is Obtained if Scattering is Elastic У $\ll \pi/2$ $\pi/2$ F π Allow spins to Rotate spins Elastic precess around z: Initially, to z and Scattering slower neutrons neutrons measure Event precess further over are polarized polarization a fixed path-length along z Rotate spins Allow spins to precess Rotate spins into through π about around z: all spins are in x-y precession plane x axis the same direction at the

Final Polarization, $P = \langle \cos(\phi_1 - \phi_2) \rangle$

echo point if $\Delta E = 0$ * F. Mezei, Z. Physik, 255 (1972) 145

Neutron Polarization is Measured using an Asymmetric Scan around the Echo Point



The echo amplitude decreases when $(Bd)_1$ differs from $(Bd)_2$ because the incident neutron beam is not monochromatic. For elastic scattering:

Echo Point

$$\langle P \rangle \sim \int I(\lambda) \cos \left[\frac{\gamma m}{h} \{ (Bd)_1 - (Bd)_2 \} \lambda \right] d\lambda$$

Because the echo point is the same for all neutron wavelengths, we can use a broad wavelength band and enhance the signal intensity

Field-Integral Inhomogeneities cause τ to vary over the Neutron Beam: They can be Corrected

 Solenoids (used as main precession fields) have fields that vary as r² away from the axis of symmetry because of end effects (div B = 0)



- According to Ampere's law, a current distribution that varies as r² can correct the field-integral inhomogeneities for parallel paths
- Similar devices can be used to correct the integral along divergent paths



Fresnel correction coil for IN15

An NRSE Triple Axis Spectrometer at HMI: Note the Tilted Coils





Neutron Spin Echo has significantly extended the (Q,E) range to which neutron scattering can be applied

Nanoscience & Biology Need Structural Probes for 1-100 nm





CdSe nanoparticles



Peptide-amphiphile nanofiber

10 nm holes in PMMA



Actin



Si colloidal crystal



Structures over many length scales in self-assembly of ZnS and cloned viruses



Thin copolymer films



- Any unscattered neutron (θ =0) experiences the same precession angles (ϕ_1 and ϕ_2) before and after scattering, whatever its angle of incidence
- Precession angles are different for scattered neutrons

 $\phi_{1} = \frac{KBd}{k\sin\chi} \text{ and } \phi_{2} = \frac{KBd}{k\sin(\chi + \theta)} \Rightarrow \cos(\phi_{1} - \phi_{2}) \approx \cos\left[\frac{KBd\cos\chi}{k\sin^{2}\chi}\theta\right]$ $P = \int dQ.S(Q).\cos\left[\frac{KBd\cos\chi}{k^{2}\sin^{2}\chi}Q\right] \qquad \text{Polarization proportional to}$ Fourier Transform of S(Q) $Spin Echo Length, \quad \zeta = KBd\cos\chi/(k\sin\chi)^{2}$

How Large is the Spin Echo Length for SESANS?

Bd/sinχ (Gauss.cm)	λ(Angstroms)	χ(degrees)	ζ (Angstroms)
3,000	4	20	1,000
5,000	4	20	1,500
5,000	6	20	3,500
5,000	6	10	7,500

It is relatively straightforward to probe length scales of ~ 1 micron

Measuring Correlations in Space & Time



- Use space and time resolving NSE simultaneously
- Large (~ 1 T.m) solenoidal field: $\tau = 186 B_0.L. \lambda^3$ ns (λ in nm)
- Triangular solenoids (~ 0.01 T): r = 147,000 B.d cot θ . λ^2 nm
- At a reactor source, we could set r and τ independently
- At a pulsed source we could keep B λ^2 constant by varying B with time after a pulse: TOF then maps out τ (might be hard)
- This would be a completely different way of measuring structural correlations using scattering – (r,τ) instead of (Q,ω)