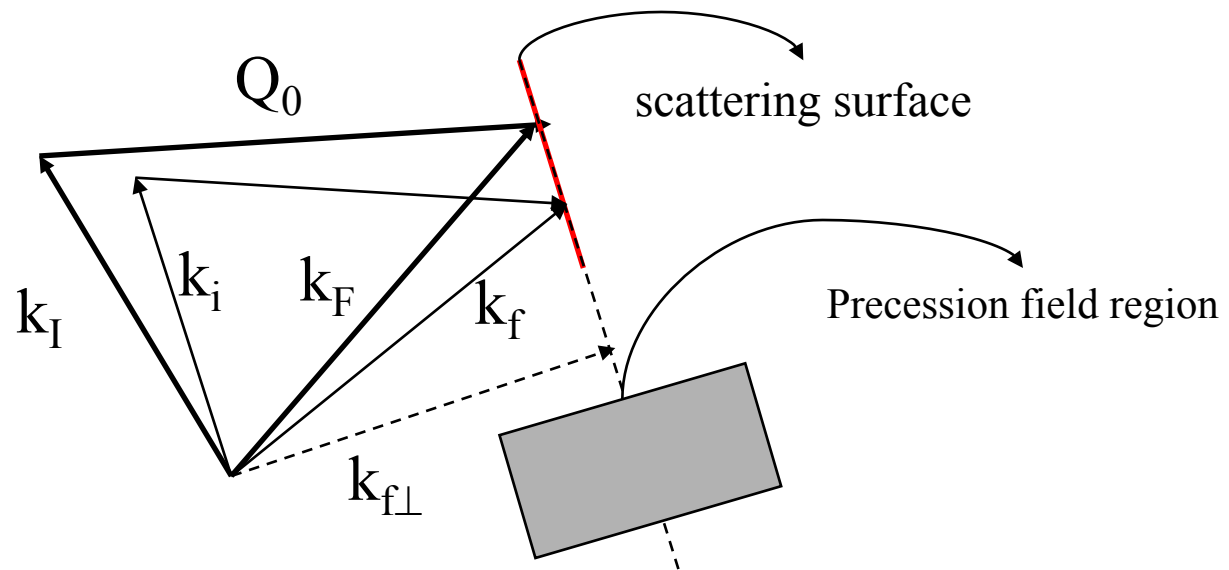


# Larmor Labelling of Scattering Angles

Roger Pynn

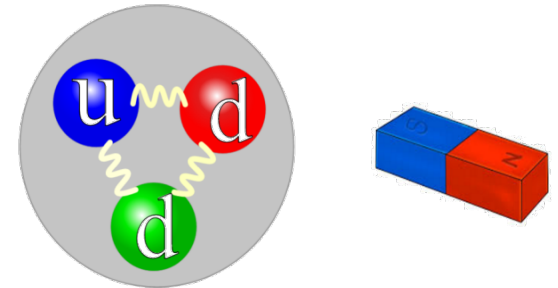
Indiana University, Bloomington &  
Oak Ridge National Laboratory

# The Picture I made for Reinhard Scherm



But I digress....

# The neutron is a spin-1/2 particle with a magnetic moment

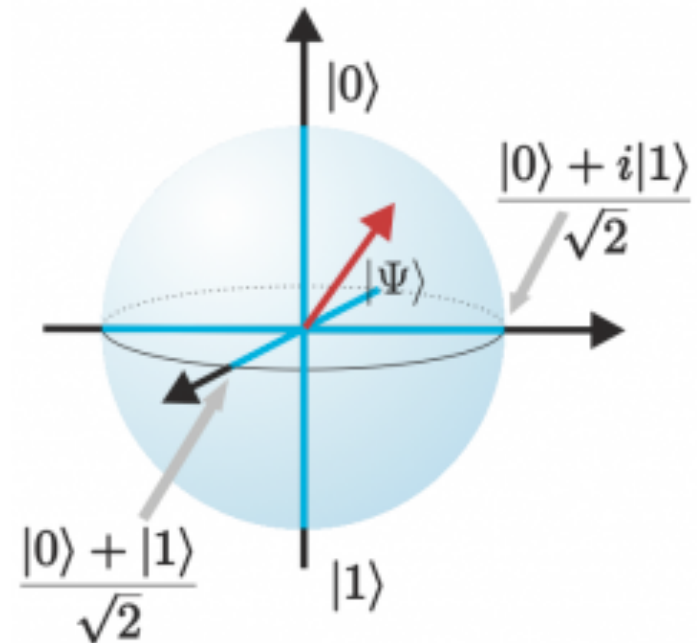


- In an applied magnetic field,  $B_z$ , the neutron has two spin eigenstates, “up” and “down” denoted  $|0\rangle$  &  $|1\rangle$ 
  - Magnetic moment either along  $z$  or  $-z$  (magnetic moment is antiparallel to the spin)
- Quantum mechanics tells us that any neutron spin state is a linear combination of the two eigenstates

$$\chi = a|0\rangle + b|1\rangle \quad \text{where } |a|^2 + |b|^2 = 1$$

$$\chi = \cos(\theta/2)|0\rangle + e^{i\phi} \sin(\theta/2)|1\rangle$$

- The general state can be viewed as a point on the Bloch sphere



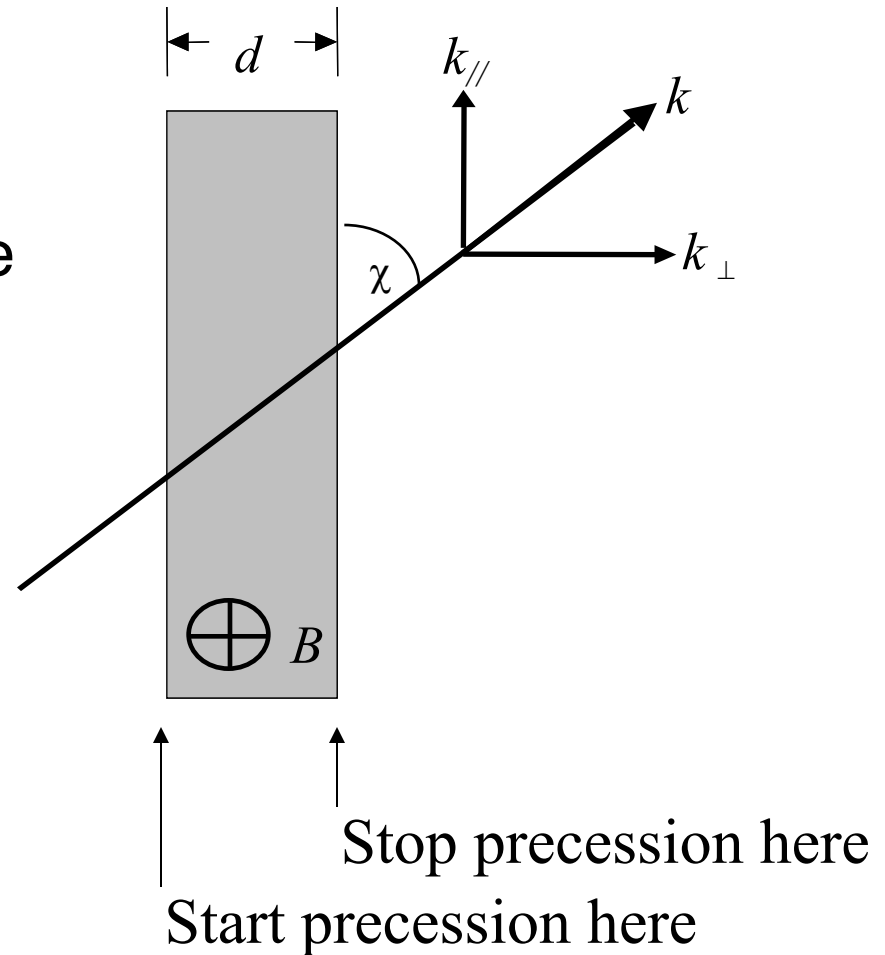
# By “Tilting” the Precession-Field Region, Spin Precession Can Be Used to Code a Specific Component of the Neutron Wavevector

If a neutron passes through a rectangular field region at an angle, its total precession phase will depend only on  $k_{\perp}$ .

$$\omega_L = \gamma B$$

$$\phi = \omega_L t = \gamma B \frac{d}{v \sin \chi} = \frac{KBd}{k_{\perp}}$$

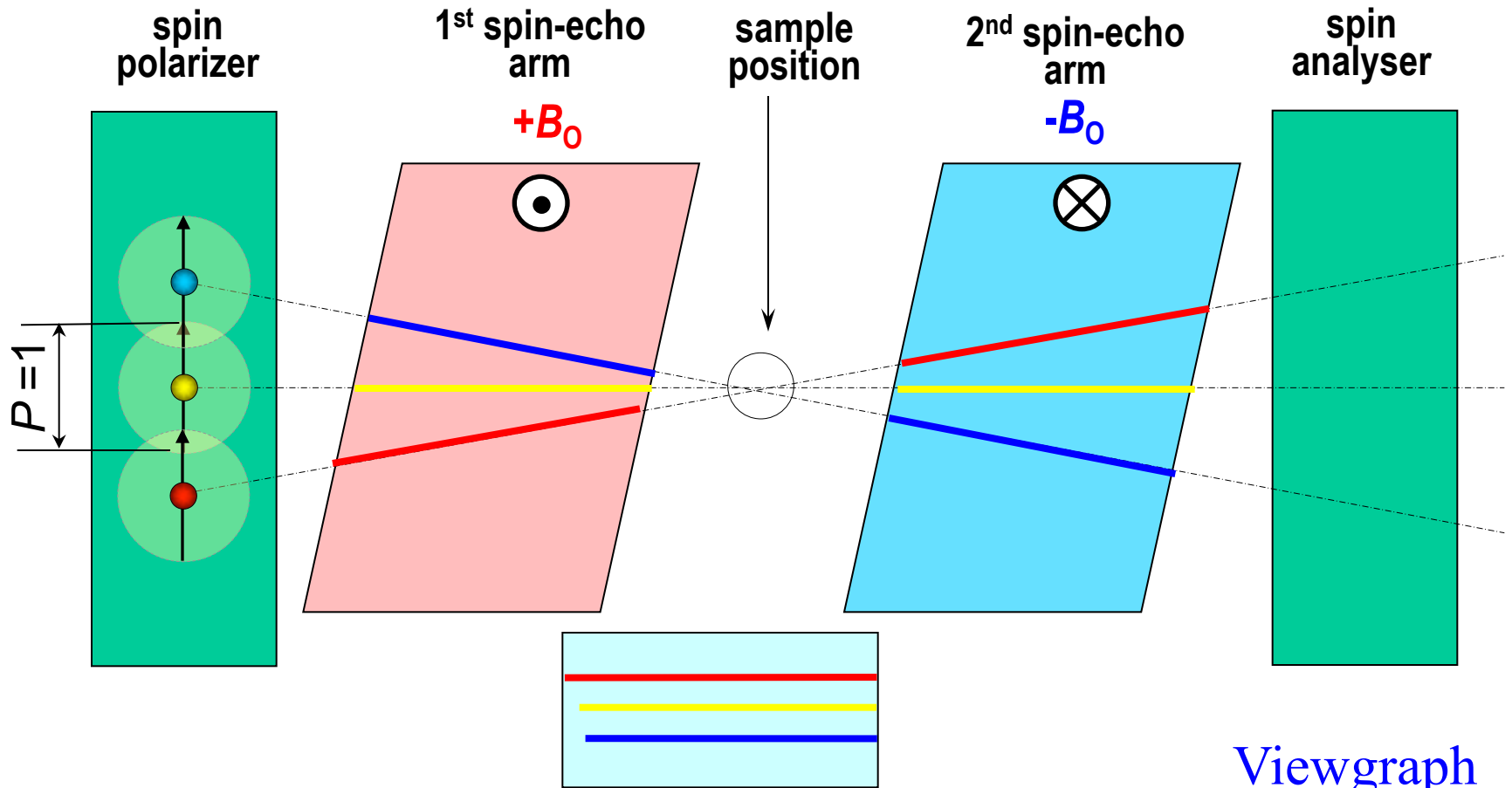
with  $K = 0.291 \text{ (Gauss.cm.}\text{\AA)}^{-1}$



