



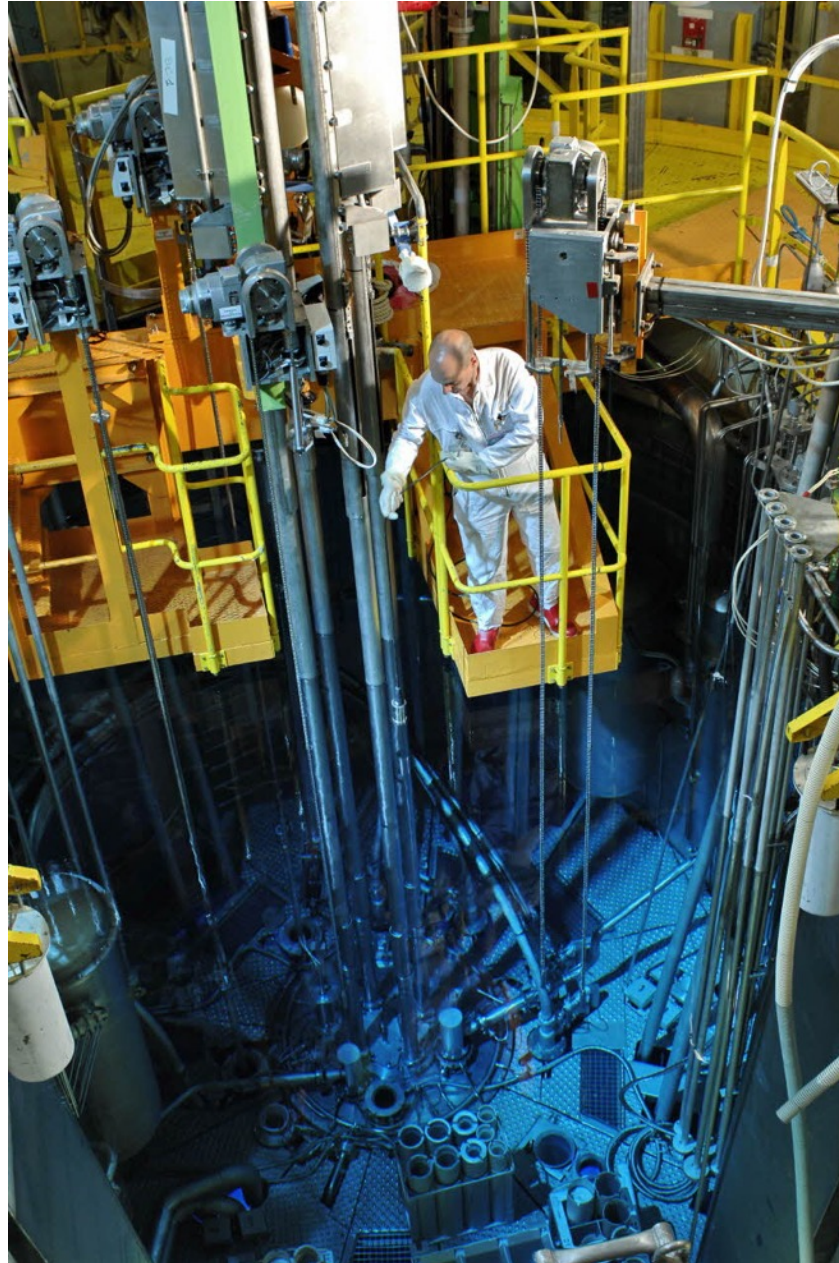
# Neutron Resonance Spin Echo at the LLB.

## Utility of transverse NRSE for quasi elastic scattering?

**S. Longeville,**  
Laboratoire Léon Brillouin (LLB)  
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France



Orphée reactor : 14 MW D<sub>2</sub>O moderated reactor

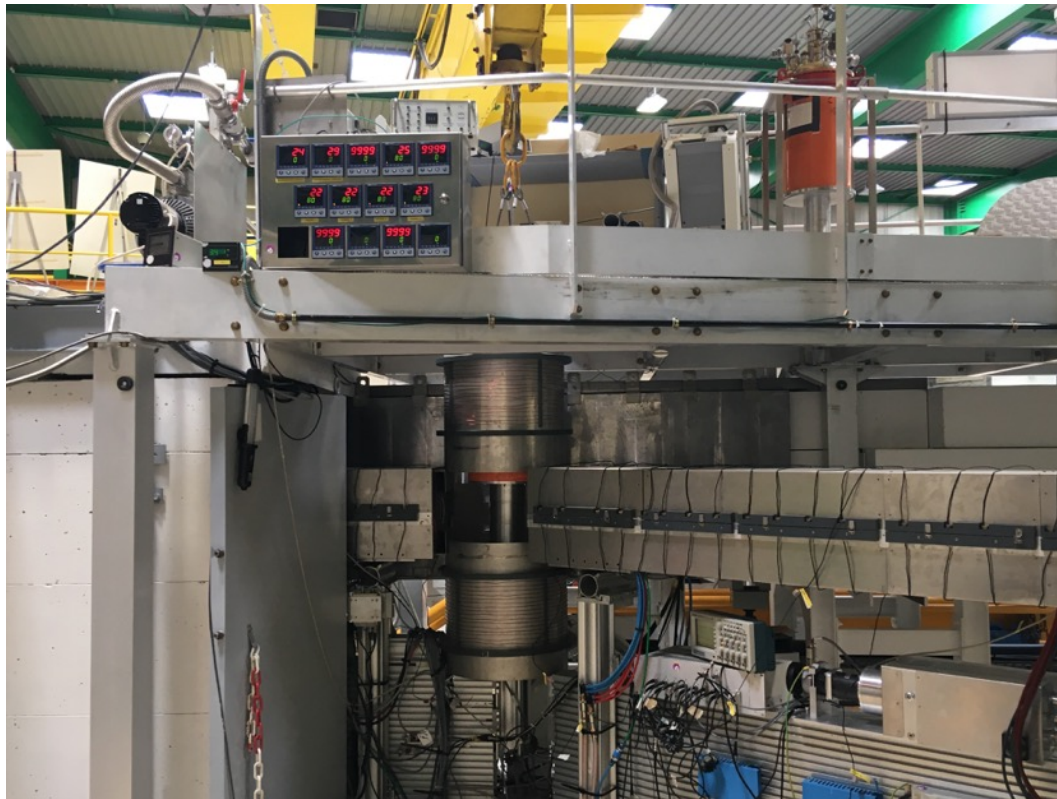


Neutron Spin-Echo Spectroscopy at Saclay  
*Mess : « Mezei » type*

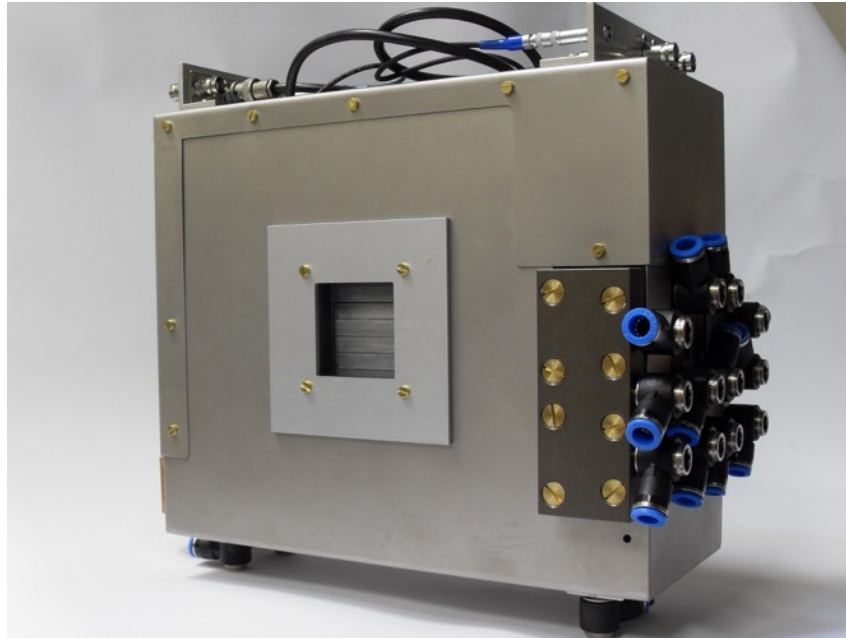


† 2005

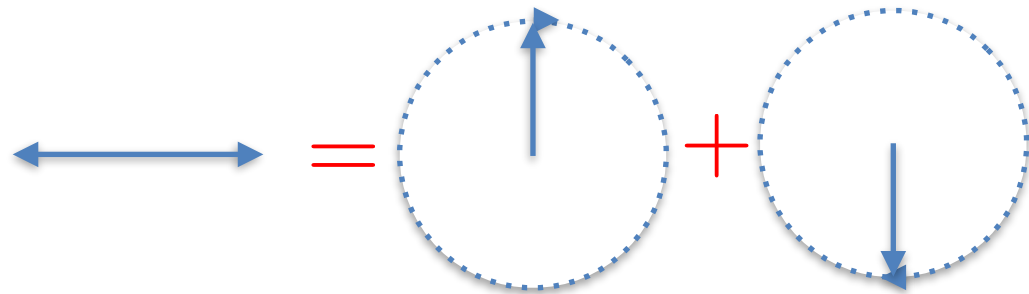
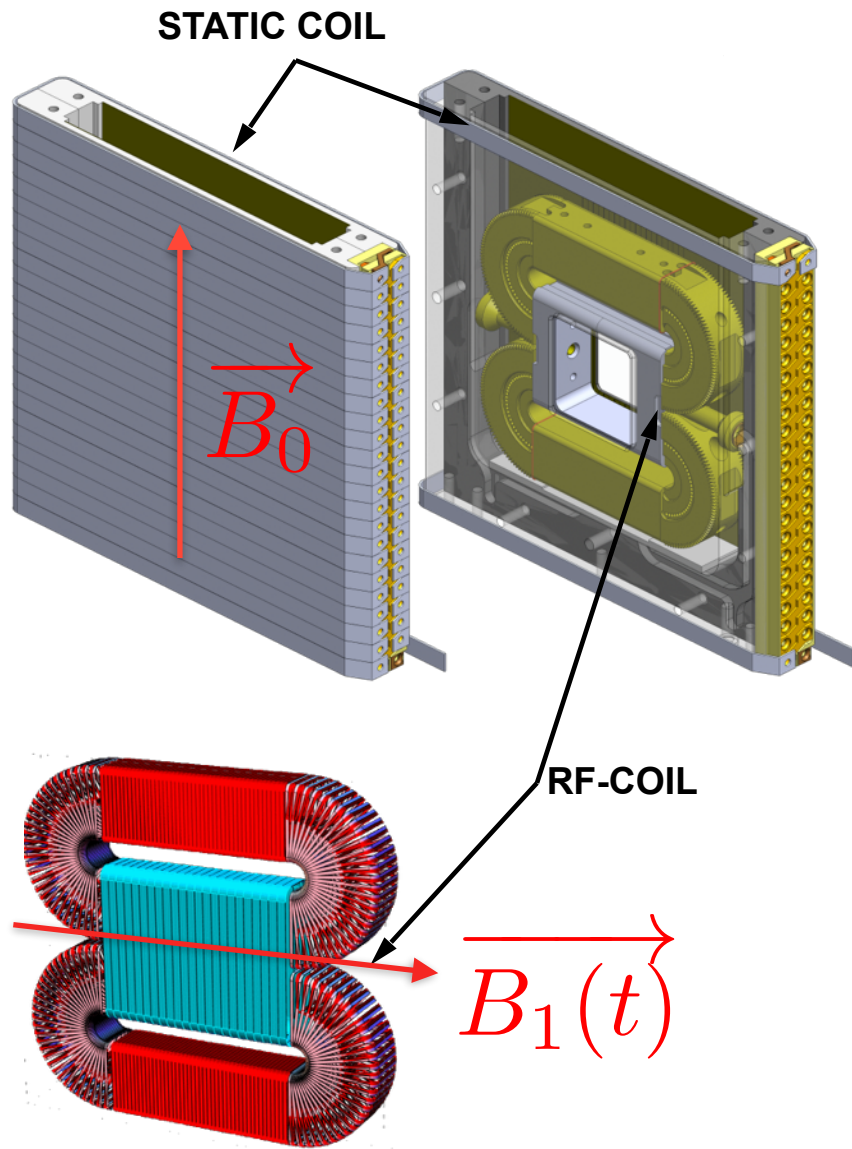
Neutron Spin-Echo Spectroscopy at Saclay  
*Mess : « Mezei » type*  
*G1bis (Muses) Neutron Resonance Spin Echo*



