

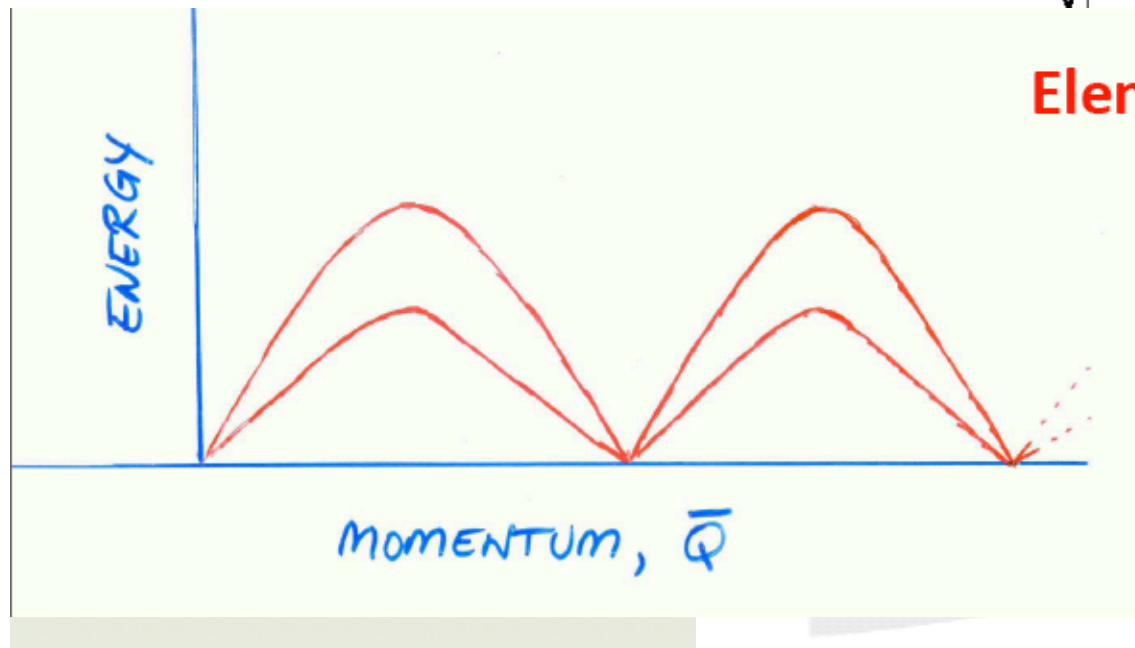
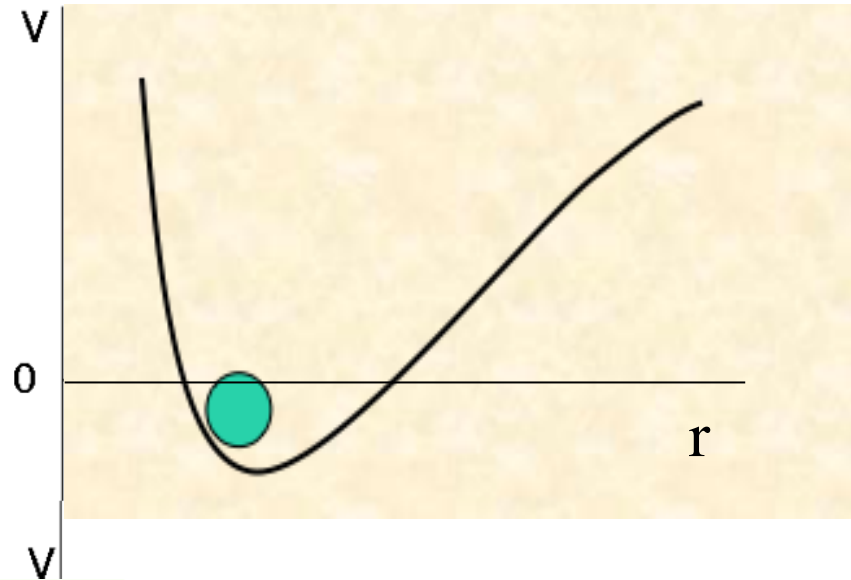
Inelastic Scattering I