Roughly 1 turn per 30 Gauss.cm for a 4  $\text{\AA}$  neutron



# The Classical Picture of SESANS

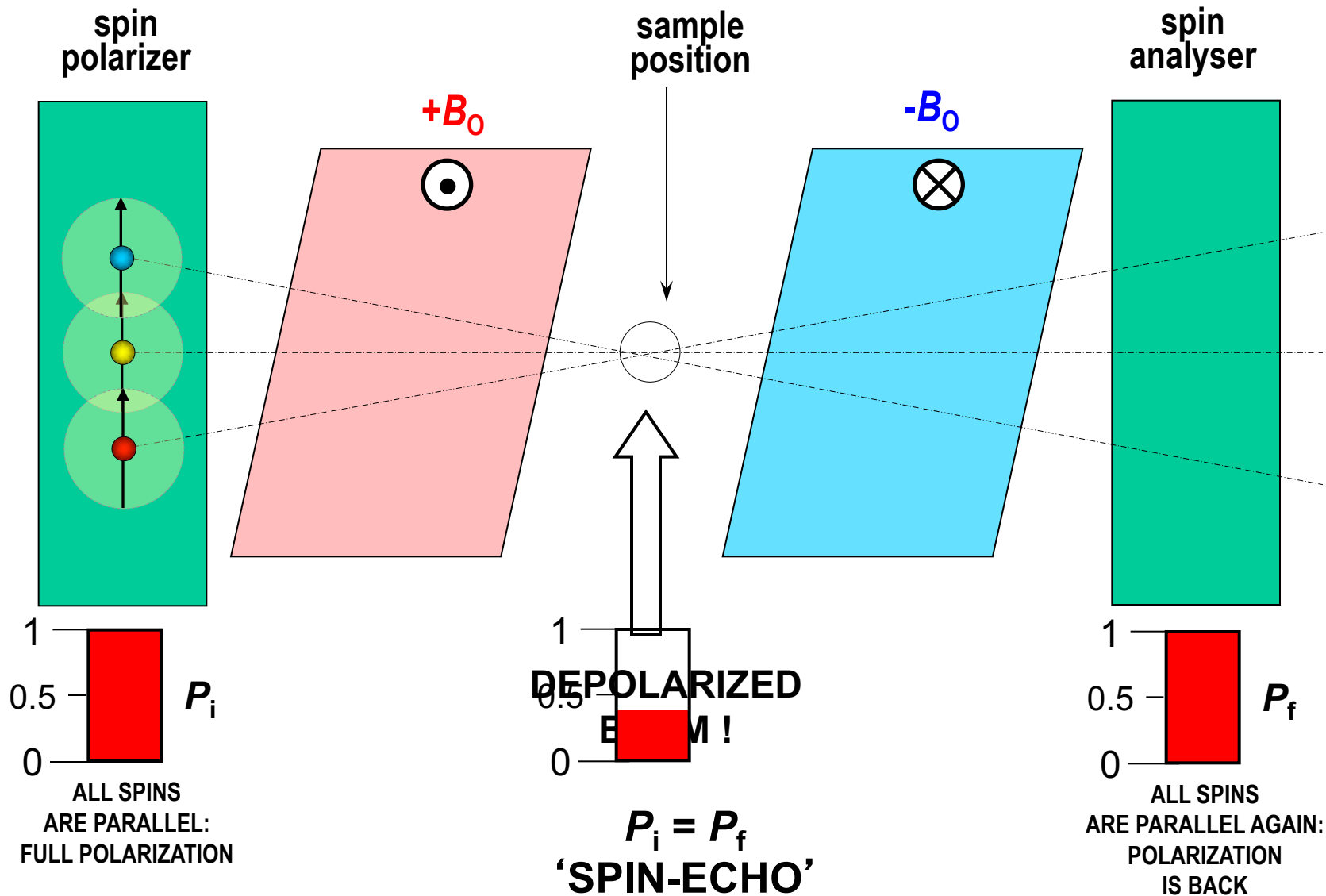


**DIFFERENT PATH LENGTHS FOR THE DIFFERENT TRAJECTORIES !**

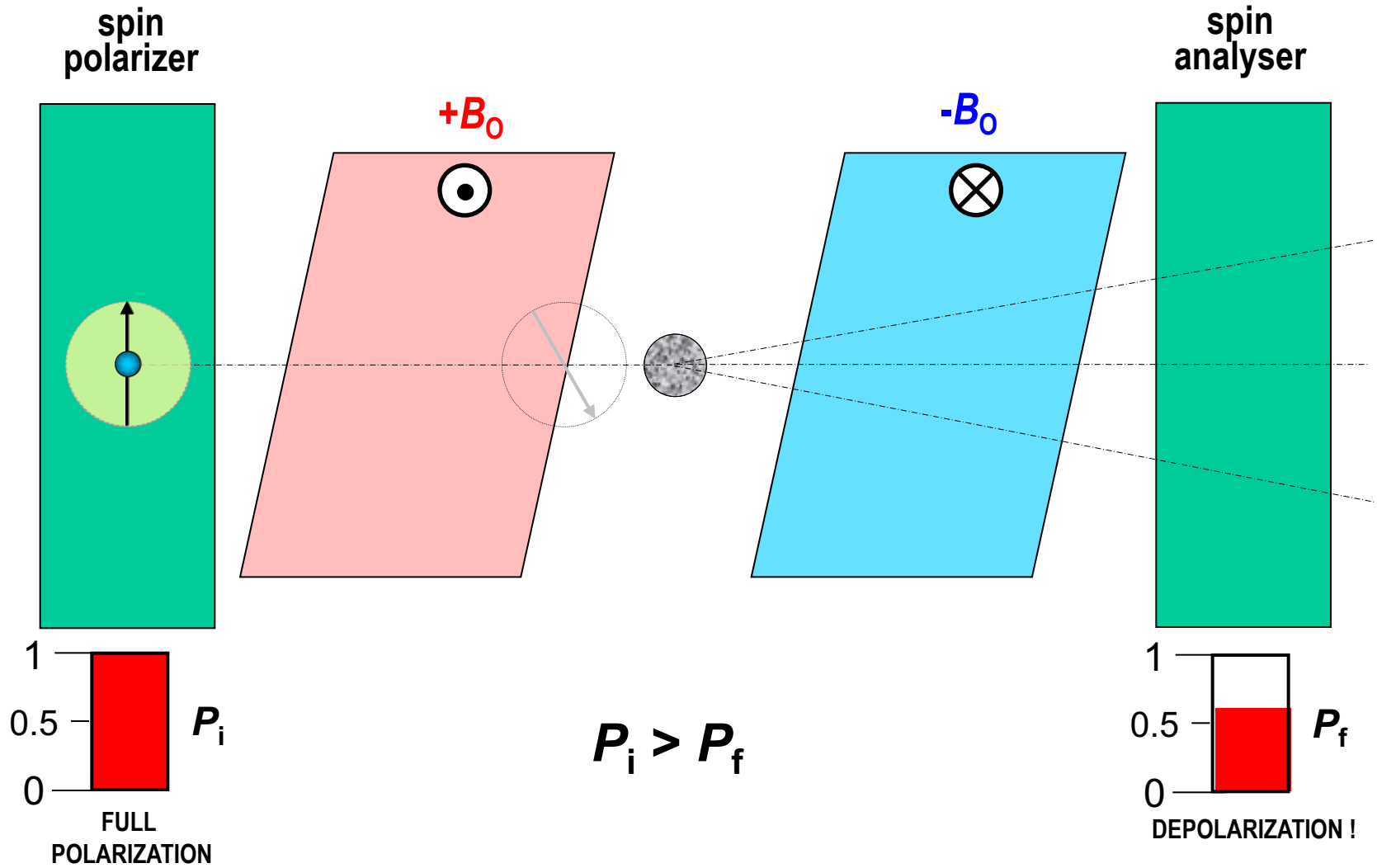
Viewgraph  
sequence by  
A. Vorobiev

# Spin Echo Scattering Angle Measurement (SESAME)

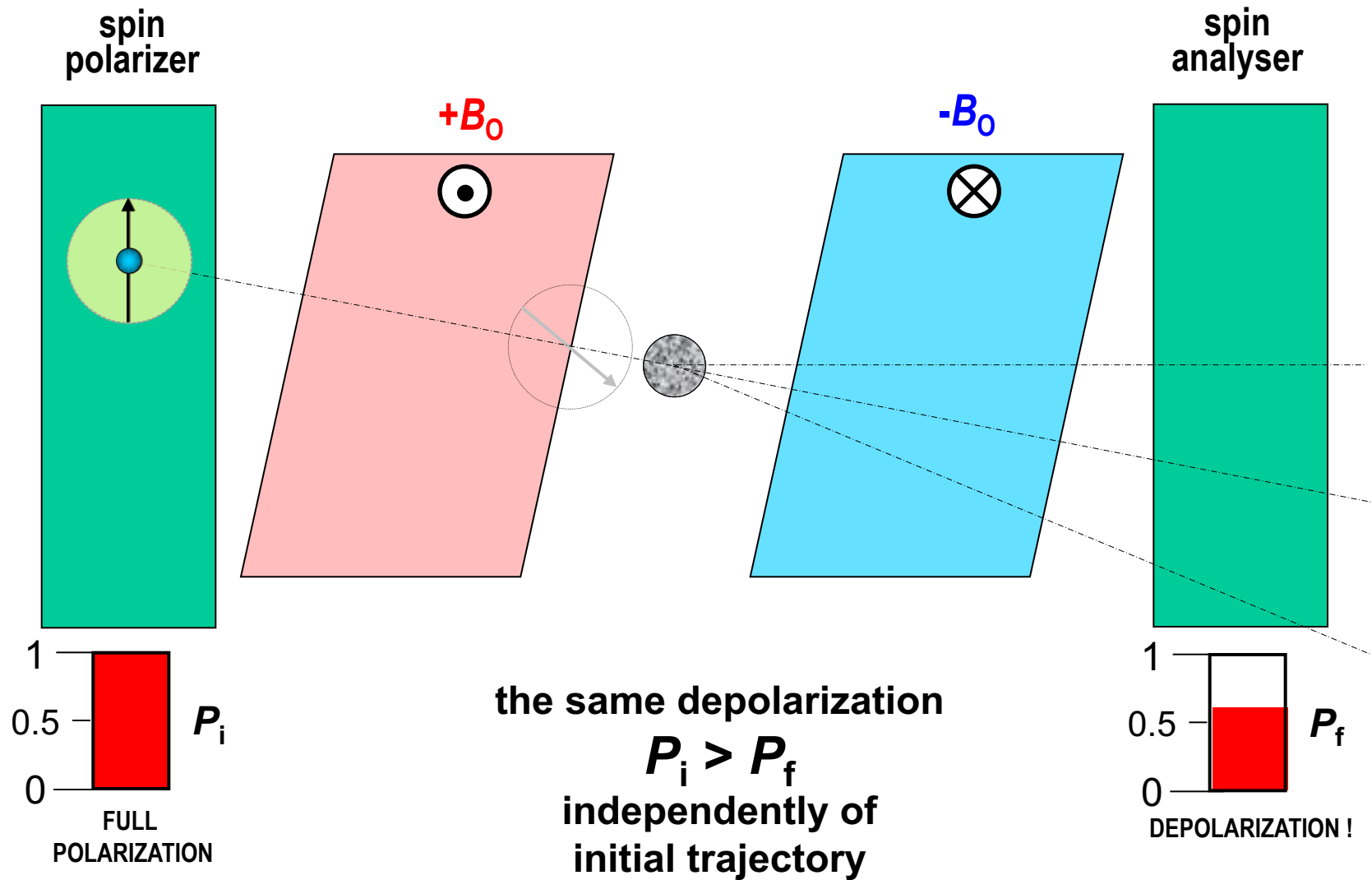
No Sample in Beam



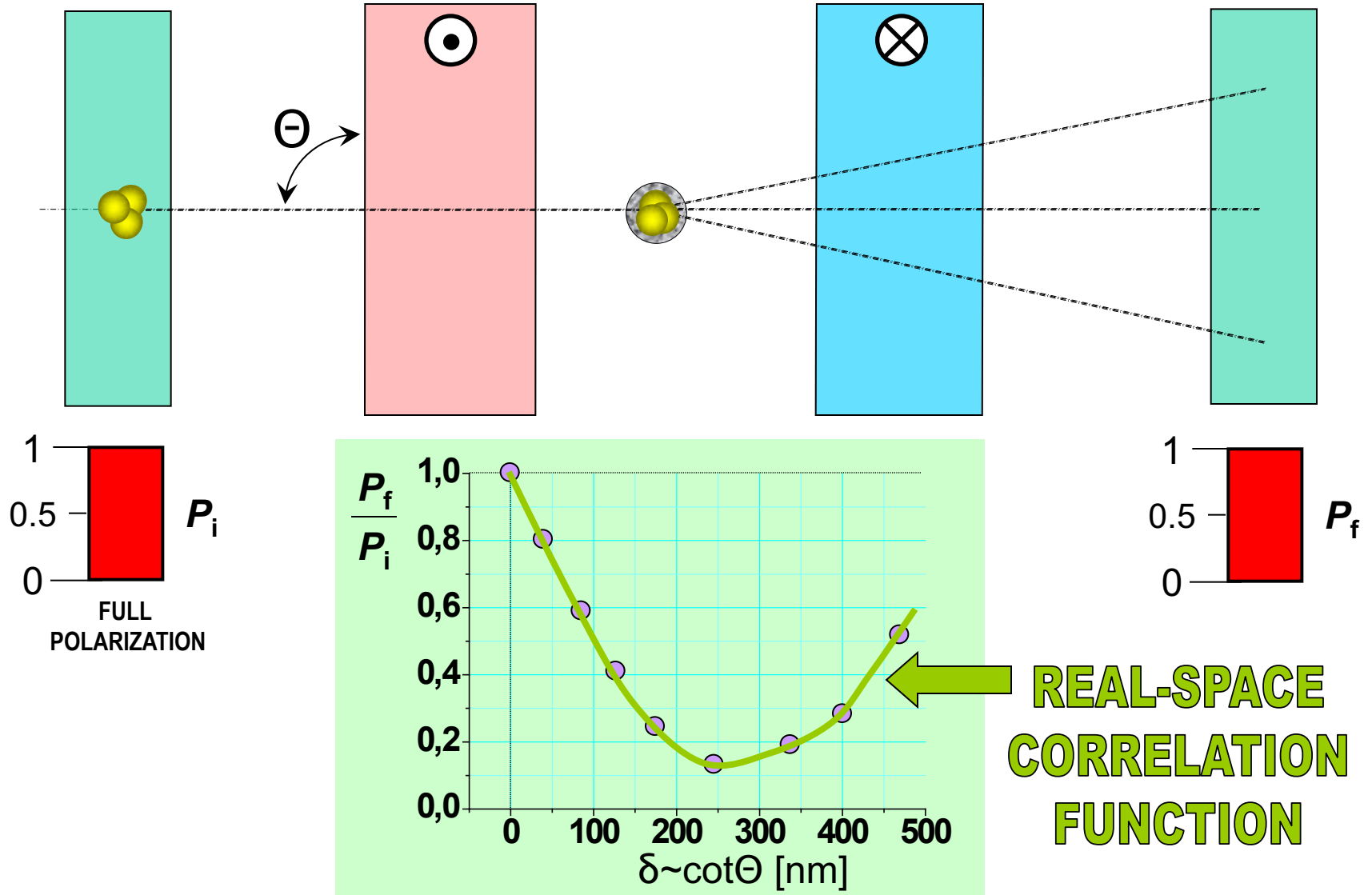
# Spin Echo Scattering Angle Measurement (SESAME) Scattering by the Sample



# Spin Echo Scattering Angle Measurement (SESAME) Scattering of a Divergent Beam

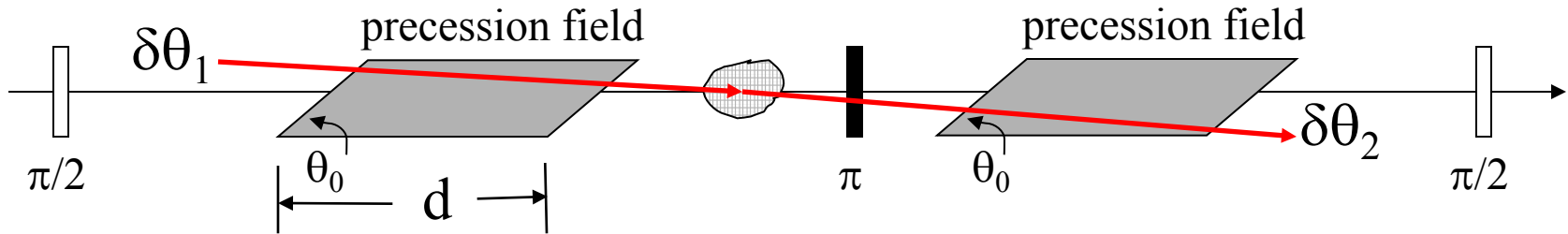


# Spin-echo angular encoding: the experiment



# NSE Angle Coding Illustrated for SESANS

- Make the number of spin precessions depend on the neutron's direction of travel instead of (only) its speed.....



$$\omega_L = \gamma B$$

$$\phi_1 = \omega_L t_1 = c B d \lambda (1 - \cot \theta_0 \cdot \delta \theta_1) \text{ where } c = \frac{\gamma m}{h} = 4.635 \times 10^{14} T^{-1} m^{-2}$$

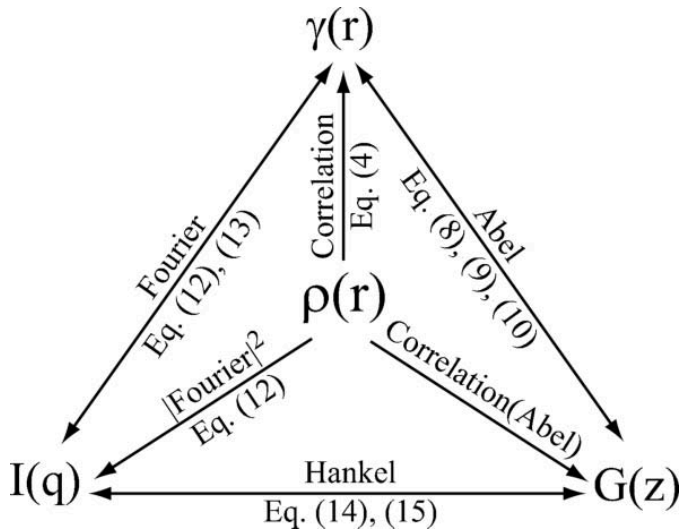
$$\phi_1 - \phi_2 = Q \frac{c B d}{4\pi} \lambda^2 \cot \theta_0 = Q \cdot \zeta$$

$$\frac{P}{P_0}(Z) = \frac{\int_0^\infty dQ_y \int_0^\infty dQ_z \frac{d\sigma}{d\Omega} \cos(Q_z Z)}{\int_0^\infty dQ_y \int_0^\infty dQ_z \frac{d\sigma}{d\Omega}} = e^{\Sigma[G(Z)-1]}$$

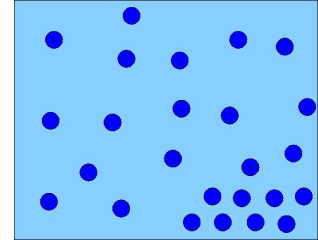
$\zeta$  is called the spin echo length