Resonance method for quasi-elastic scattering ?



*courtesy : S. Klimko*

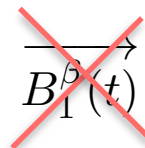


$$\vec{B}_1(t) = \vec{B}_1^\alpha(t) + \vec{B}_1^\beta(t)$$

$\vec{B}_0$  spin precession counter clockwise  
 $\gamma_n < 0$

if  $\vec{B}_1^\alpha(t)$  is in the direction of rotation then  
 $\vec{B}_1^\beta(t)$  is in the opposite direction

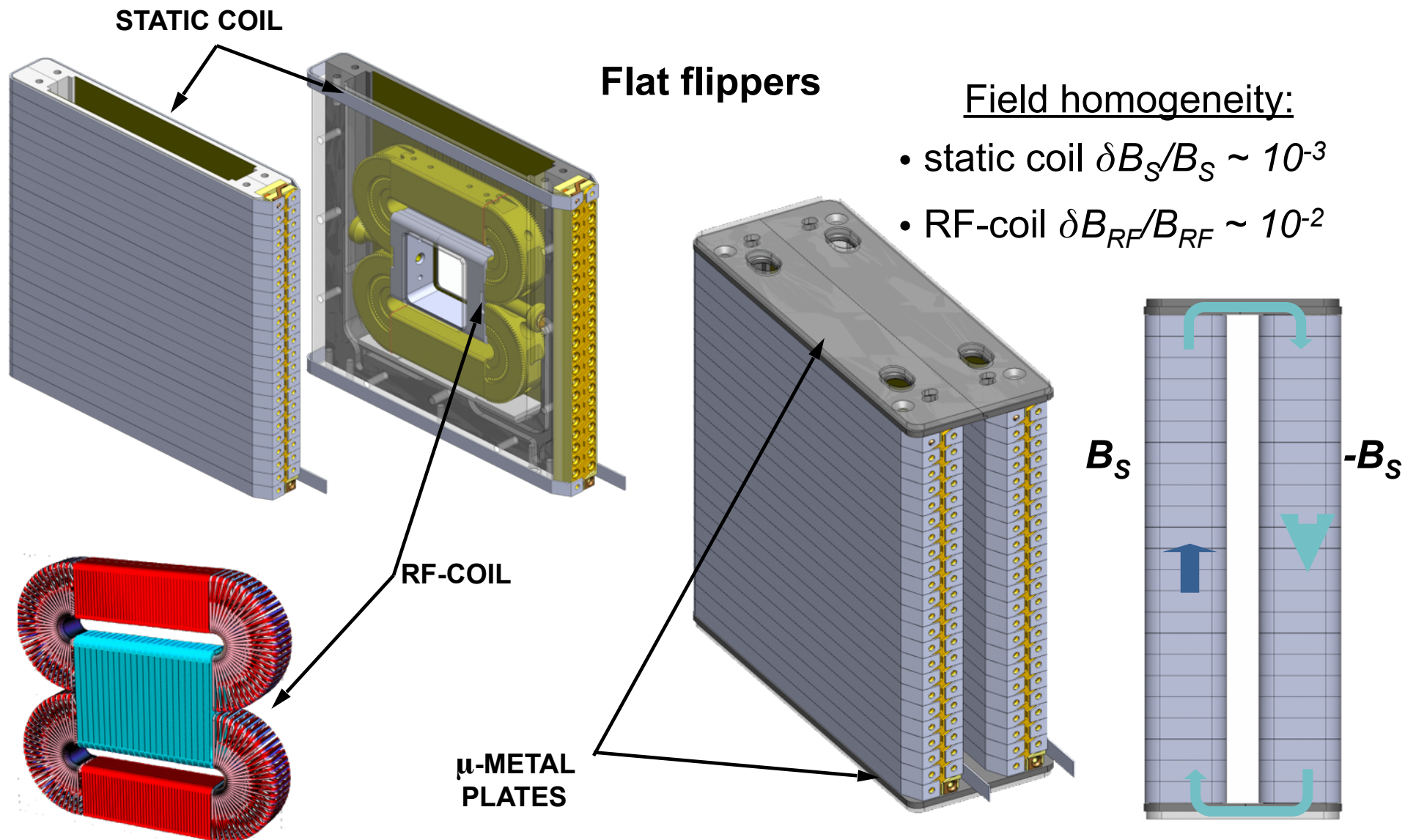
$$\|\vec{B}_0\| \gg \|\vec{B}_1(t)\|$$



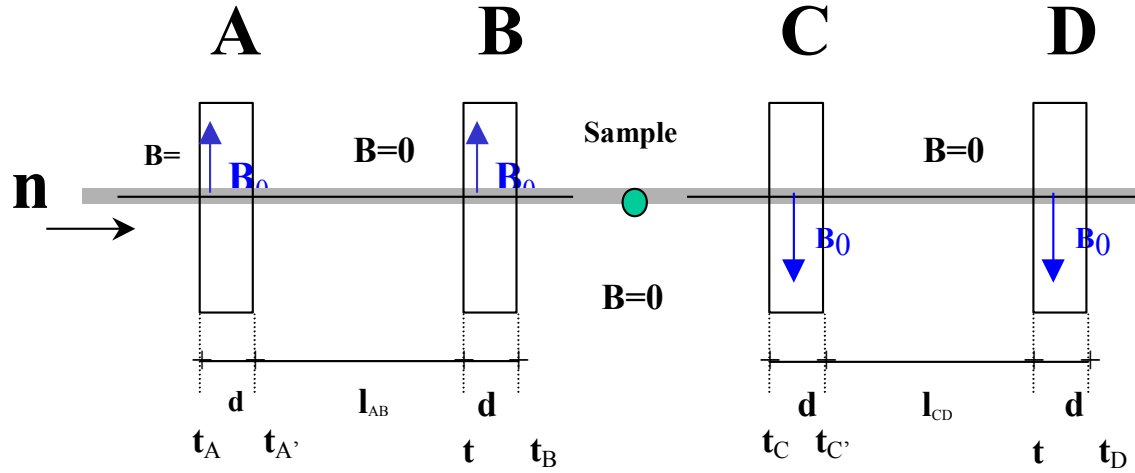
Not short times

Bloch-Siegert shift

## New Flat Resonance Flippers (1st Arm)



## NRSE spectrometer



1

**Table 1.** Spin orientation

	Time $t$	Phase field $B_r$	neutron Spin phase $S$
A	$t_A$	$\omega t_A$	0
A'	$t_{A'} = t_A + \frac{d}{v}$	$\omega t_{A'}$	$2\omega t_A + \omega \frac{d}{v}$
B	$t_B = t_A + \frac{l_{AB}+d}{v}$	$\omega t_B$	$2\omega t_A + \omega \frac{d}{v}$
B'	$t_{B'} = t_A + \frac{l_{AB}+2d}{v}$	$\omega t_{B'}$	$2\omega \frac{l_{AB}+d}{v}$
C	$t_C$	$-\omega t_C$	$2\omega \frac{l_{AB}+d}{v}$
C'	$t_{C'} = t_C + \frac{d}{v}$	$-\omega t_{C'}$	$-\omega \frac{d}{v'} - 2\omega t_C - 2\omega \frac{l_{AB}+d}{v}$
D	$t_D = t_C + \frac{l_{CD}+d}{v'}$	$-\omega t_D$	$-\omega \frac{d}{v'} - 2\omega t_C - 2\omega \frac{l_{AB}+d}{v}$
D'	$t_{D'} = t_C + \frac{l_{CD}+2d}{v'}$	$-\omega t_{D'}$	$2\omega \left( \frac{l_{AB}+d}{v} - \frac{l_{CD}+d}{v'} \right)$