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e
Centro NAST

INELASTIC NEUTRON SCATTERING

Fundamental Information on Interactions in Materials

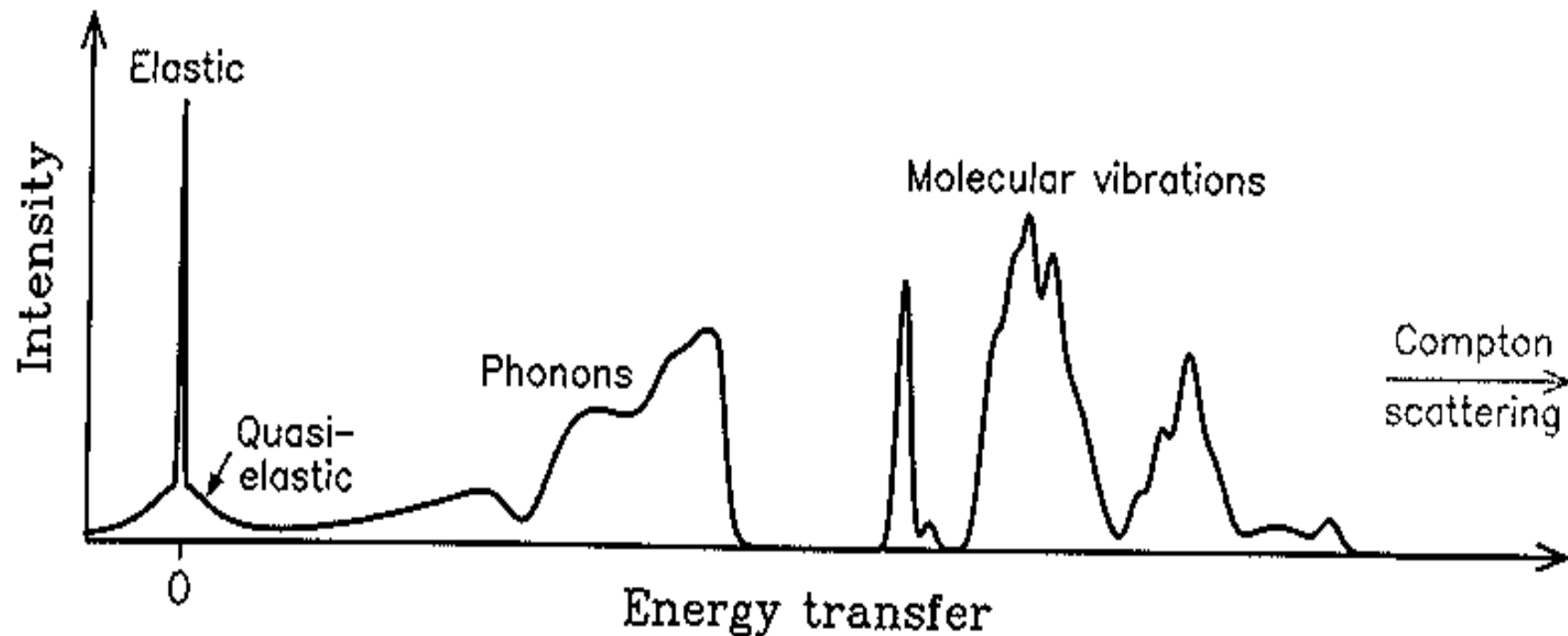
Structural probes yield indirect information on interactions in materials by locating the minimum of the potential. Dynamical probes, including neutron scattering, reveal information on the shape of the potential.