$\Sigma$  = total probability for single scattering;  $G(Z)$  is a correlation function

# Correlation Function



Local Density Distribution  $\rho(\vec{r})$



Debye Correlation Function

$$\gamma(r) = \frac{1}{V} \left\langle \int_V \rho(\vec{r}') \rho(\vec{r}' + \vec{r}) d^3 \vec{r}' \right\rangle$$

Andersson et al J. Appl. Cryst. (2008). 41, 868–885

SANS

SESANS

$$I(Q) = \frac{d\Sigma}{d\Omega}(Q) = 4\pi \int_0^\infty \gamma(r) J_0(Qr) r^2 dr$$

$$G(z) = 2 \int_z^\infty \gamma(r) \frac{r}{\sqrt{r^2 - z^2}} dr$$

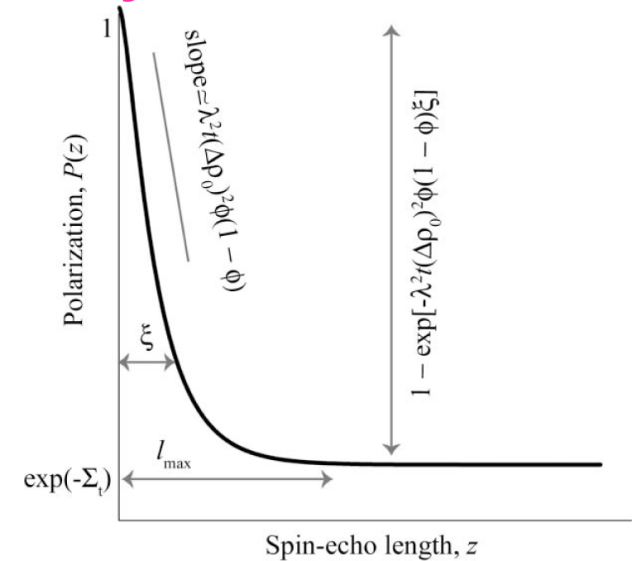
Fourier

Abel

# What we Actually Measure with SESANS for a Bulk Two-Phase System

$$\frac{1}{\lambda^2} \ln \left( \frac{P(z)}{P_0} \right) = \Sigma_t [G(z) - 1]$$

$$\text{since } G(\infty) = 0, \quad \Sigma_t = -\frac{1}{\lambda^2} \ln \left( \frac{P(\infty)}{P_0} \right)$$

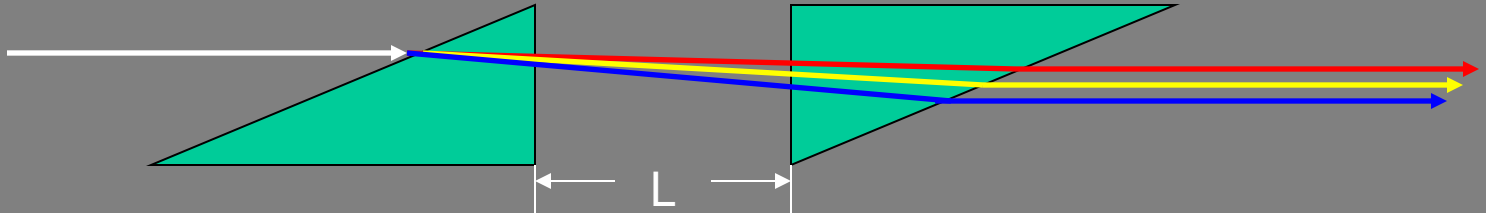


also,  $\Sigma_t = t \Delta \rho^2 \phi(1 - \phi) \xi$  where the correlation length  $\xi = 2 \int \gamma(r) dr$

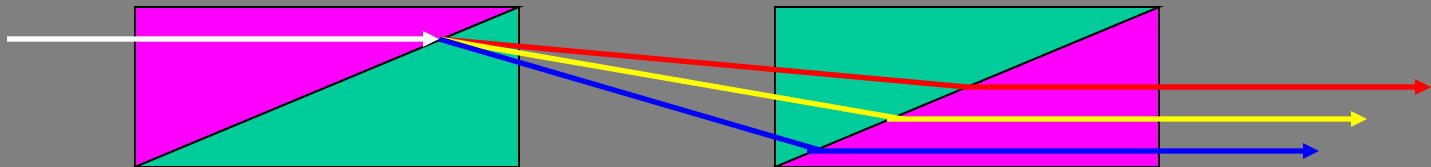
- $G(z)$  is a projection of the usual Debye density-density correlation function: it is a **real space correlation function**
- The value of  $z$  where  $\ln(P)/\lambda^2$  becomes flat tells us the maximum size of density correlations in the system (particle diameter for dilute spherical objects)
- The value of  $\ln(P)/\lambda^2$  at  $z \rightarrow \infty$  measures the total scattering