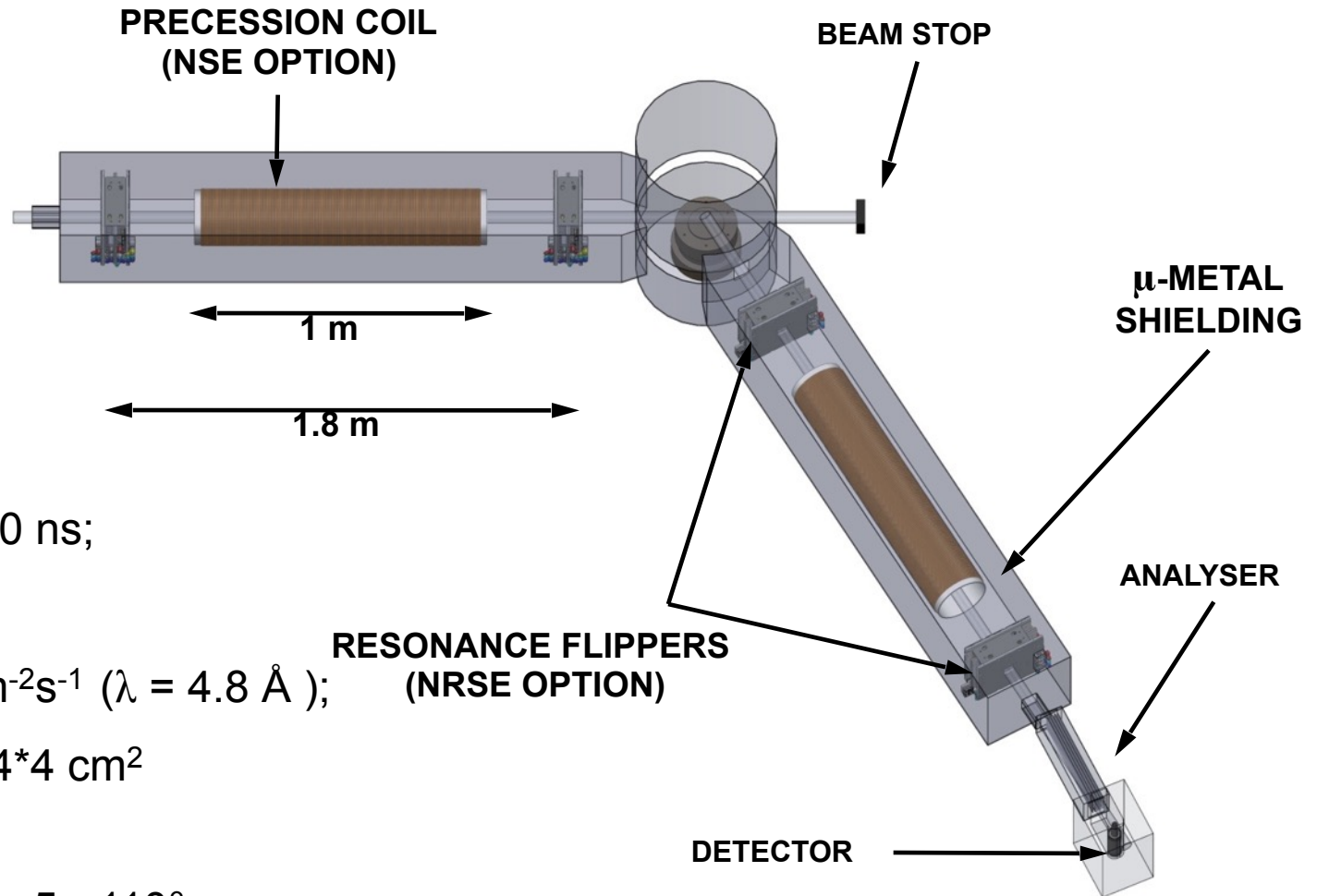


$$\text{Echo condition NRSE} \quad \frac{L_{AB}}{v} = \frac{L_{CD}}{v'} \qquad \text{NSE} \quad \frac{\int B_0 dl}{v} = \frac{\int B_1 dl}{v'}$$

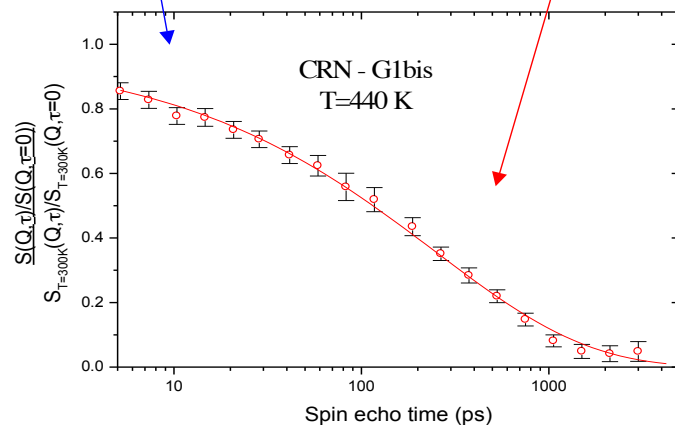
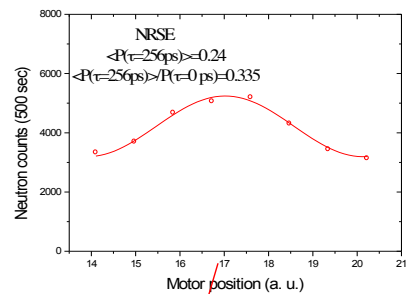
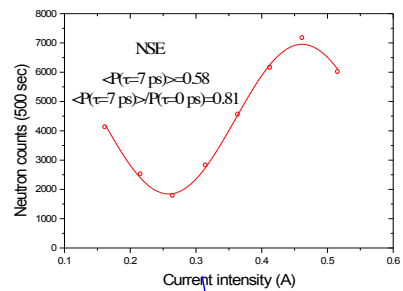
$$\langle P \rangle \cong \int_0^{\infty} I(\lambda) \int_{-\infty}^{\infty} S(Q, \omega) \cos(\omega \tau_{NRSE}) d\omega d\lambda$$

$$\tau_{NRSE} = 2\omega_f \frac{l+d}{2\pi} \frac{m^2}{h^2} \lambda^3$$

## MUSES (G1<sub>bis</sub>) - mixed NSE / NRSE



- $\tau_{\text{NSE}} : 0.2 \text{ ps} - 20 \text{ ns};$
- $\Delta\lambda/\lambda \approx 15\%;$
- $\Phi_s = 2 \cdot 10^7 \text{ n cm}^{-2}\text{s}^{-1}$  ( $\lambda = 4.8 \text{ \AA}$ );
- beam section:  $4 \times 4 \text{ cm}^2$
- $\lambda: 3.5 - 14 \text{ \AA}$
- Scattering angle:  $5 - 110^\circ$
- Q range:  $0.05 - 4 \text{ \AA}^{-1}$



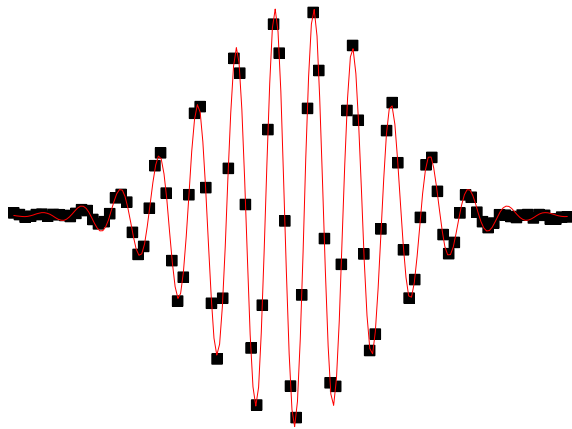
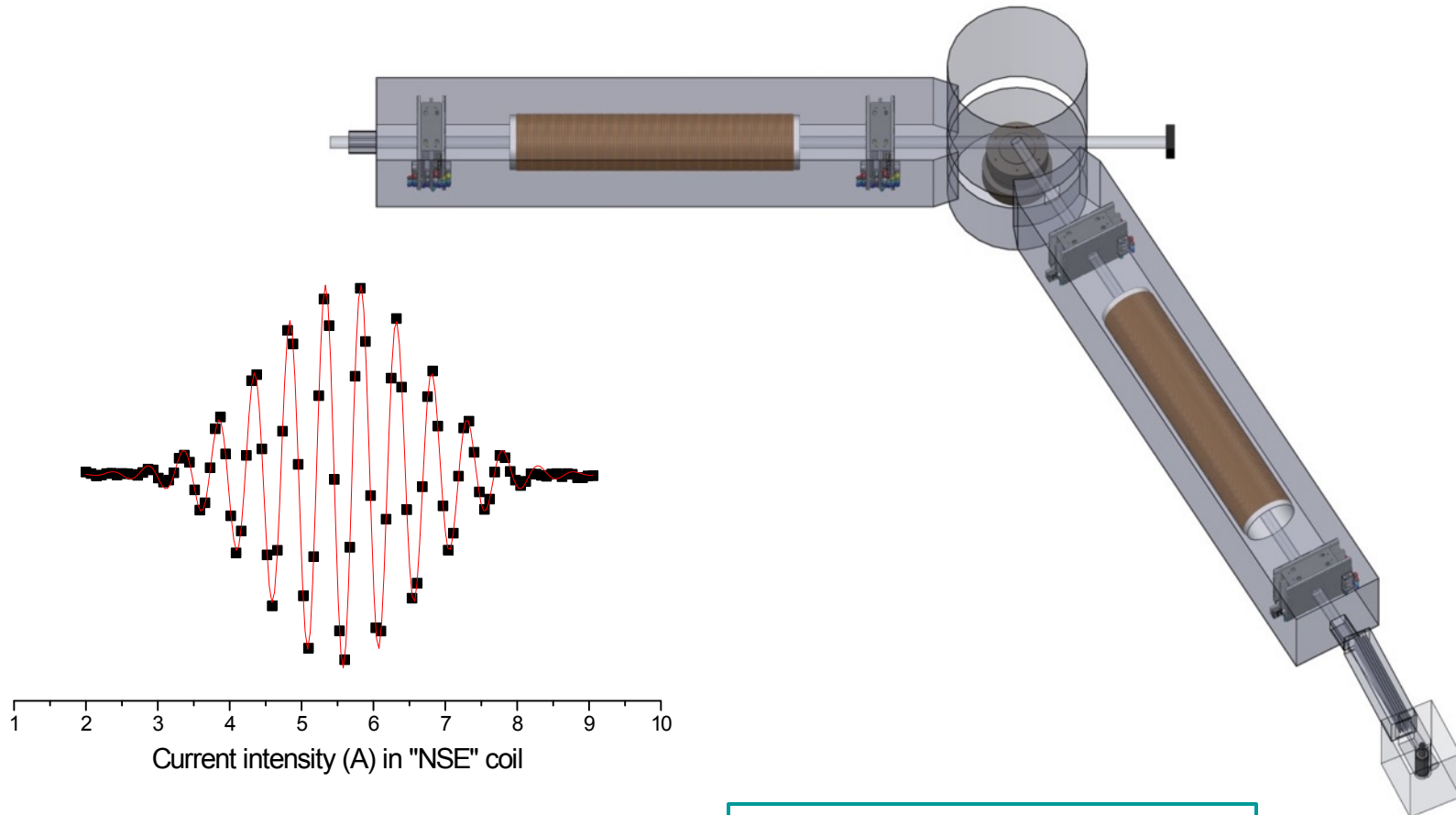
Current intensity (A)

Motor position (a. u.)

Spin echo time (ps)

# Neutron spin echo for quasi-elastic scattering

Small times  $\rightarrow$  NSE option !