Elementary Excitations in Solids

- Lattice Vibrations (Phonons)
- Spin Fluctuations (Magnons)

A sketch of collective and single-particle excitations



From: "Elementary Scattering Theory For X-ray and Neutron Users" D.S. Sivia OUP (2011)

XII School of Neutron Scattering F. P. Ricci at Ettore Majorana Foundation and Centre for Scientific Culture (2014)

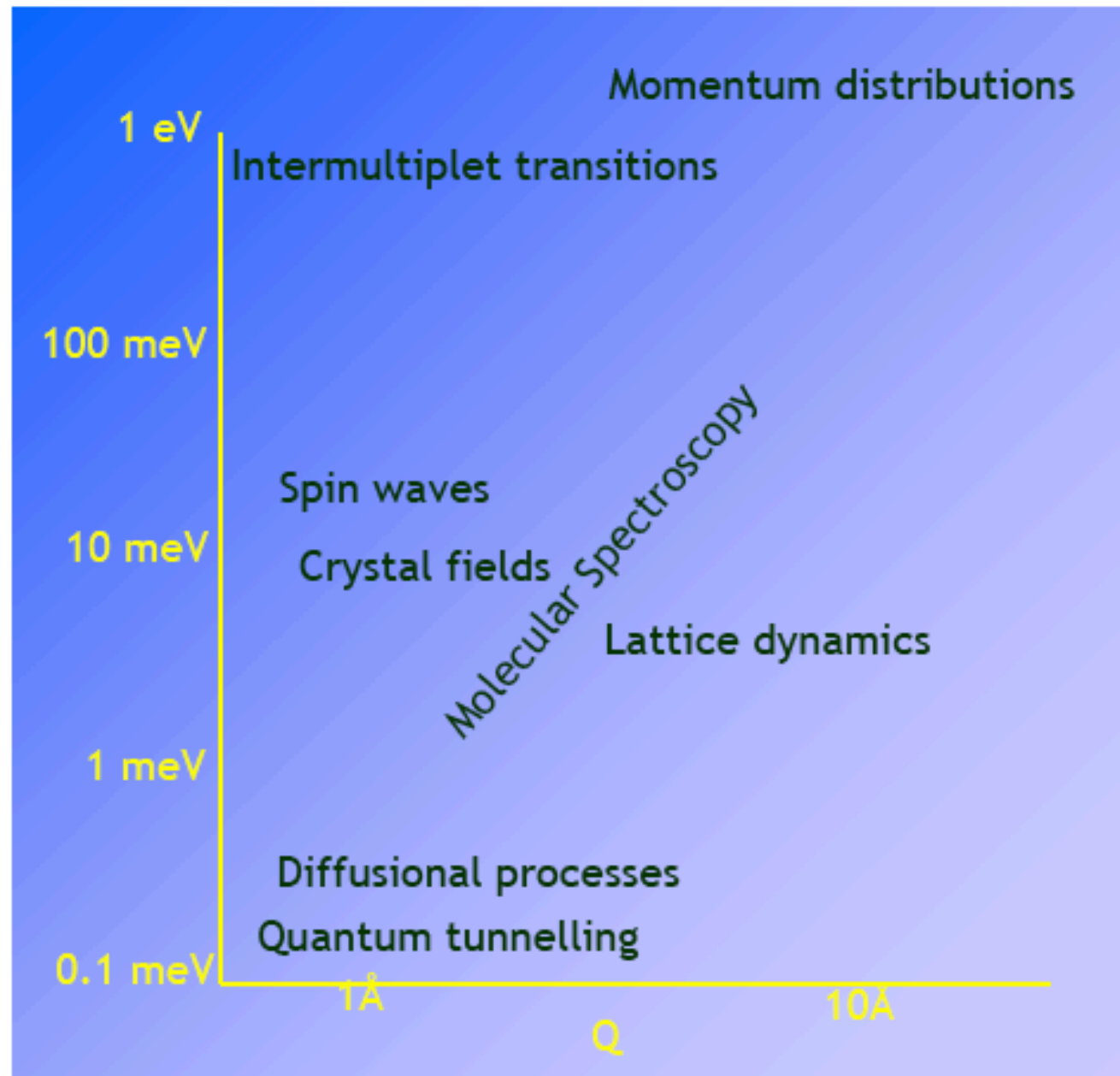
INELASTIC NEUTRON SCATTERING EXPERIMENT

- In the inelastic neutron scattering experiment, the quantity we measure is the double differential cross section