# Increasing the Spin Echo Length

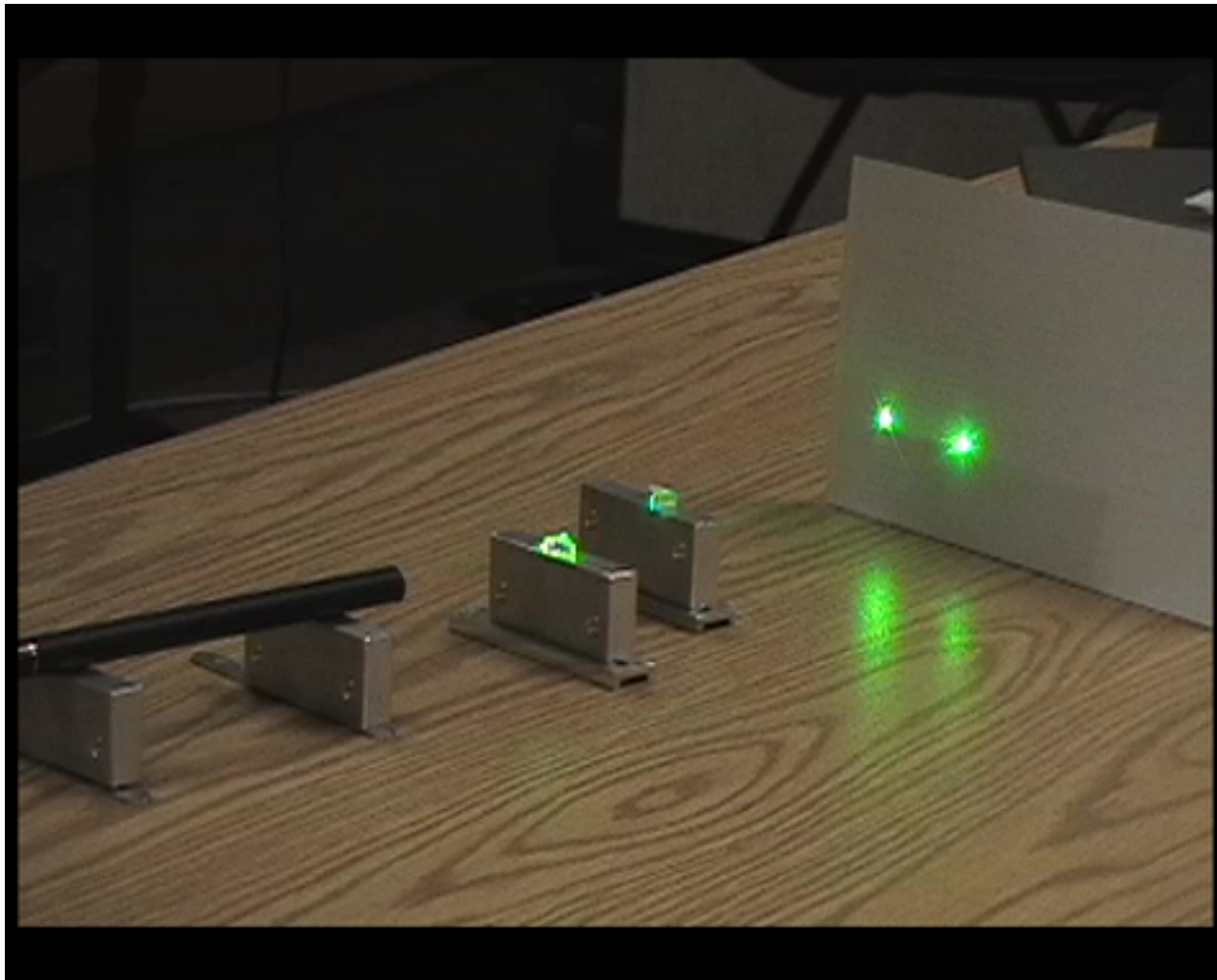


Cutting the parallelogram into prisms increases the length of the parallelogram ( $d$  in the previous VG) and adds a rectangular region that does no angle coding. In QM language, it allows the spatial separation up (red) and down (blue) states to be controlled by the separation of the prisms



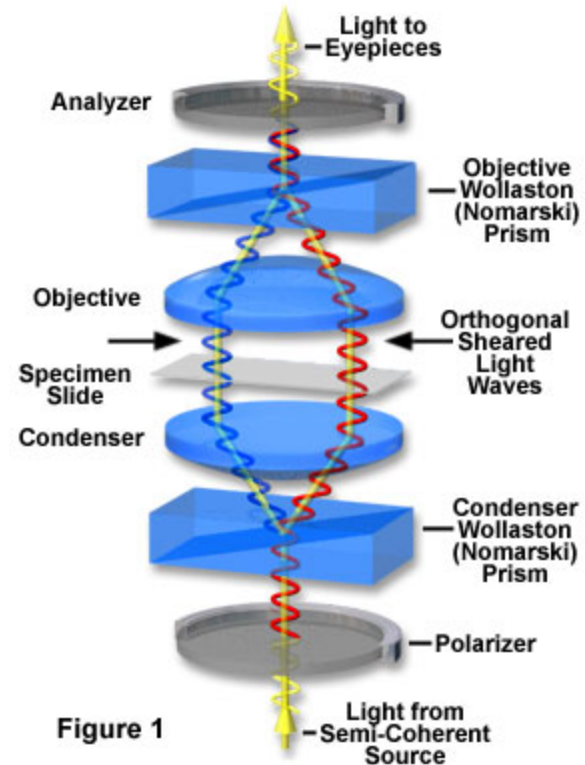
Adding prisms with the opposite dispersion (i.e. opposite magnetic fields) doubles the separation of the red and blue rays (i.e. doubles the spin echo length)

# The Quantum View

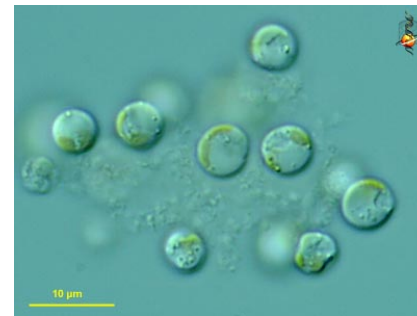


# Differential Interference Contrast Microscopy

Differential Interference Contrast Schematic



on Prism

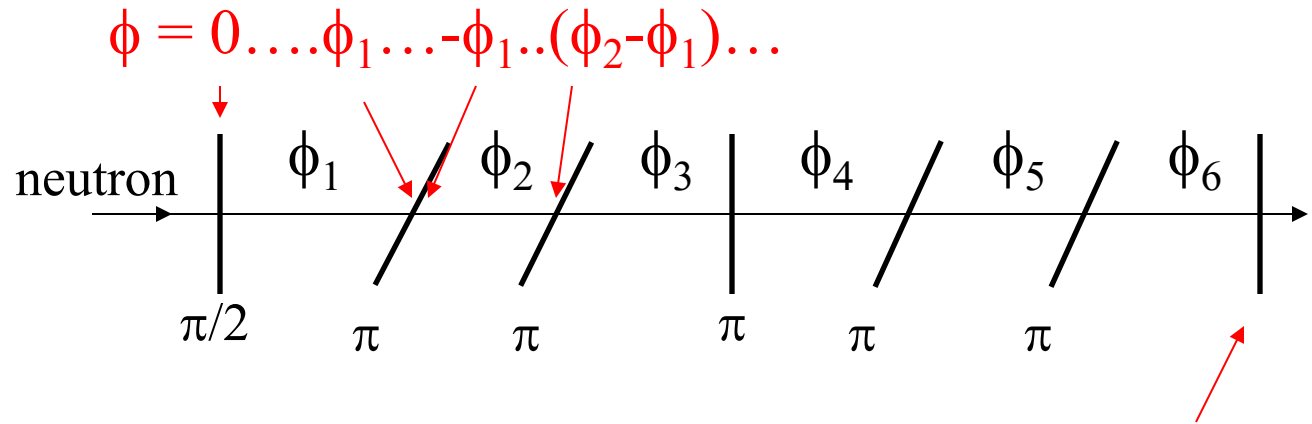
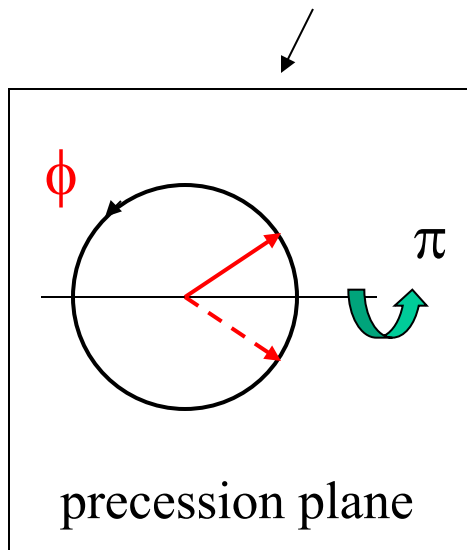


Eukaryotic Algae

Two polarization states of light “visit” neighboring parts of a sample and interfere to produce contrast that depends on the phase difference between the paths.

# Using Sequential $\pi$ Flippers to Implement SESANS

An element that performs a  $\pi$  rotation about an axis in the precession plane changes the sign of prior precession angles

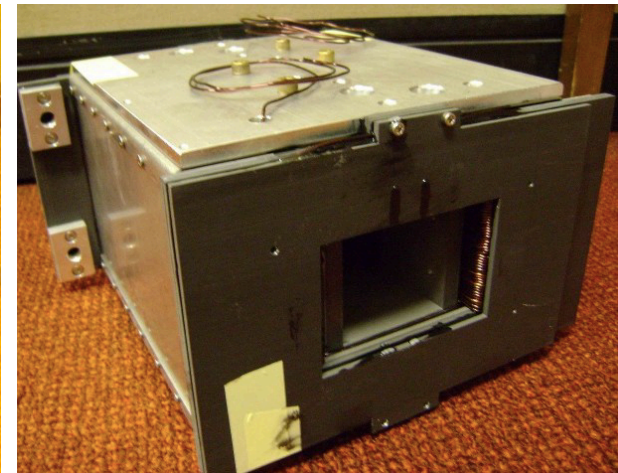
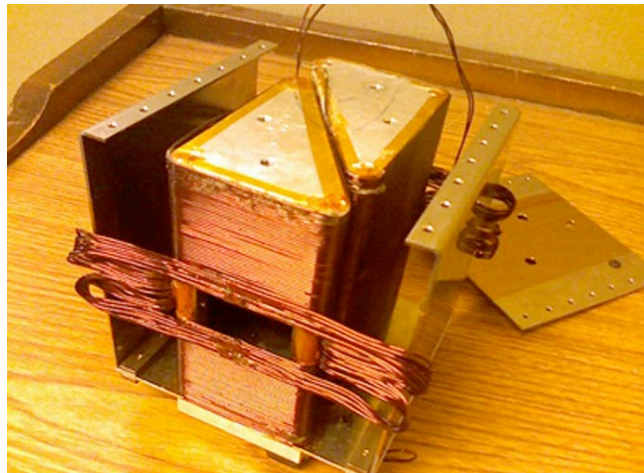


$$\text{Net Precession} = -\phi_1 + \phi_2 - \phi_3 + \phi_4 - \phi_5 + \phi_6$$

- Total net precession angle =  $(\phi_4 + \phi_5 + \phi_6) - (\phi_1 + \phi_2 + \phi_3) - 2(\phi_5 - \phi_2)$
- The first two terms are the familiar spin-echo terms
  - Depend only on neutron velocity as shown here
- Last term depends on velocity *and* angle of neutron trajectory
  - Picks up an extra factor of 2 compared to using a parallelogram-shaped field
  - The Wollaston prism is a pi-flipper