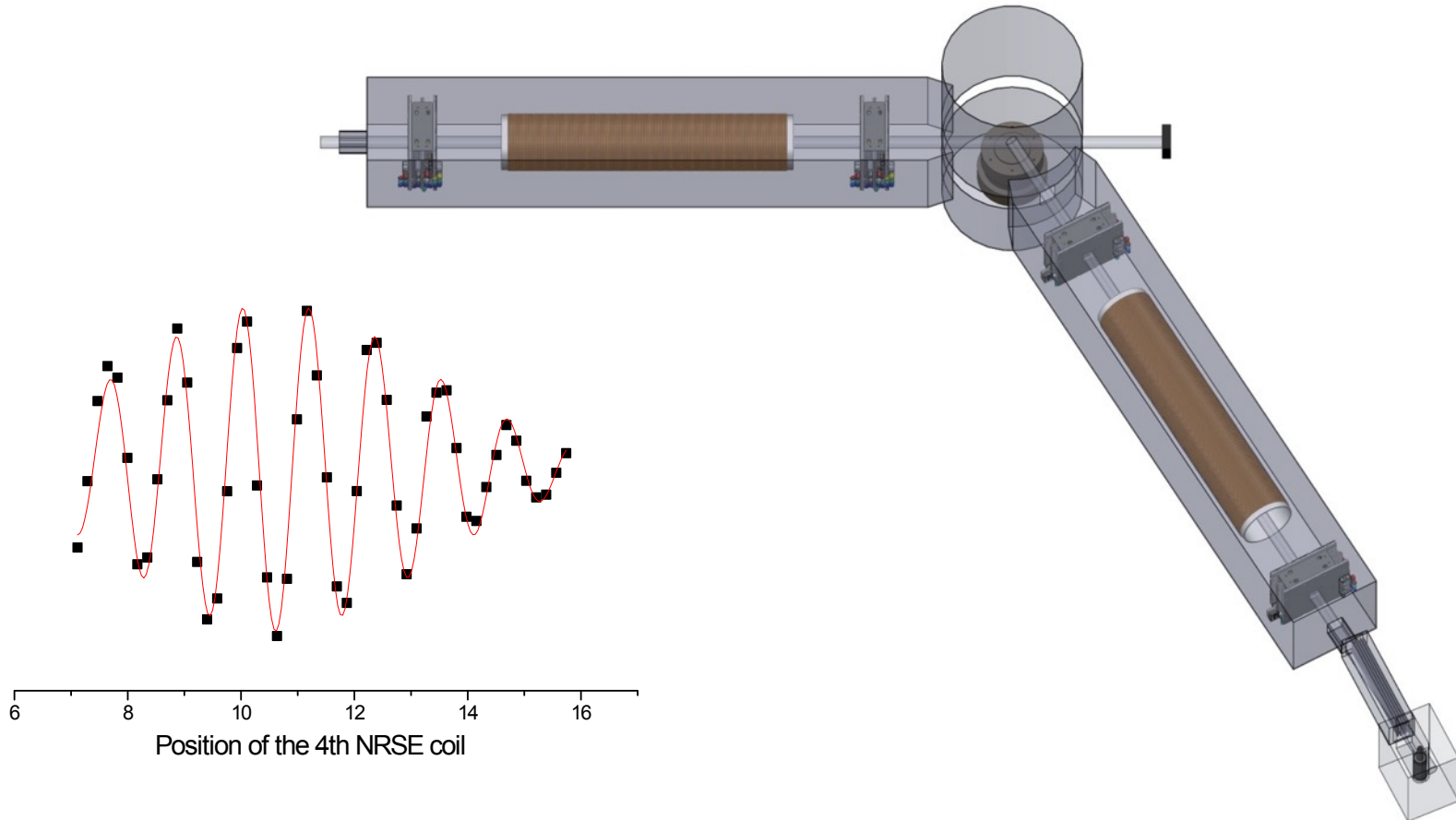


1 2 3 4 5 6 7 8 9 10  
Current intensity (A) in "NSE" coil

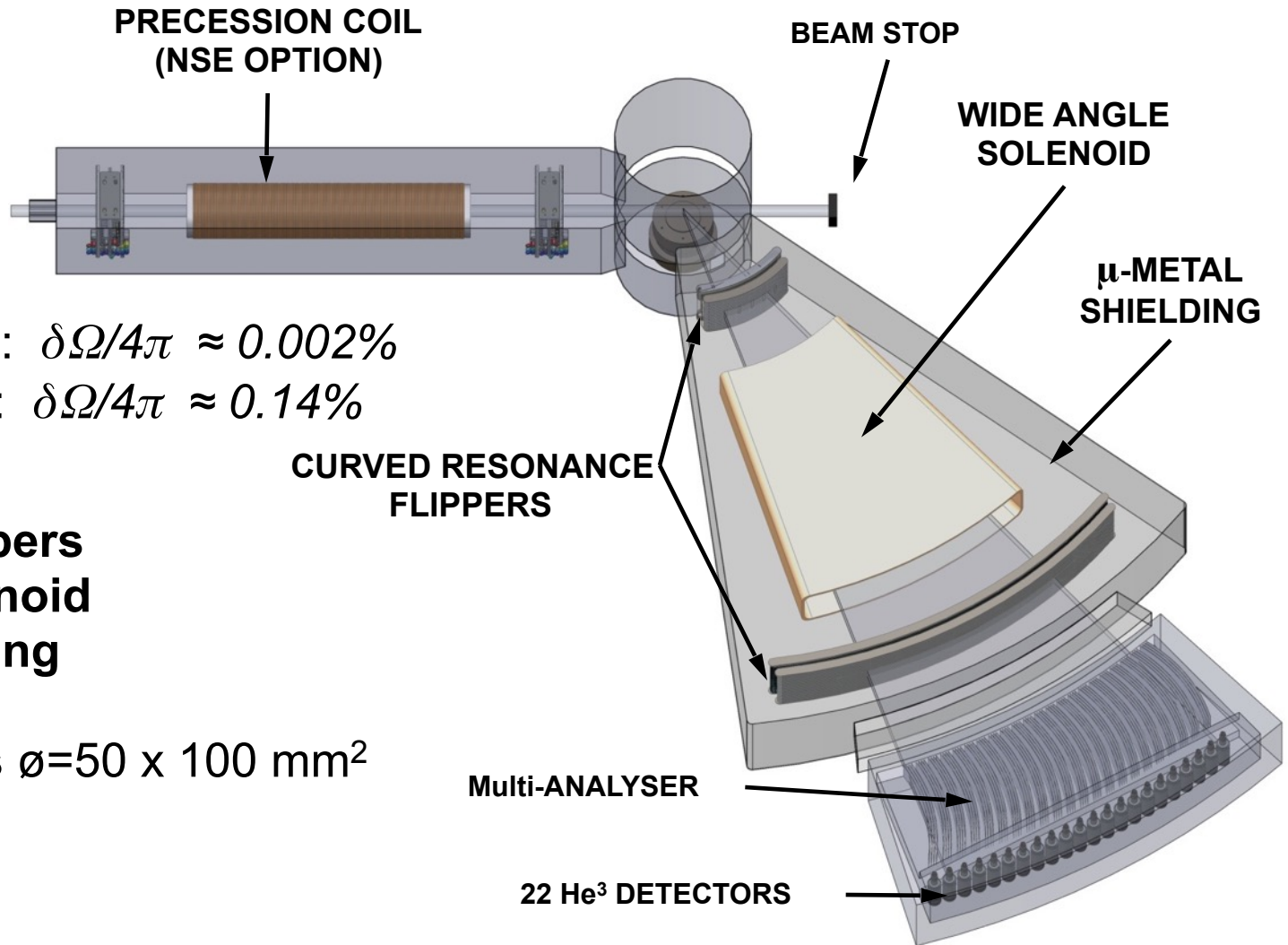
$$\int |B_1| dl = \int |B_2| dl$$

# Neutron spin echo for quasi-elastic scattering

long times  $\rightarrow$  NRSE option !



# Multi - MUSES



Single detector :  $\delta\Omega/4\pi \approx 0.002\%$

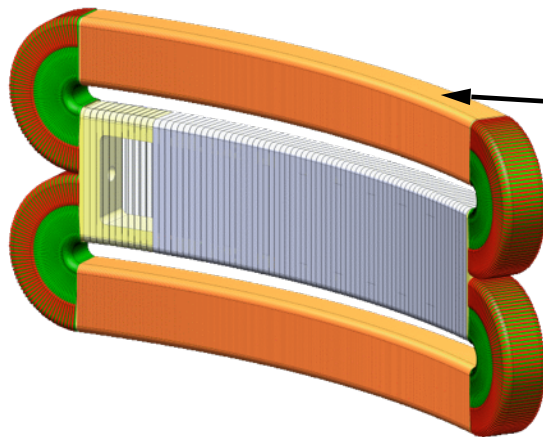
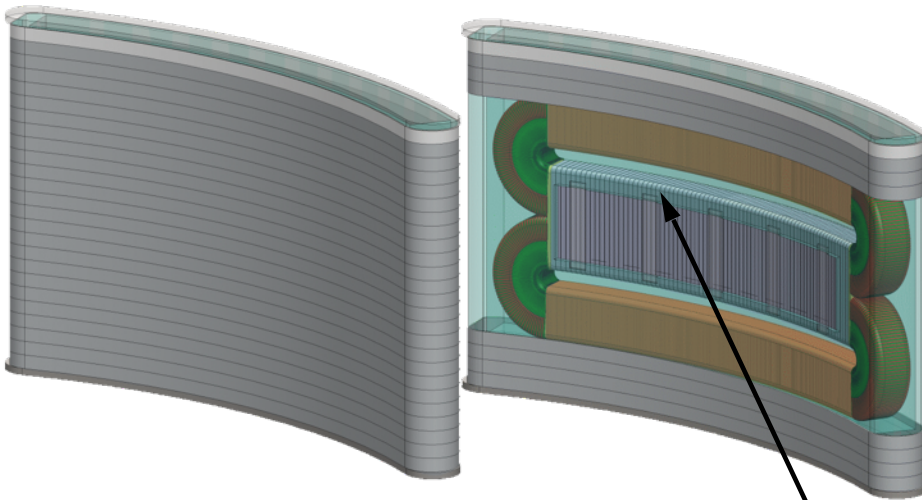
'Multi'- detector :  $\delta\Omega/4\pi \approx 0.14\%$

- Resonance flippers
- Wide angle solenoid
- Magnetic shielding
- Multi – Analyser
- 22 He<sup>3</sup> detectors  $\varnothing=50 \times 100 \text{ mm}^2$