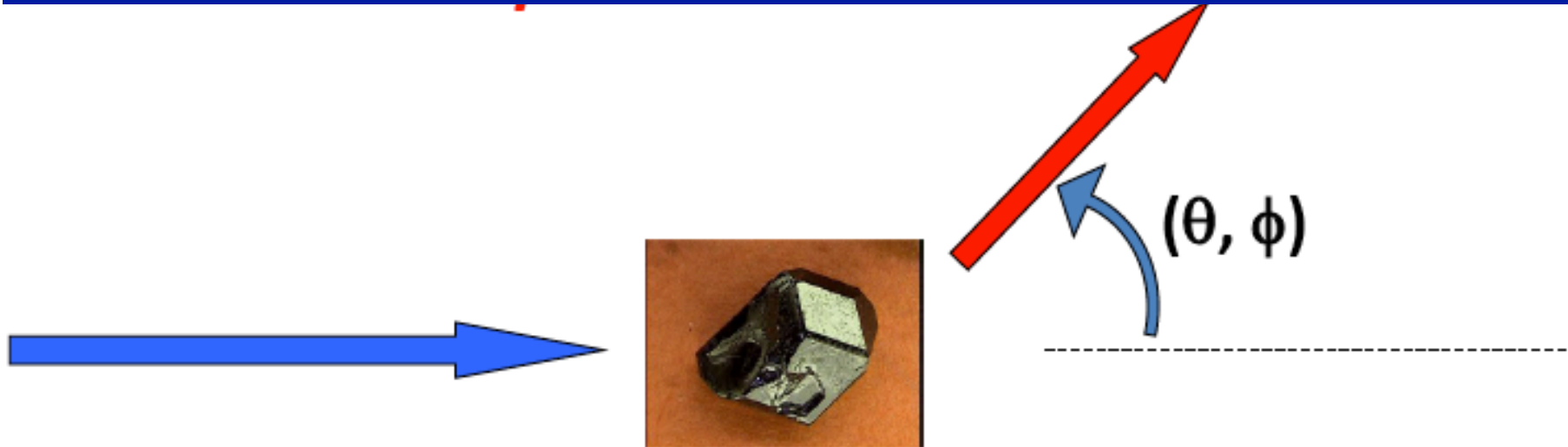
$$\frac{d^2\sigma}{d\Omega dE} = \frac{k_f}{k_i} b^2 S(\mathbf{Q}, \omega)$$

- It is the scattering function $S(\mathbf{Q}, \omega)$ which provides the link between the scattering data and the physical system being studied.
- The type of experiment will dictate the portions of (\mathbf{Q}, ω) space which are to be probed.

Regions of interest in (Q, ω) space



INELASTIC NEUTRON SCATTERING: THE BASIC EXPERIMENT



Incident Beam:

- monochromatic
- “white”

Scattered Beam:

- Resolve its energy
- Don't resolve its energy
- Filter its energy

ANALYSING THE SCATTERING TRIANGLE

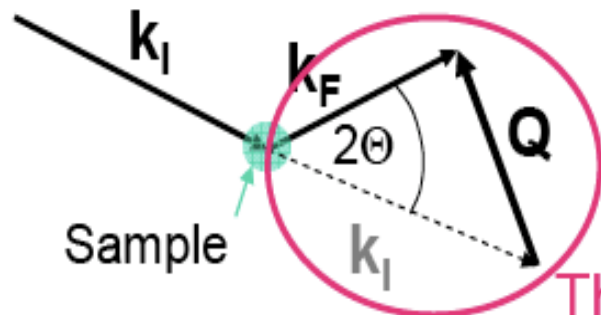
Determine the probability of finding a given change $\hbar\mathbf{Q}$ from the momentum $\hbar\mathbf{k}_i$ of a neutron incident on the specimen to the momentum $\hbar\mathbf{k}_f$ of the neutron scattered from the specimen.