# How can we realize the magnetic Wollaston prism in practice?

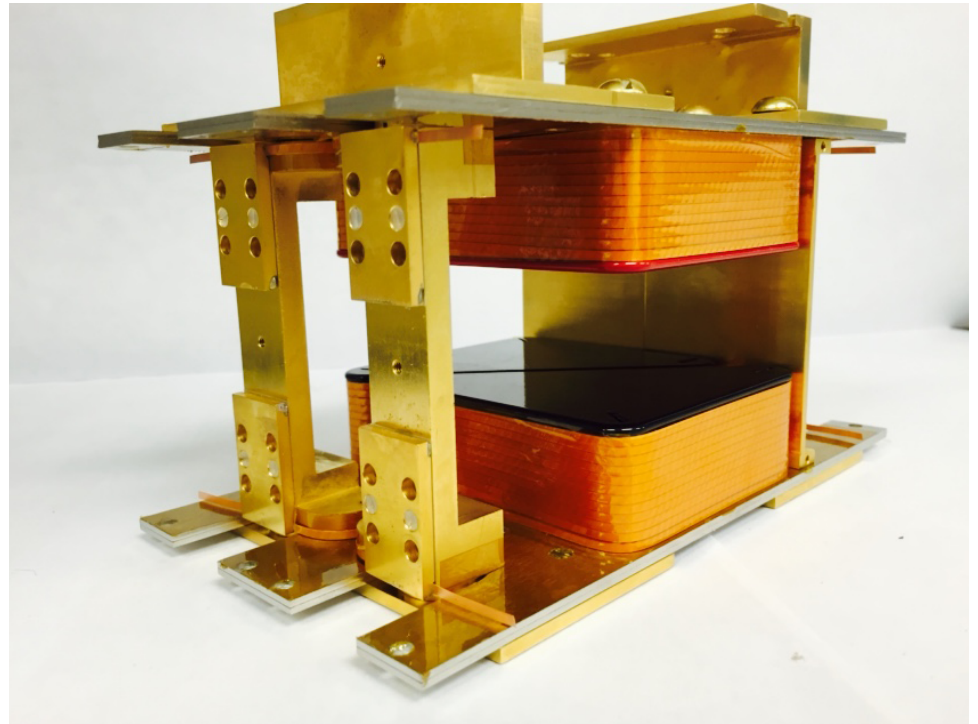
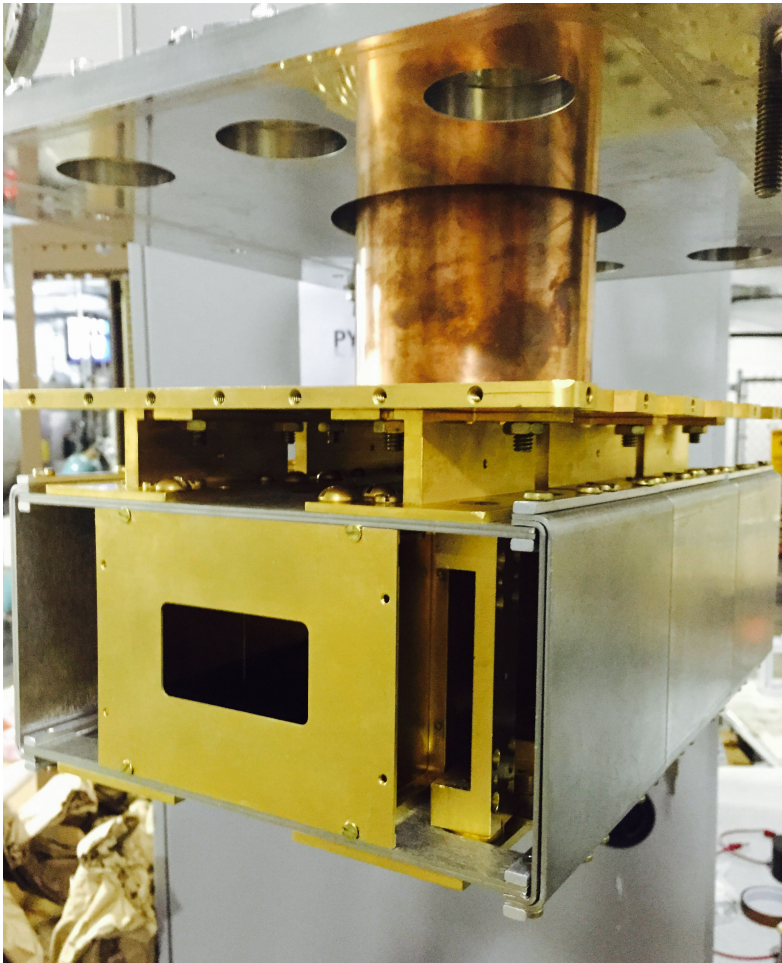
- We built several generations of air- and water-cooled WPs and (re)discovered the importance of using symmetry ( $\pi$  flip) to cancel decoherence effects (Larmor phase aberrations)
- Ultimately, room temperature coils are limited by:
  - Flatness of coils (waviness, discrete wires)
  - Neutron absorption by wires
  - Maximum achievable magnetic field



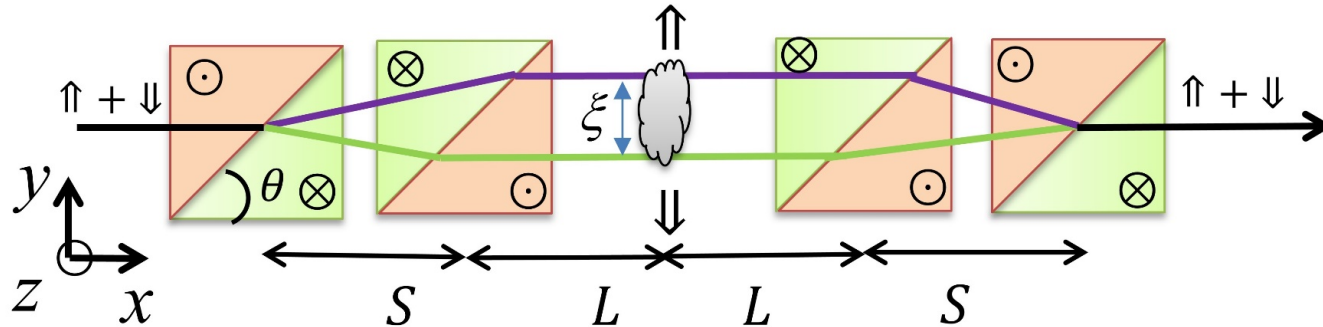
Paul Stonaha



Finally we built Wollaston prisms using HTS tape and HTS films surrounding the field regions to achieve high fields & good uniformity



# Spin echo small angle neutron scattering- a polarized neutron interferometer



$$\zeta = c\lambda^2 B S \cot\theta / 2\pi \quad c = 1.47 \times 10^{14} T^{-1} m^{-2}$$

$$\zeta = 590 \text{ nm for } \lambda = 4 \text{ \AA}, B = 0.1 \text{ T \& } S = 0.5 \text{ m}$$

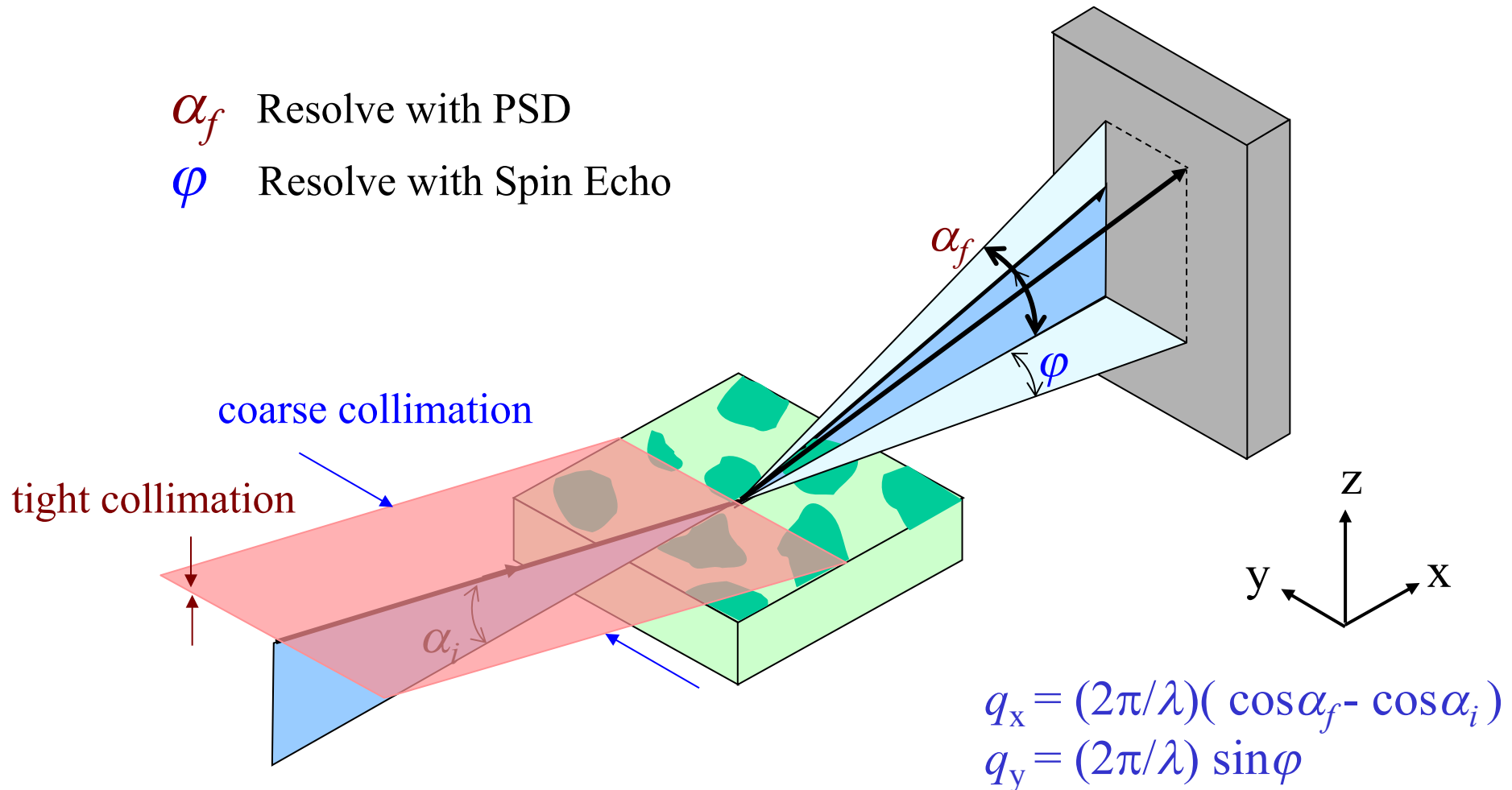
$\xi$ : spin echo length-max accessible length scale, which depends on,

- The separation of the two devices,  $S$
- Inclination angle,  $\theta$
- Field intensity

The final neutron polarization is a measure of the degree of correlation between scattering from points separated by a distance  $\xi$ .