# Curved Resonance Flippers

## Curved flippers

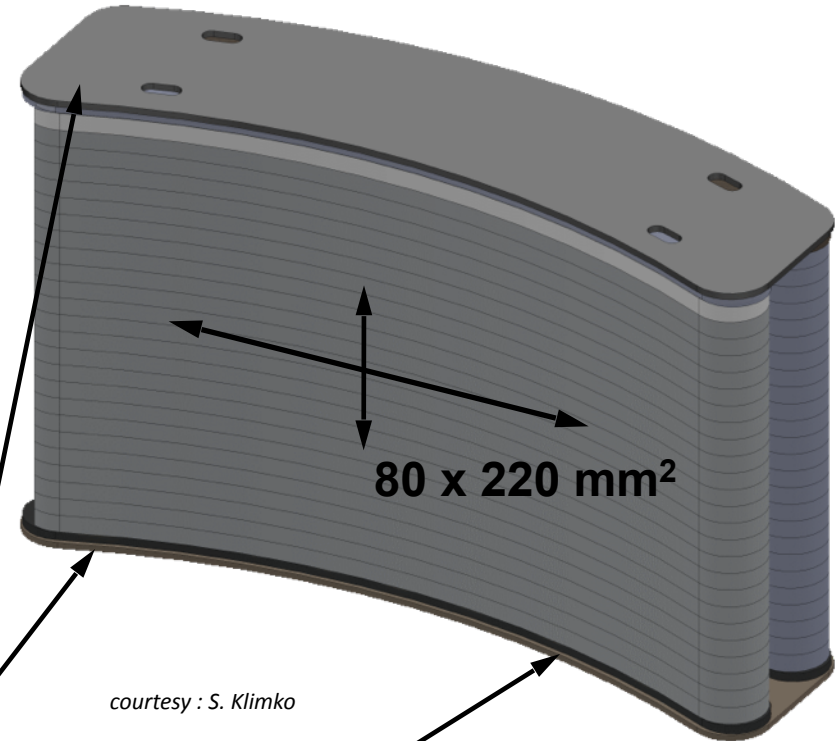
STATIC COILS



RF-COIL

$\mu$ -METAL  
PLATES

BOOTSTRAP-PAIR

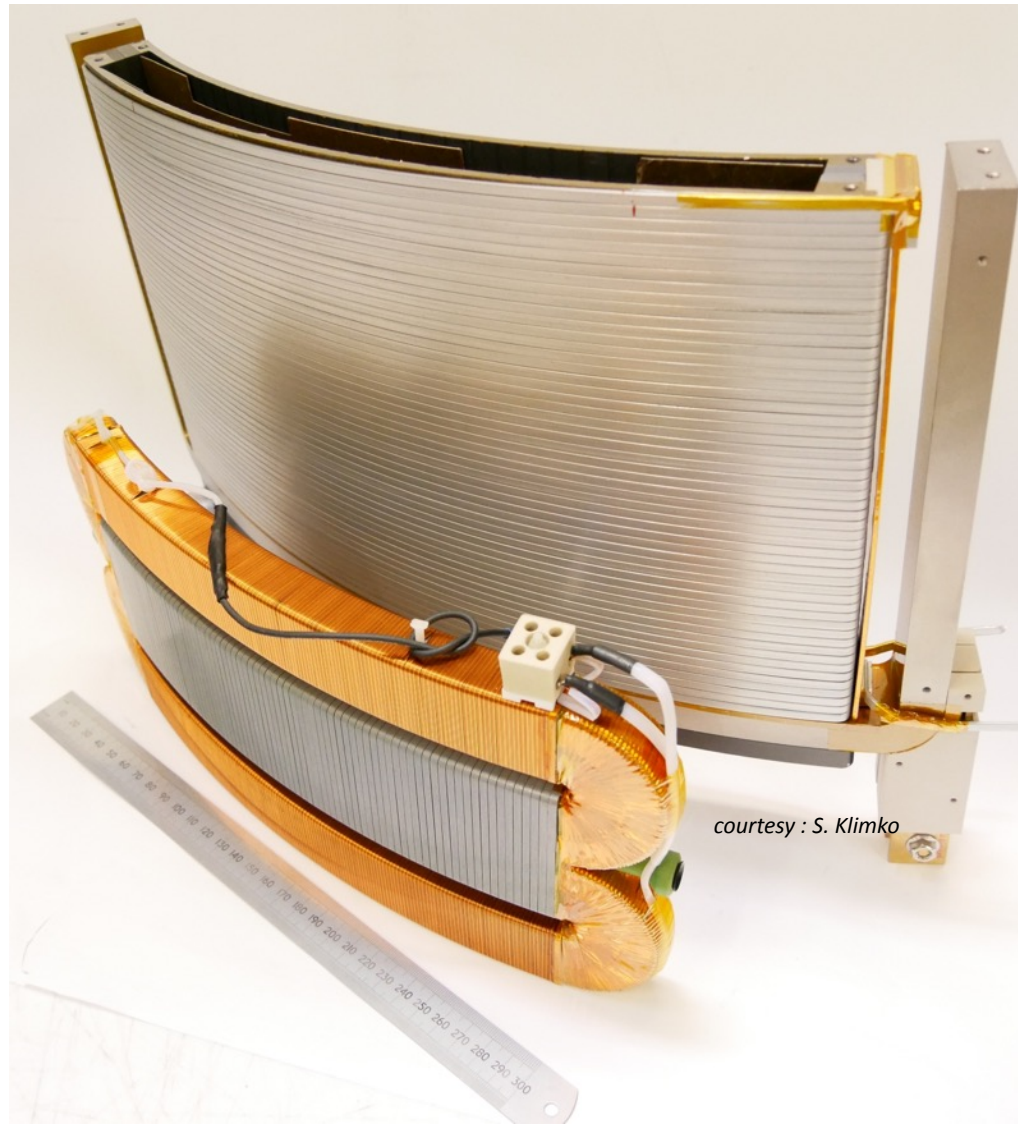


$80 \times 220 \text{ mm}^2$

*courtesy : S. Klimko*

$R_{\text{CUR}} = 400 \text{ mm}$

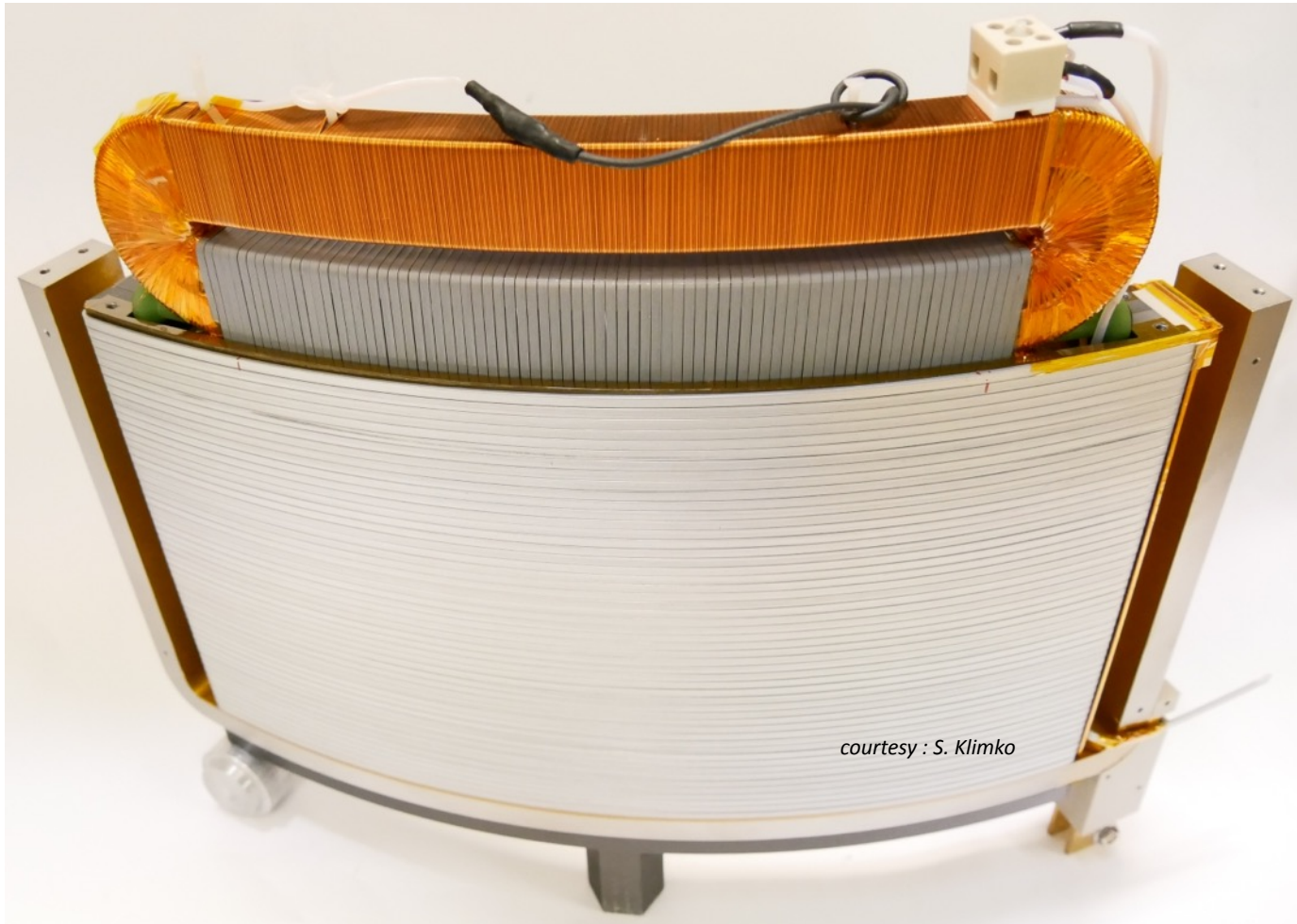
## Curved Resonance Flippers



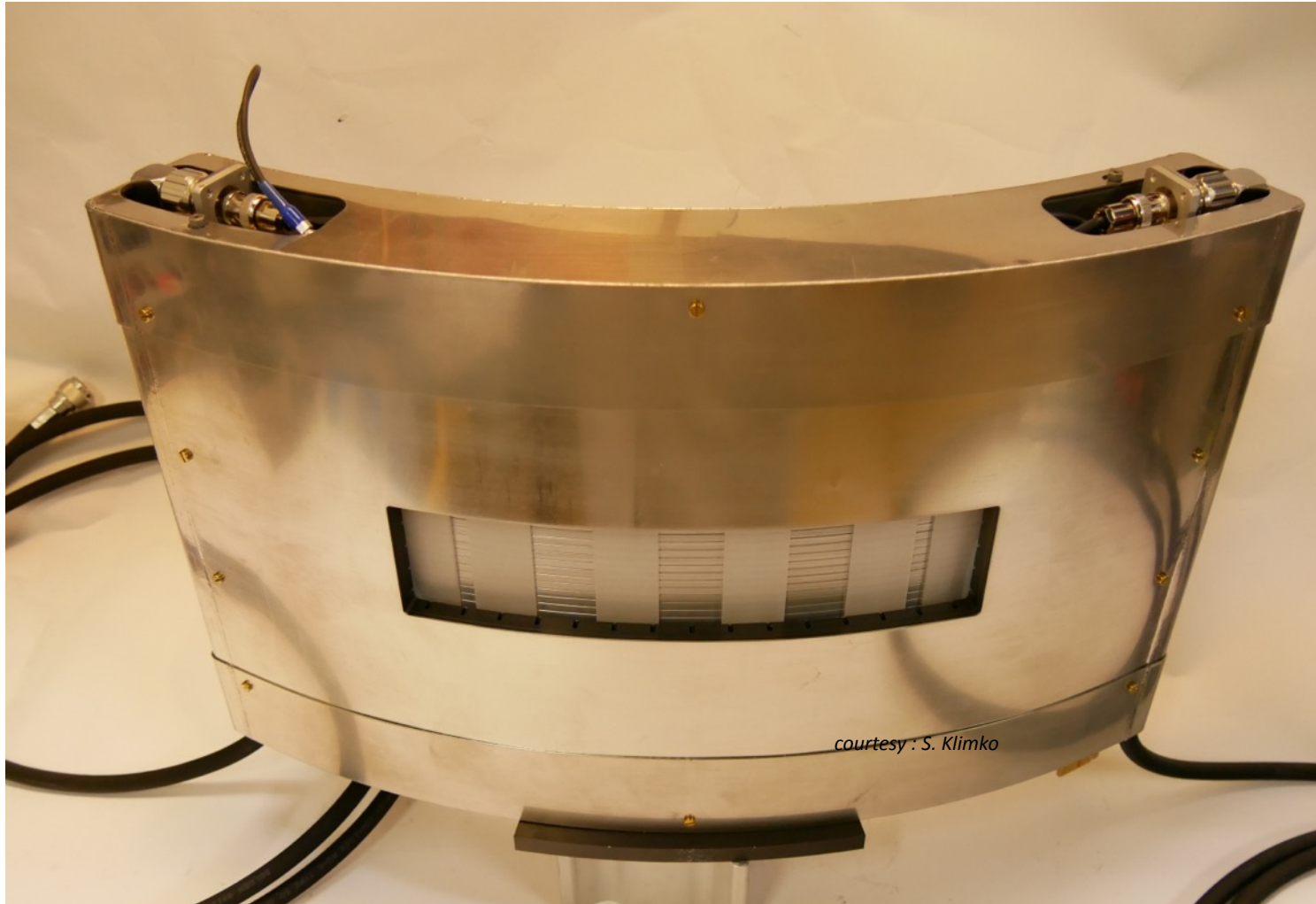
*courtesy : S. Klimko*



## Curved Resonance Flippers



## Curved Resonance Flippers

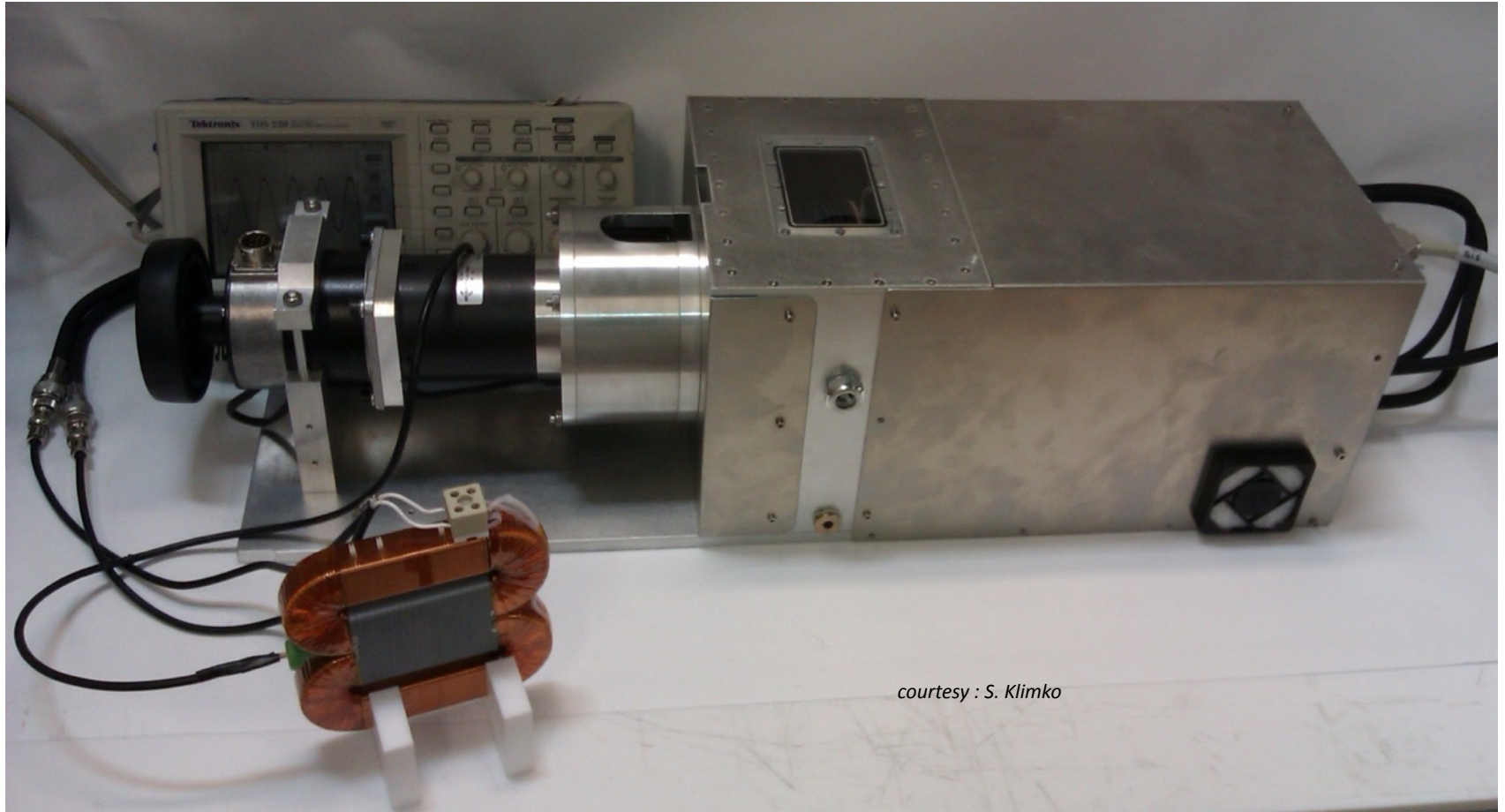


## Curved Resonance Flippers

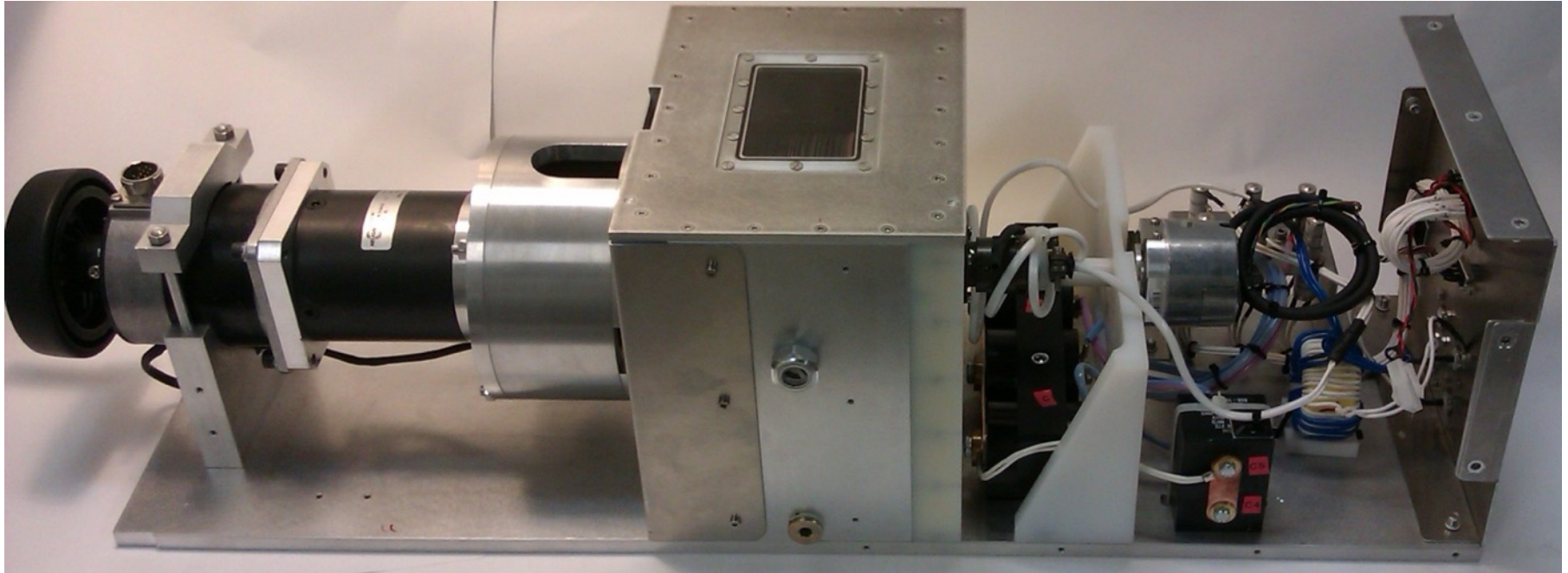


*courtesy : S. Klimko*

# HF circuit : impedance adaptation

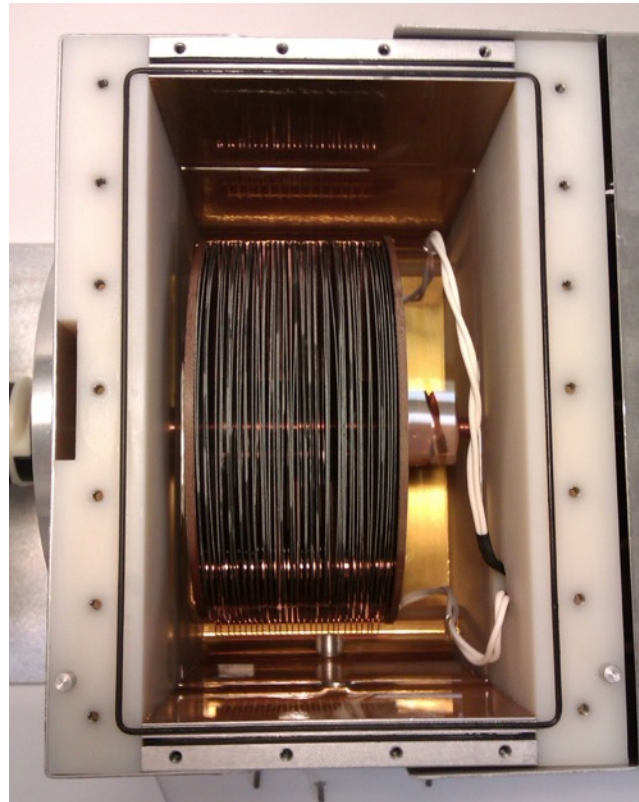


# HF circuit : impedance adaptation



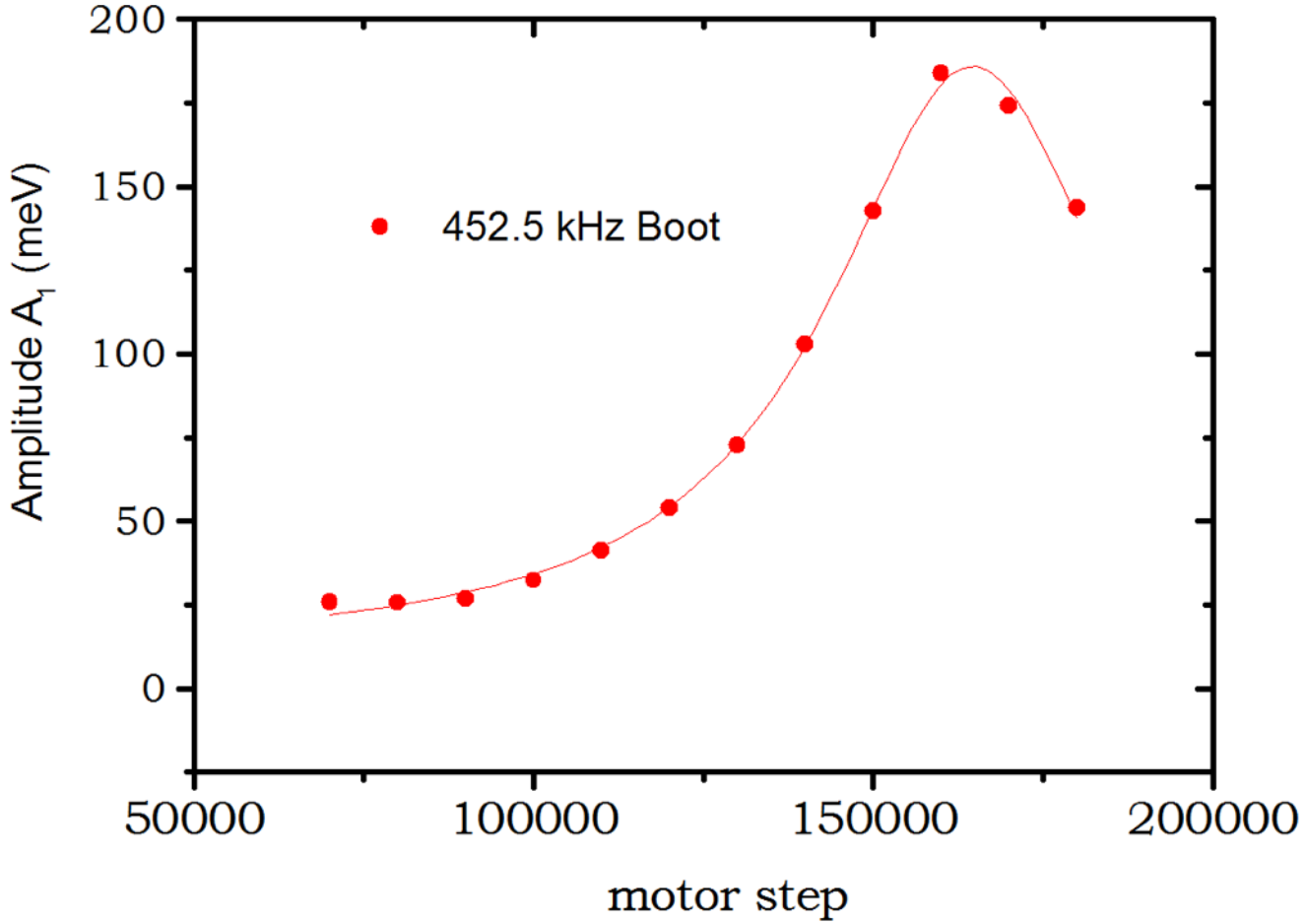
*courtesy : S. Klimko*

HF circuit : impedance adaptation



*courtesy : S. Klimko*

HF circuit : impedance adaptation



## NRSE versus NSE@ Saclay

- Mixed spectrometer
  - Not the short times (Bloch-Siegert shift)
  - Not the long times (no correction coils)
- NSE inside the mu-metal shielding : short times (no depolarization of the beam)
- Small to medium times (0.5 ps to 20 ns)
- Compact spectrometer : High flux ( $2 \cdot 10^7$  n.cm<sup>-2</sup>.s<sup>-1</sup> polarized neutrons @ sample position 4\*4 cm<sup>2</sup>)
- Zero field in the sample position :
  - resolution is very little Q dependent
  - not sensible to flight path distribution in the sample
- Multi detector is limited
- Stability of the RF current requires appropriate design of the circuits
- high Q measurements
- Beam is passing through 3\*2 mm aluminium per coil \*8 = 48 mm without sample environment



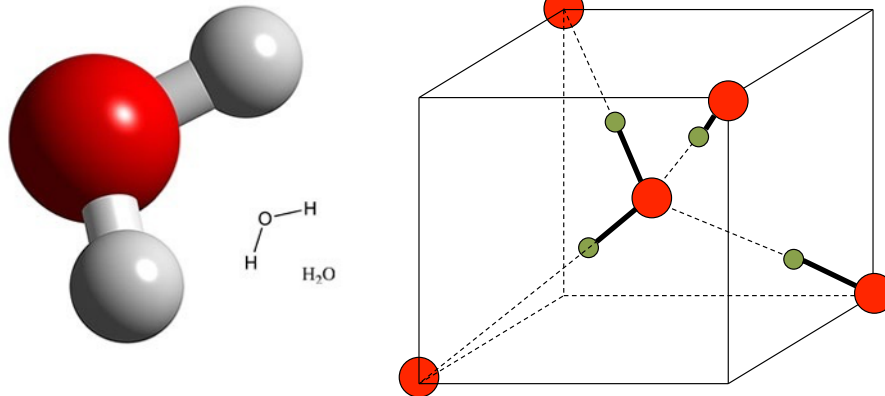


# 1- The dynamics of supercooled water - hydrogen bond lifetime

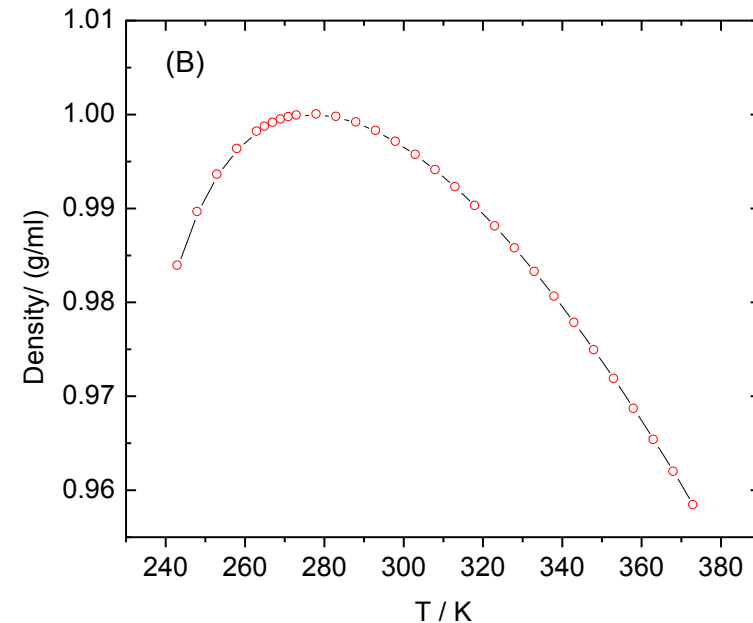
*J. Teixeira, A. Luzar and S. Longeville, J. Phys.: Condens. Matter* **18** (2006) S2353–S2362