# SESANS in reflectometry - SERGIS

- $\alpha_f$  Resolve with PSD
- $\phi$  Resolve with Spin Echo

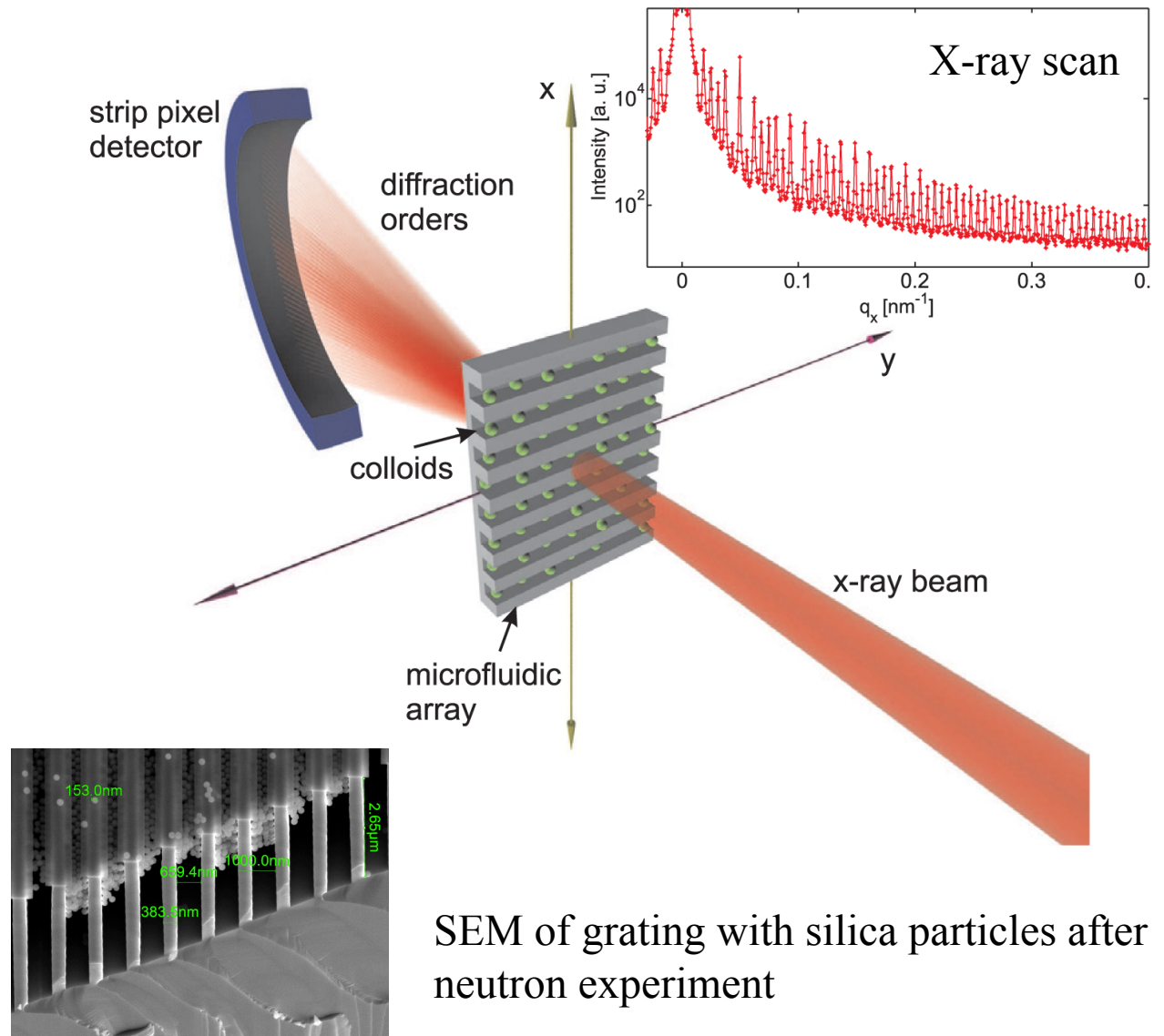


Note that scattering at a particular value of  $q$  is spread over larger angles  $\alpha_f$  than  $\phi$



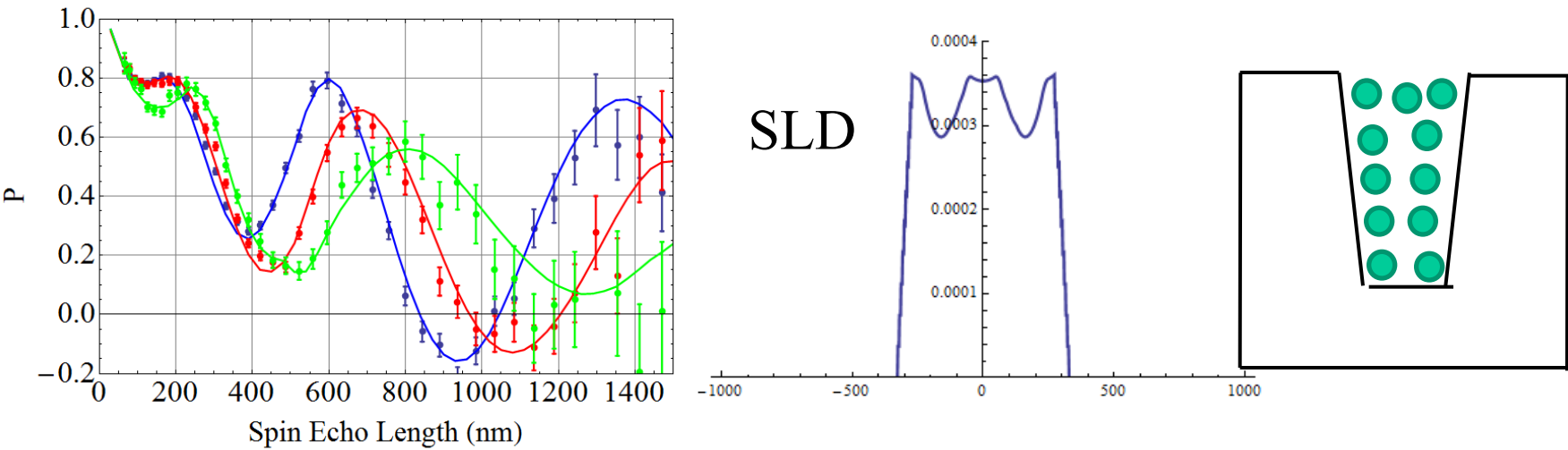
# X-Ray & Neutron Grating Holography

- Grating creates near-field modulation of amplitude & phase
- Contents of grooves act as a weak-phase scattering object
- Deconvolve pattern produced by bare grating to deduce structure in grooves
- X-rays measure Bragg peaks & calculate Patterson function
- Neutrons measure Patterson function using SESANS

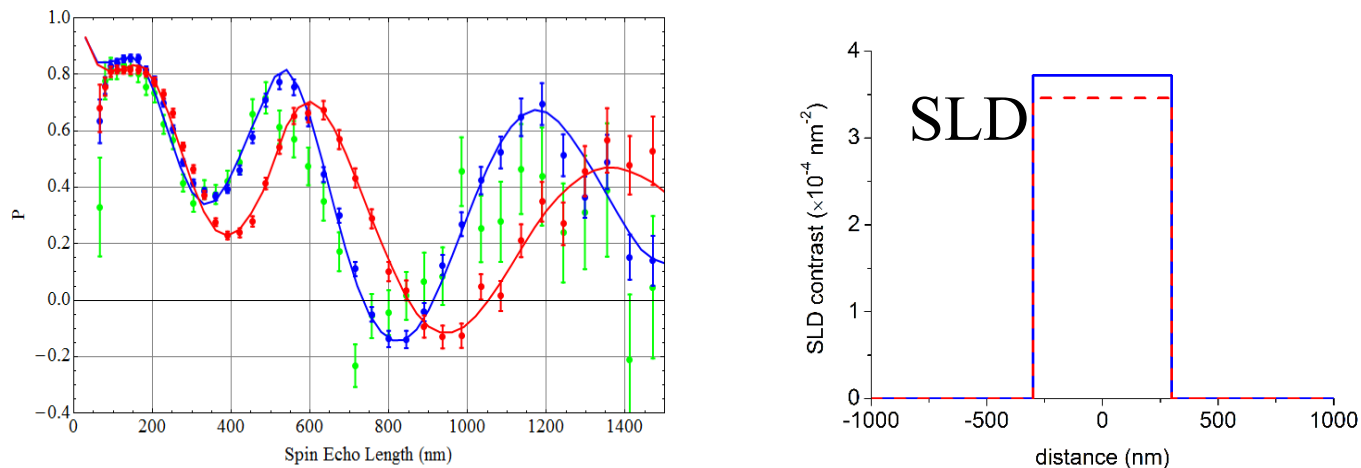


SEM of grating with silica particles after neutron experiment

# Neutron Holography Results



Above: Blue: carrier fluid only; Red: with colloid at pH = 9.7; Green: same as red 24 hours later  
Below: Blue: carrier fluid only; Red: with colloid at pH = 10.6

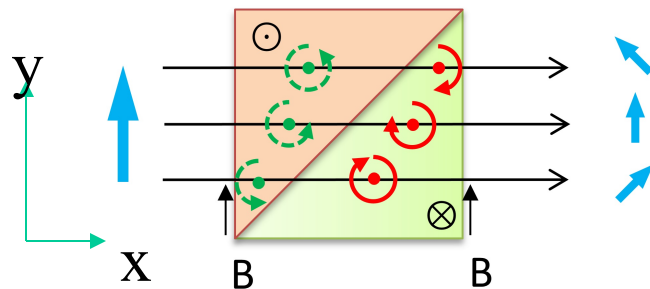


# Magnetic Wollaston Prisms

- Label a neutron's distance from the beam center ( $y-y_0$ ) by its Larmor precession angle
- Each prism consists of two triangular magnetic-field regions with oppositely directed B fields and sharp boundaries

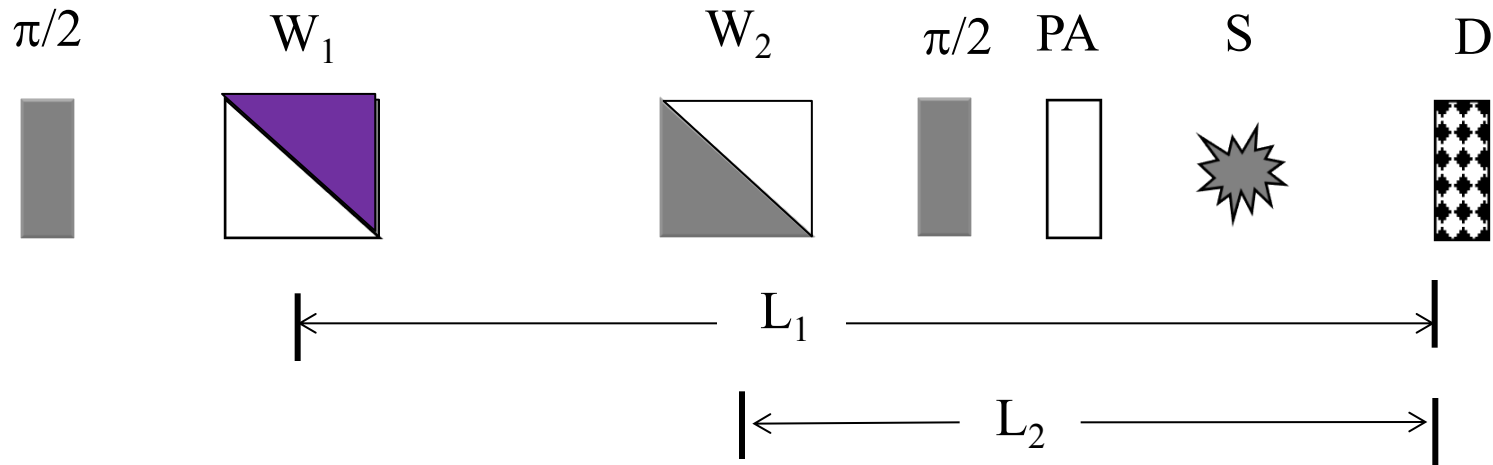
$$\phi_{total} = 2cB\lambda(y - y_0)cot\theta$$

$\sim 4$  rads/mm for 4 Å neutrons and  $B = 100$  G



If a polarization analyzer is placed after the prism, the measured intensity varies as  $1 + \cos(\phi_{total})$