**Water:** a common liquid with « anomalous properties »

Why is water so complex ?



- Large number of hydrogen bonds
- Local tetrahedral order
- Short life time of H bonds



*G. Johari and J. Teixeira,  
J. Phys. Chem. B* **119**, 14210 (2015)

*Density, compressibility, heat capacity  
viscosity, isotopic effects ...*



# 1- The dynamics of supercooled water - hydrogen bond lifetime

*J. Teixeira, A. Luzar and S. Longeville, J. Phys.: Condens. Matter* **18** (2006) S2353–S2362

PHYSICAL REVIEW A

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## Experimental determination of the nature of diffusive motions of water molecules at low temperatures

J. Teixeira and M.-C. Bellissent-Funel

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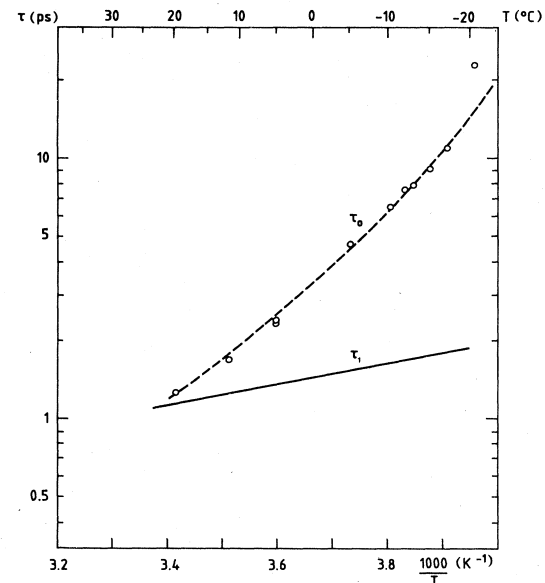
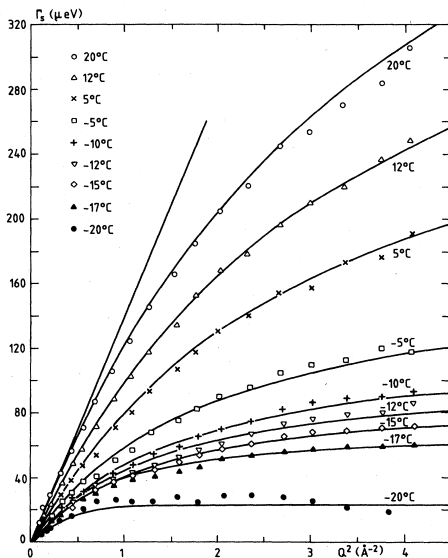
S. H. Chen

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(Received 27 April 1984)