# Spin Echo Modulated SANS (SEMSANS)



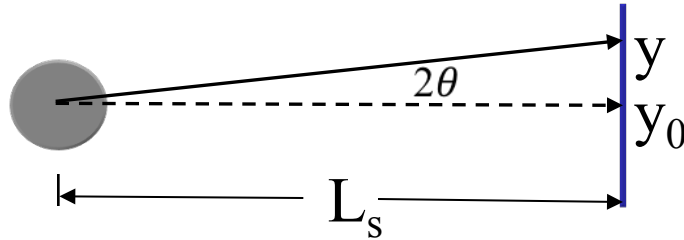
- Consider a neutron arriving at a point on the detector at position  $y$  (relative to the center) with divergence angle,  $\theta$

$$\varphi_1 = c\lambda B_1(y + L_1\theta)$$

$$\varphi_2 = -c\lambda B_2(y + L_2\theta)$$

- The total phase is independent of  $\theta$  if  $B_1L_1 = B_2L_2$  so we get a cleanly modulated intensity on the detector **even with a divergent beam**
- Period is:  $p = \pi \cdot \tan\theta / c\lambda(B_2 - B_1)$
- I.e  $p \sim 1.7$  mm for 4 Å neutrons and  $B_2 - B_1 = 100$  G

# What do we measure with SEMSANS?



$$y - y_0 = 2\theta \cdot L_s = \frac{Q\lambda}{2\pi}$$

$$P(y) = \int_{det} \cos\left(\frac{2\pi y_0}{p}\right) \frac{d\sigma}{d\Omega} dy_0 / \int_{det} \frac{d\sigma}{d\Omega} dy_0$$

$$P(y) = \frac{\int_{det} \cos\left(\frac{2\pi(y - \frac{Q\lambda L_s}{2\pi})}{p}\right) \frac{d\sigma}{d\Omega} dQ}{\int_{det} \frac{d\sigma}{d\Omega} dQ} = P(y_0) \frac{\int_{det} \cos(Q \zeta) \frac{d\sigma}{d\Omega} dQ}{\int_{det} \frac{d\sigma}{d\Omega} dQ}$$

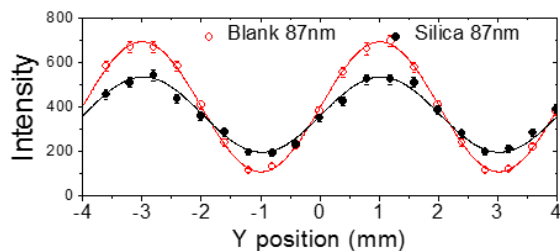
With  $\zeta = \frac{\lambda L_s}{p}$  as the spin echo length

# SEMSANS for a homogeneous sample: every point on the detector measures the same thing

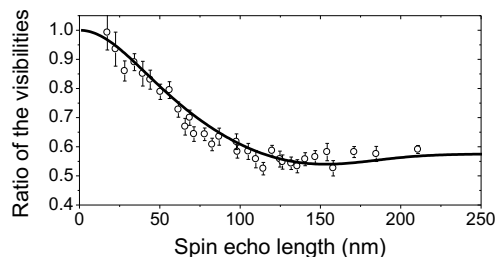
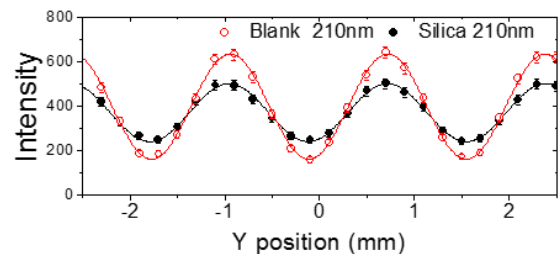
data obtained at NCNR

$$\frac{P_s}{P_0} = e^{\Sigma_t(G(\xi)-1)} \quad \xi = \lambda L_s / p$$

(a)



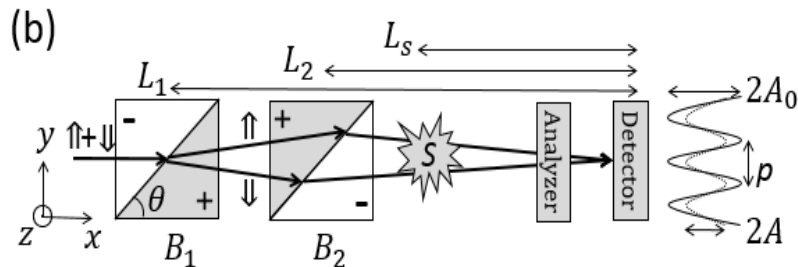
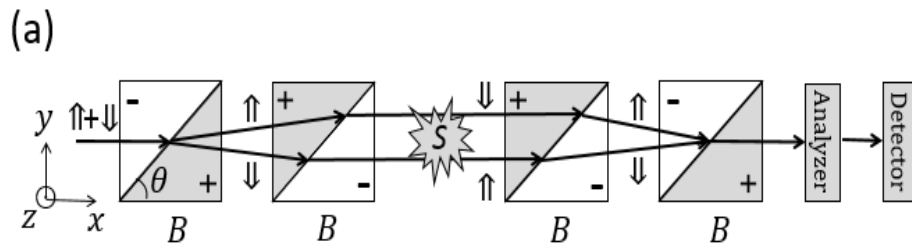
(b)



silica colloid (NCNR)

- $P_s$  = polarization at any point on detector with sample
- $P_0$  = polarization without sample
- $\Sigma_t$  is the fraction of the beam (single) scattered by the sample
- $G(Z)$  is a real space correlation function measured at a spin echo length,  $Z$ . It is the same function as is measured using SESANS
- Data analysis can be tricky because of zero-crossing of polarization
- Can also implement dark-field radiography this way
  - each point on the detector measures a SESANS curve for a small area of the scattering object

# Comparison of SESANS & SEMSANS



## SESANS

- Four Wollaston prism needed
- All the fields are equal and balanced
- Paths at the sample are parallel
- Non-magnetic sample environment
- Low resolution detector is OK

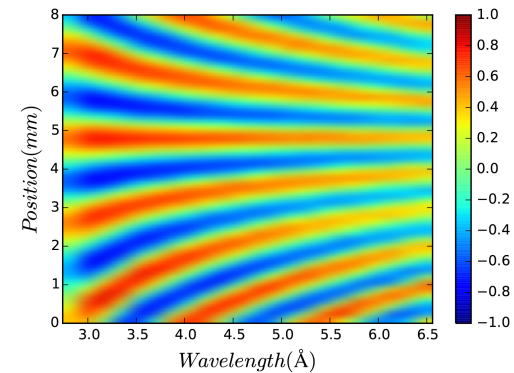
## SEMSANS

- Two Wollaston prism needed
- The fields are not balanced
- Focused towards the sample
- Sample environment relaxed
- High resolution detector is required

They both measure the correlation function of the sample.



# Double WP unit on LARMOR with $^3\text{He}$ polarization analyzer and $55\text{-}\mu\text{m}$ -pixel, $28\times 28\text{ mm}^2$ Tremsin Detector



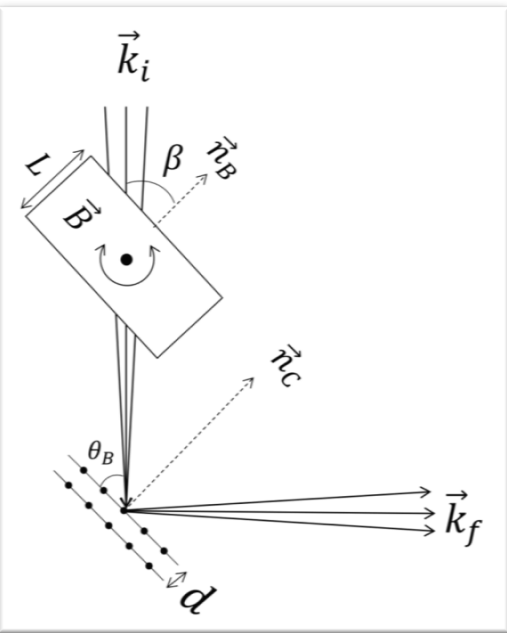
Polarization projected on to horizontal coordinate as a function of position and wavelength.

Fringe period is chromatic  
Visibility is achromatic  
(opposite of grating-based far-field interferometer)

Independent of beam divergence – would work with a lens



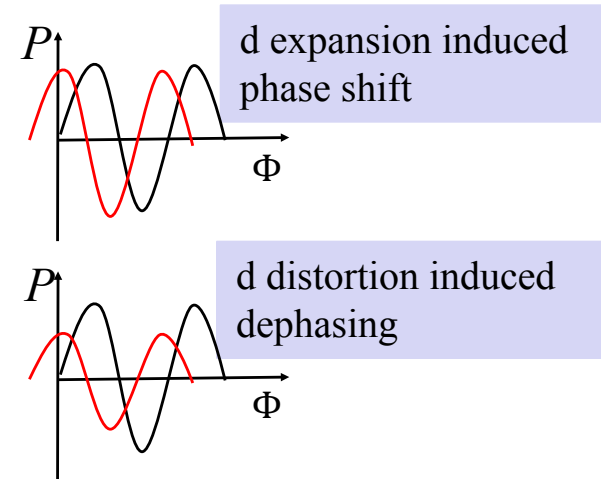
# Larmor diffraction to achieve higher resolution



$$\Phi = \frac{\gamma_N m B L}{h k_{i,\perp}}$$

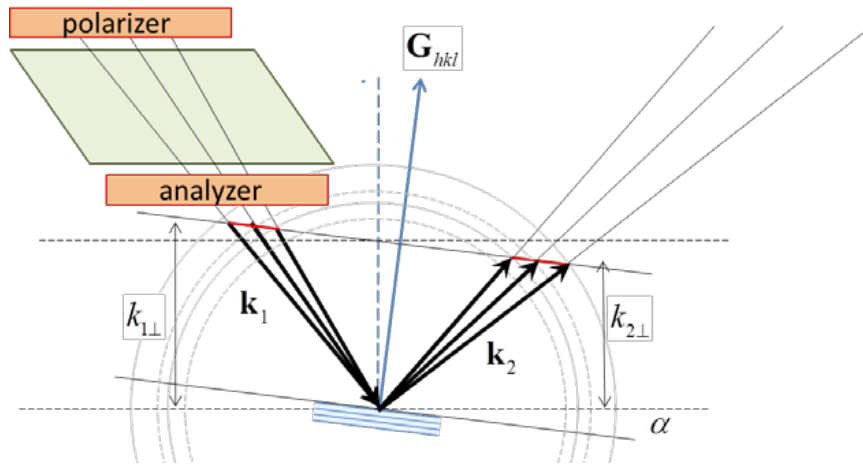
$$k_{i,\perp} = \frac{2\pi}{d}$$

$$\Phi = \frac{\gamma_N m B L d}{2\pi h}$$



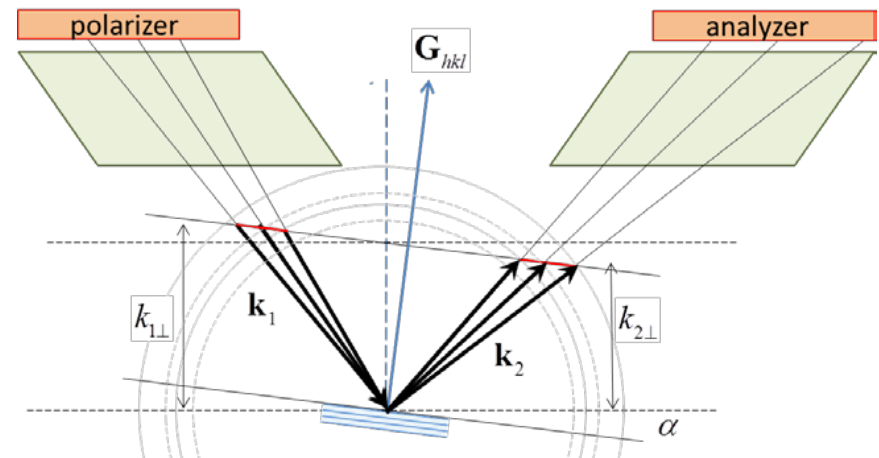
- ❖ The  $d$  spacing distribution is encoded into Larmor phase
- ❖ High resolution for  $d$ -spacing expansion  $\Delta d/d \sim 10^{-6}$ .
- ❖ The tilting angle has to match the crystal plane.