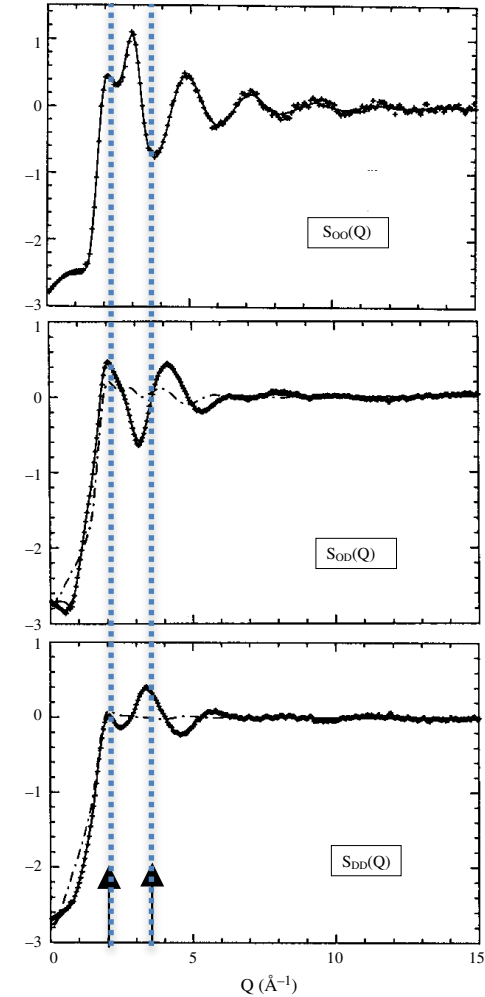
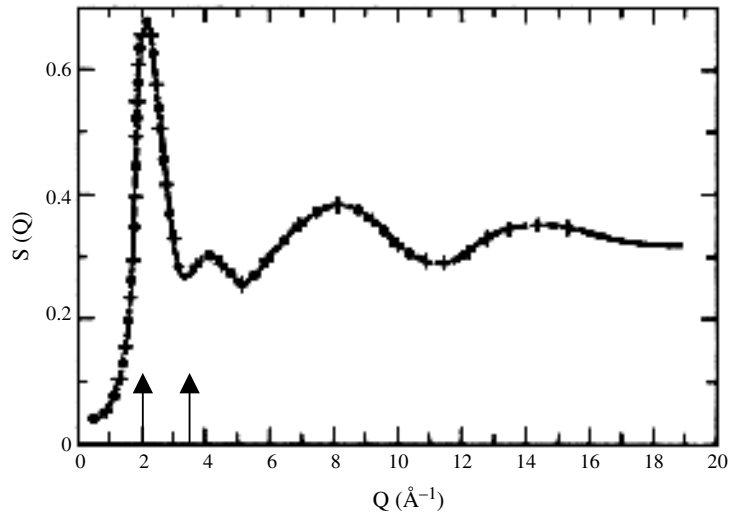


# 1- The dynamics of supercooled water - hydrogen bond lifetime

*J. Teixeira, A. Luzar and S. Longeville, J. Phys.: Condens. Matter* **18** (2006) S2353–S2362

Can we have a more precise measurement of  $\tau_{HB}$

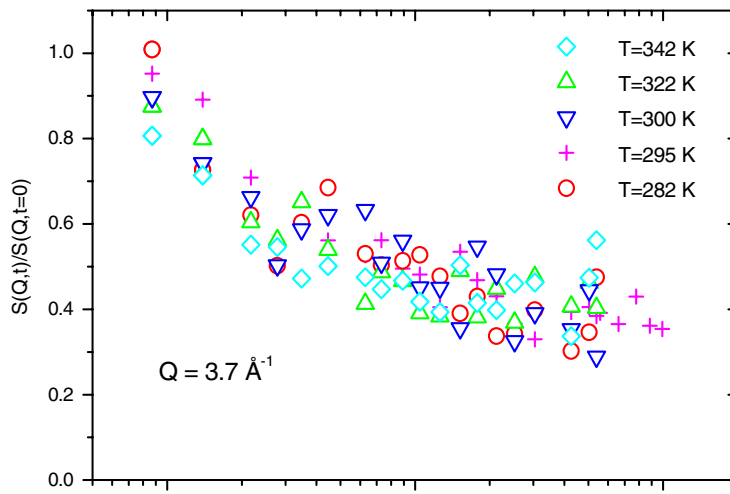
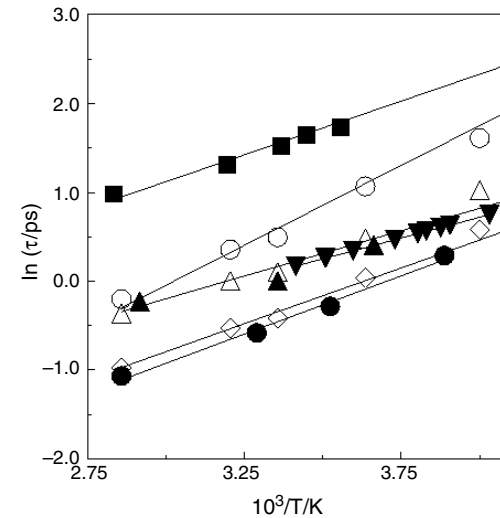
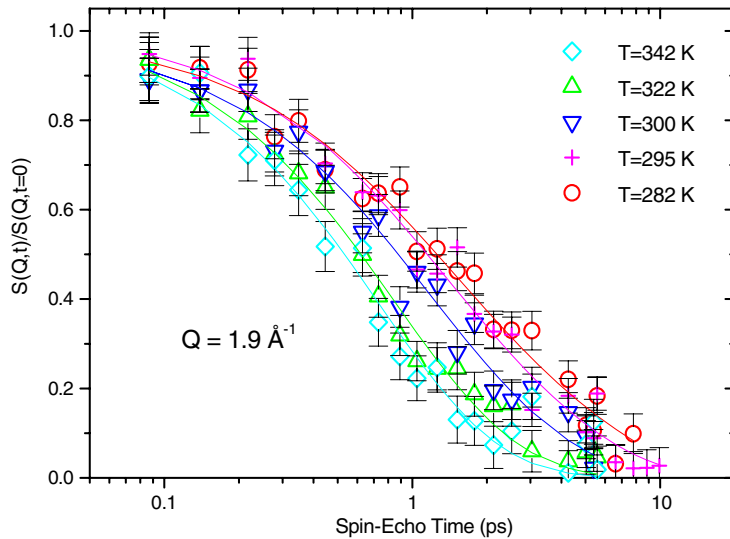
D-D





# 1- The dynamics of supercooled water - hydrogen bond lifetime

*J. Teixeira, A. Luzar and S. Longeville, J. Phys.: Condens. Matter* **18** (2006) S2353–S2362



**Figure 7.** Arrhenius temperature dependence of the hydrogen bond dynamics determined by several different experimental techniques (filled symbols) and theoretical calculations (open symbols); experimental points: coherent QENS (this work, squares), incoherent QENS ([4], triangles down), IR transient hole burning ([11], triangles up), depolarized Rayleigh light scattering ([7], circles). Theoretical points obtained by means of molecular dynamics and the reactive flux correlation function approach [23], using the SPC model of water: hydrogen bond lifetime (circles), hydrogen bond reforming time (triangles up), time of switching hydrogen bond partners [32]. Slopes represent activation energies between 8 and 11  $\text{kJ mol}^{-1}$ .



## 2- Structural relaxation in supercooled (Na<sub>2</sub>O–Li<sub>2</sub>O)–2P<sub>2</sub>O<sub>5</sub> : a neutron spin-echo study

B. Rufflé

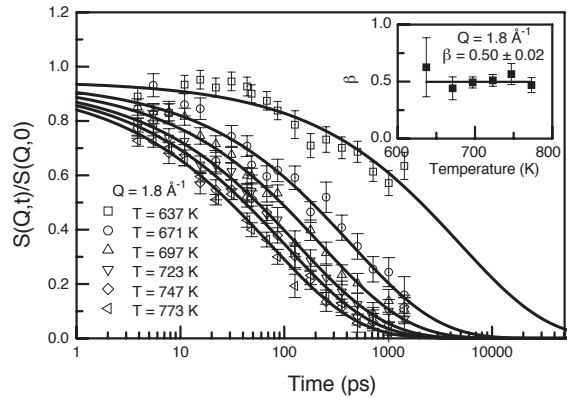


Fig. 1. Intermediate scattering function  $\phi(Q, t) = S(Q, t)/S(Q, t = 0)$  measured on (Na<sub>2</sub>O–Li<sub>2</sub>O)–2P<sub>2</sub>O<sub>5</sub> at  $Q = 1.8 \text{ \AA}^{-1}$  for six temperatures ranging from 637 to 773 K. Lines are the resulting KWW fit curves (Eq. (1)) described in the text. The inset shows the temperature dependence of the Kohlrausch stretching exponent.

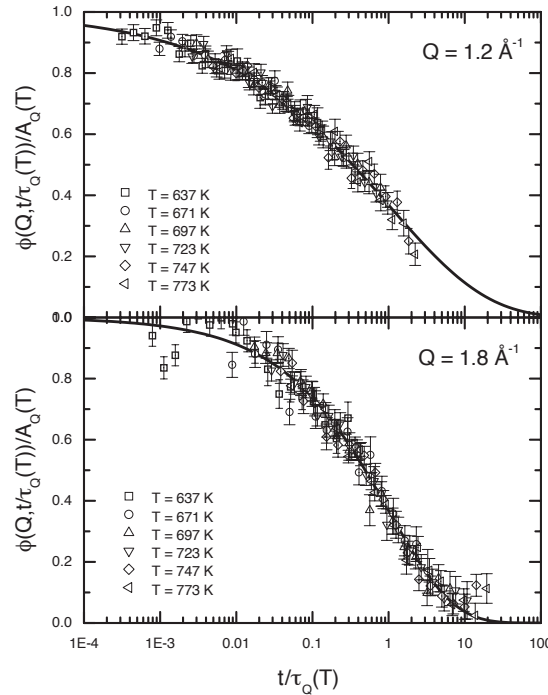


Fig. 2. Scaling analysis of the  $\alpha$  relaxation in (Na<sub>2</sub>O–Li<sub>2</sub>O)–2P<sub>2</sub>O<sub>5</sub>: master curves  $\phi(Q, t/\tau_Q(T))/A_Q(T)$ , where  $A_Q(T)$  and  $\tau_Q(T)$  are obtained from iterative fits with Eq. (1).

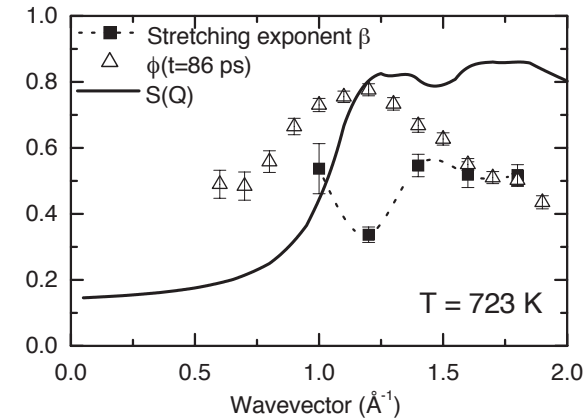


Fig. 3. Dependence on scattering vector  $Q$ : (■) stretching exponent  $\beta_Q$  obtained by the three-step procedure described in the text; ( $\Delta$ ) intermediate scattering function at a fixed time  $\phi(t = 86 \text{ ps})$  revealing a maximum of the structural relaxation time around  $Q = 1.2 \text{ \AA}^{-1}$ , the position of the pre-peak; (solid line) static structure factor  $S(Q)$  of (Na<sub>2</sub>O–Li<sub>2</sub>O)–2P<sub>2</sub>O<sub>5</sub> at  $T = 723 \text{ K}$ .

$$\Phi(\tau = 86ps) = \frac{S(Q, \tau = 86ps)}{S(Q, \tau = 0)}$$