## Single arm Larmor diffraction



Drawback: Lower resolution and not for mosaic sample  
Advantage: OK for magnetic sample

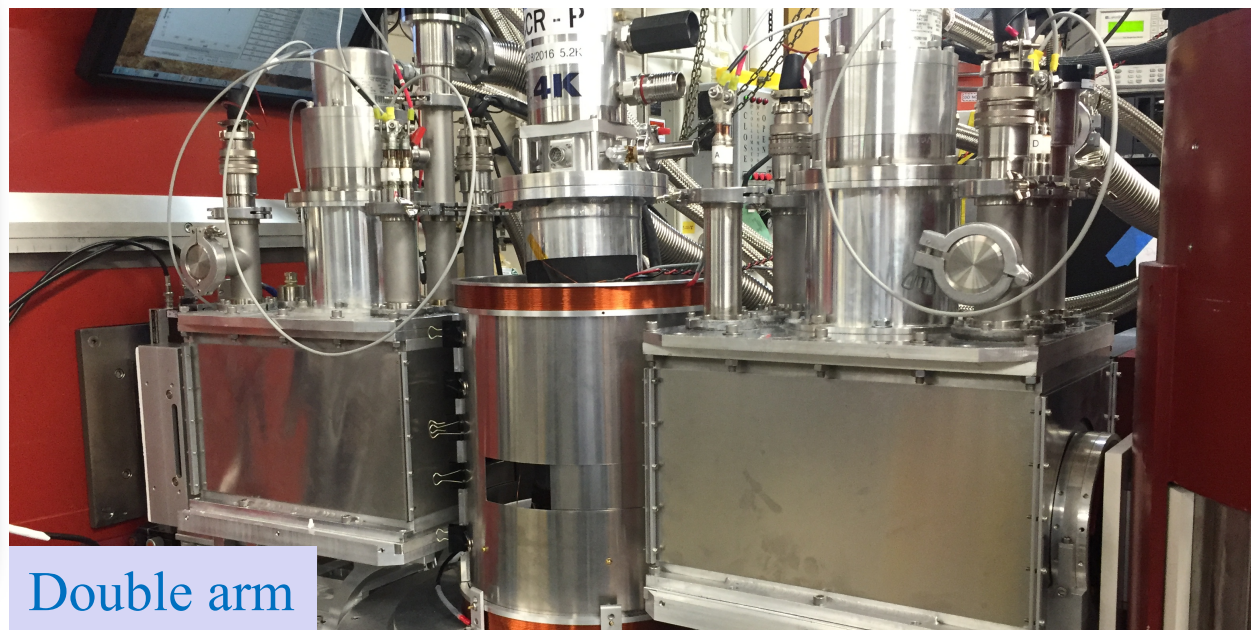
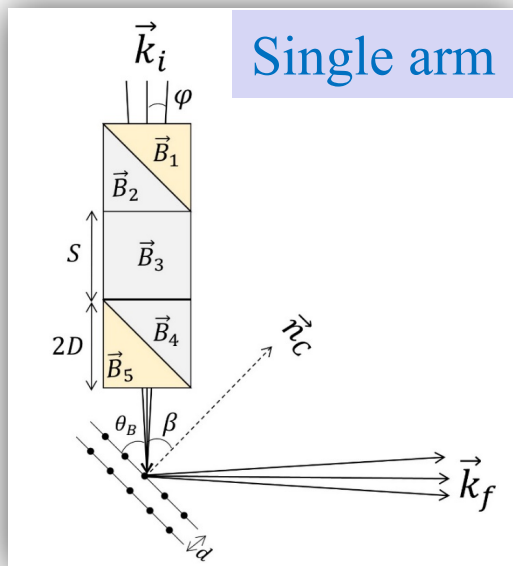
## Double arm Larmor diffraction



Drawback: problematic for magnetic sample  
Advantage: higher resolution, OK with mosaic sample

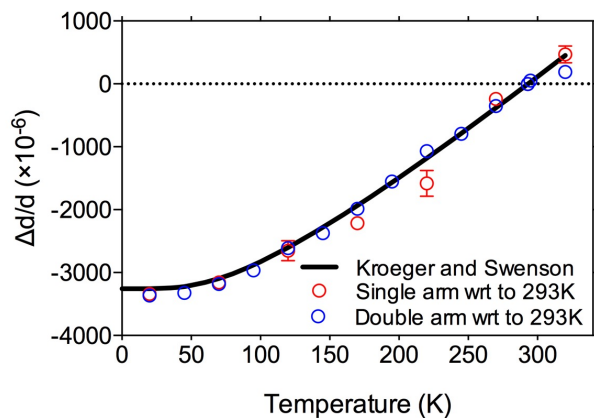
*Courtesy of Dr. Ad van Well*

# Larmor diffraction with Wollaston Prisms: all electromagnetic tuning

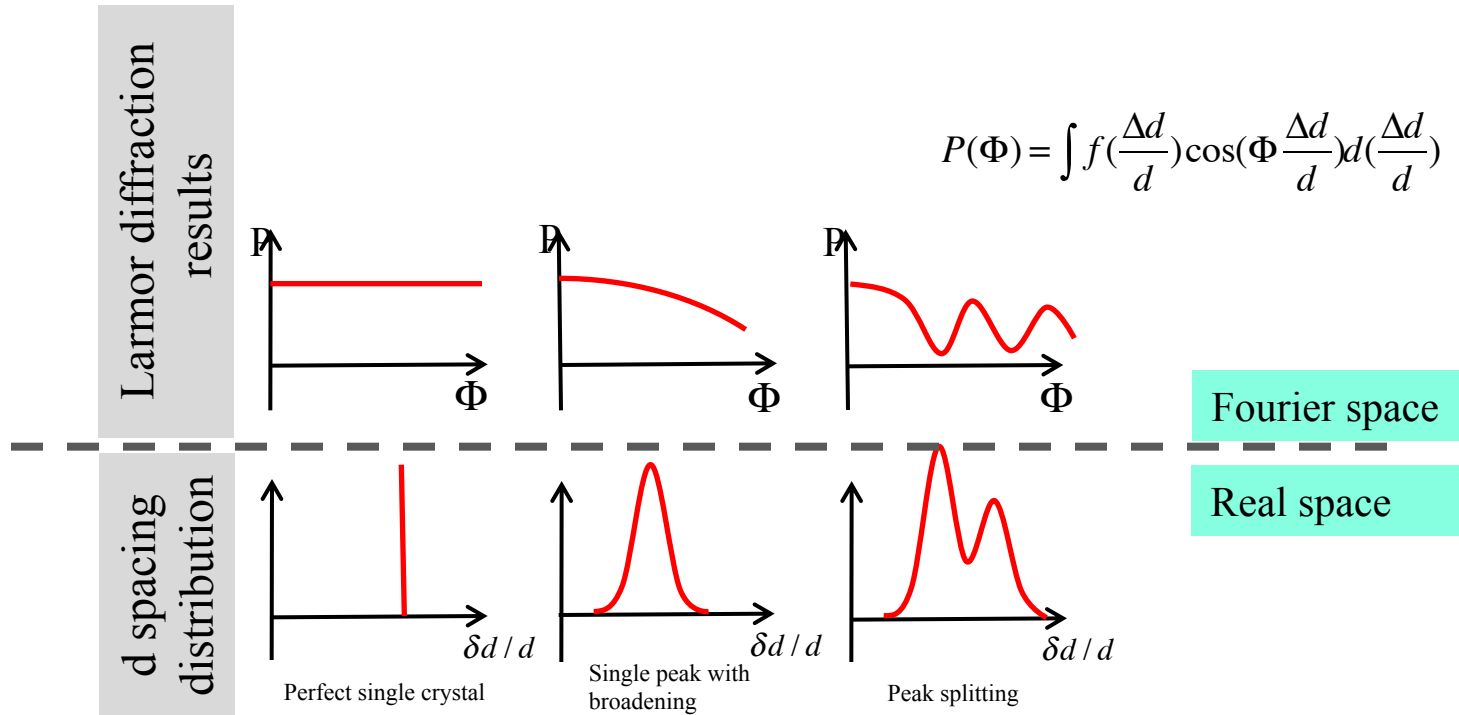


HB-1 @ ORNL  
Feb. 2017

Thermal  
expansion of  
silicon



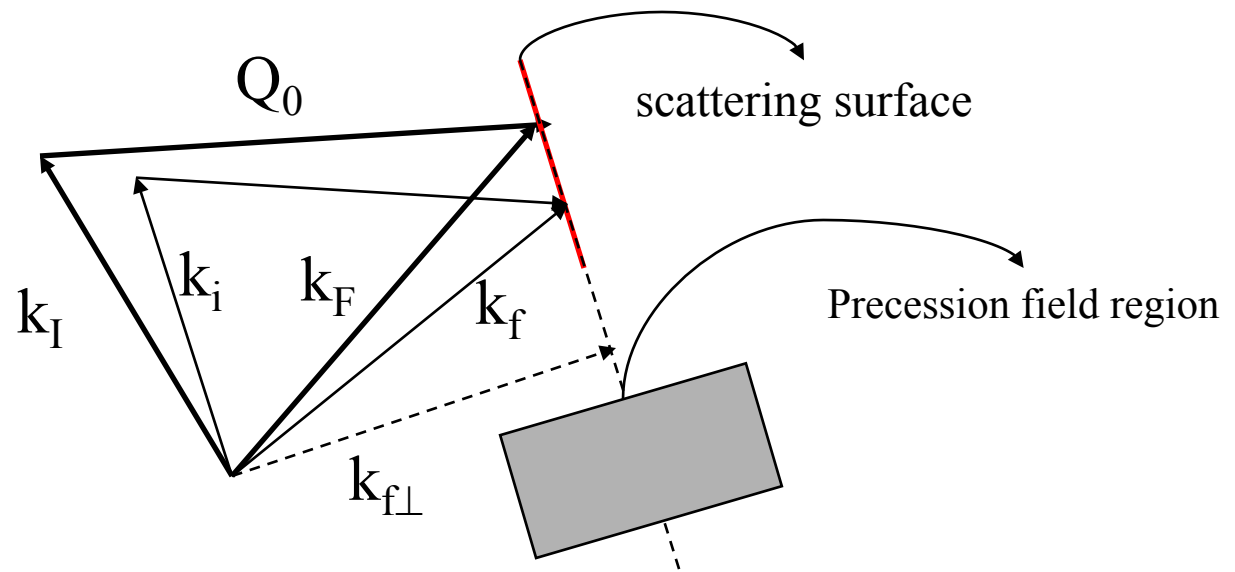
# Rules of thumb for Larmor diffraction



❖ The maximum range of the total Larmor phase denotes the capability of the setup to see small lattice splitting.

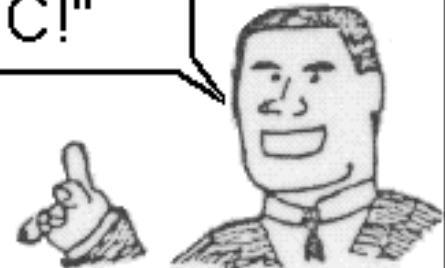
# “Phonon Focusing”

- For a single incident neutron wavevector,  $\mathbf{k}_I$ , neutrons are scattered to  $\mathbf{k}_F$  by a phonon of frequency  $\omega_0$ . The “scattering surface” is the locus in Q space of all phonons of frequency  $\omega_0$ .
- Provided the edges of the NSE precession field region are parallel to the scattering surface, all neutrons with scattering wavevectors on the scattering surface will have equal spin-echo phase



"In just 38 months, you can earn big PROFIT\$  
as a fully trained QUANTUM MECHANIC!"

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in your spare time, without quitting your present job!



Thank You

Questions?