In other words:

Measure the momentum transfer

$$\hbar\mathbf{Q} = \hbar\mathbf{k}_i - \hbar\mathbf{k}_f$$

$$\text{or: } \mathbf{Q} = \mathbf{k}_i - \mathbf{k}_f \quad (\text{wave number notation})$$



2θ is called the “scattering angle”

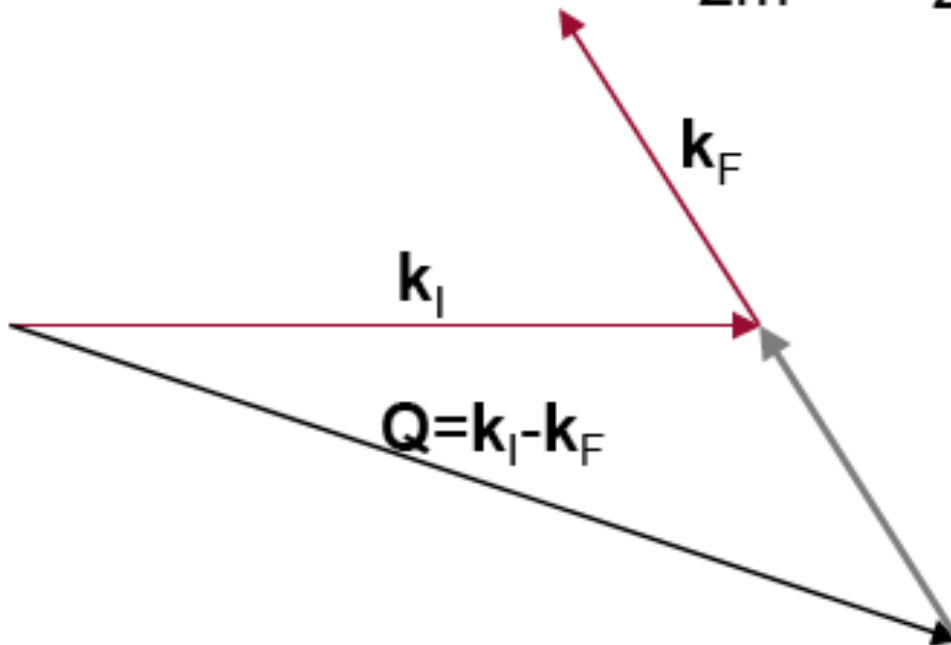
The scattering triangle

ANALYSING THE SCATTERING TRIANGLE

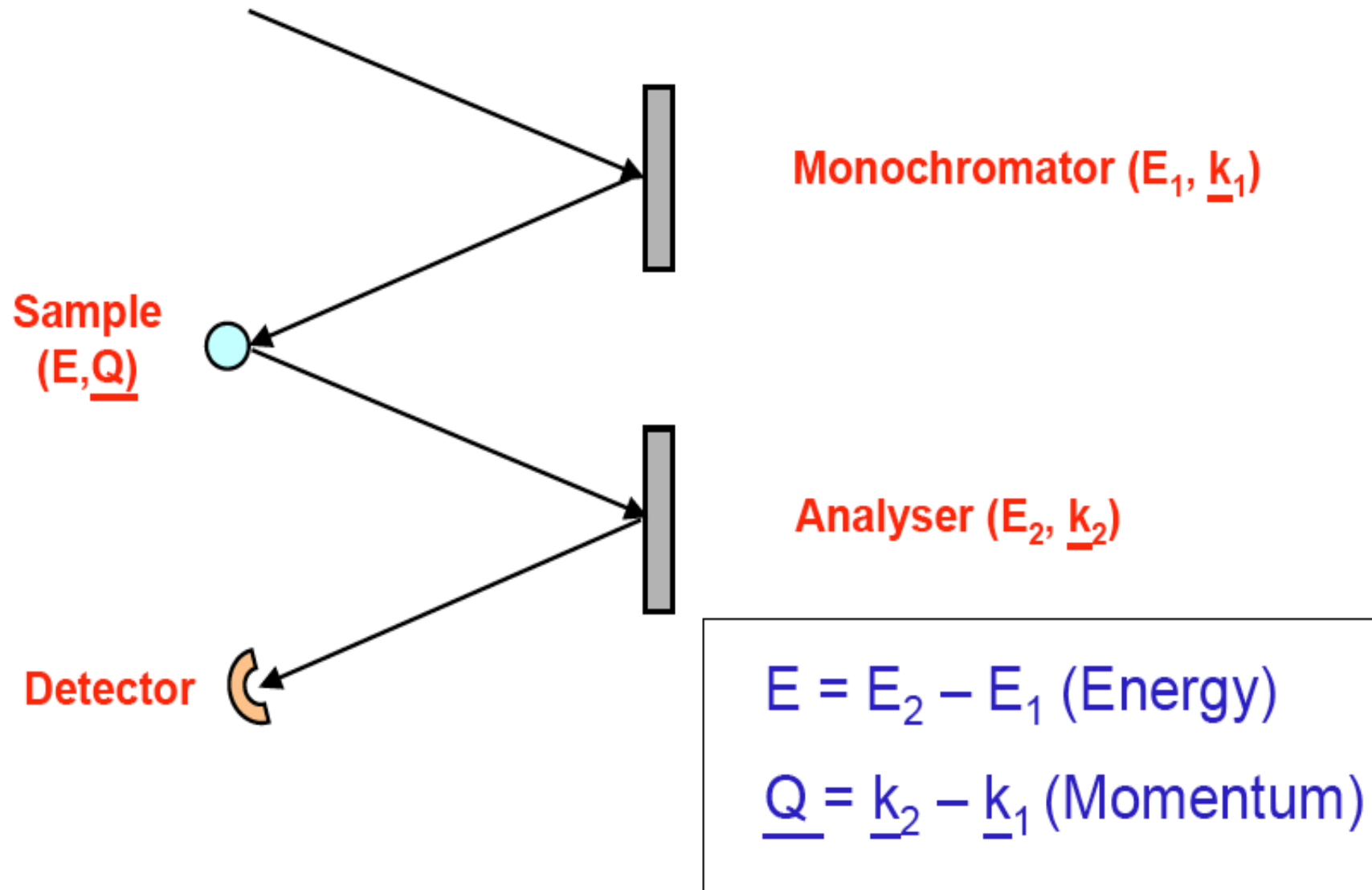
Basic relations:

momentum transfer: $\mathbf{Q} = \mathbf{k}_I - \mathbf{k}_F$

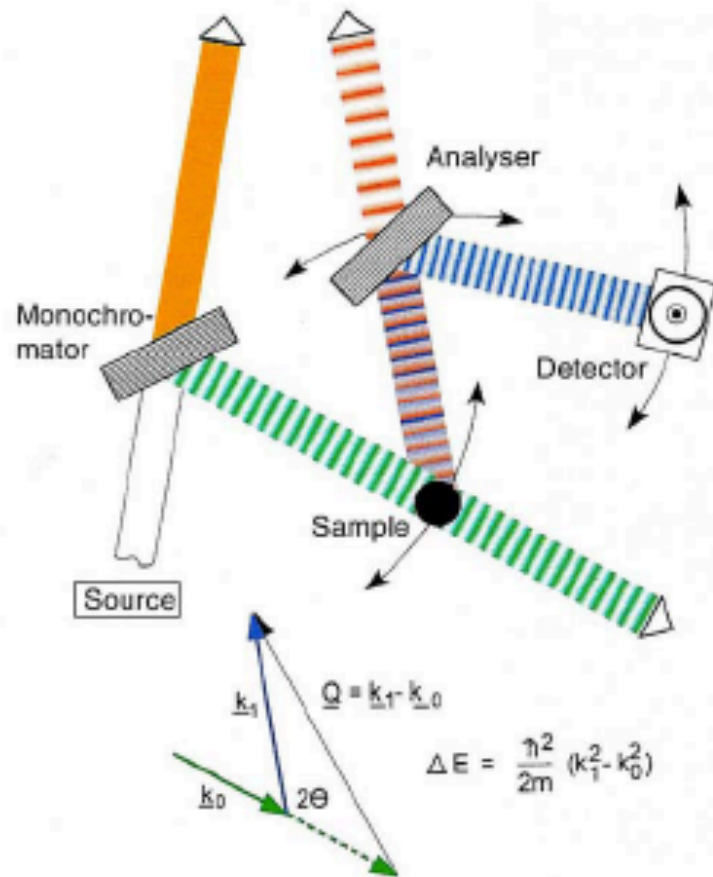
energy transfer: $\Delta E = \frac{\hbar^2 k_I^2}{2m} - \frac{\hbar^2 k_F^2}{2m} = \hbar\omega$



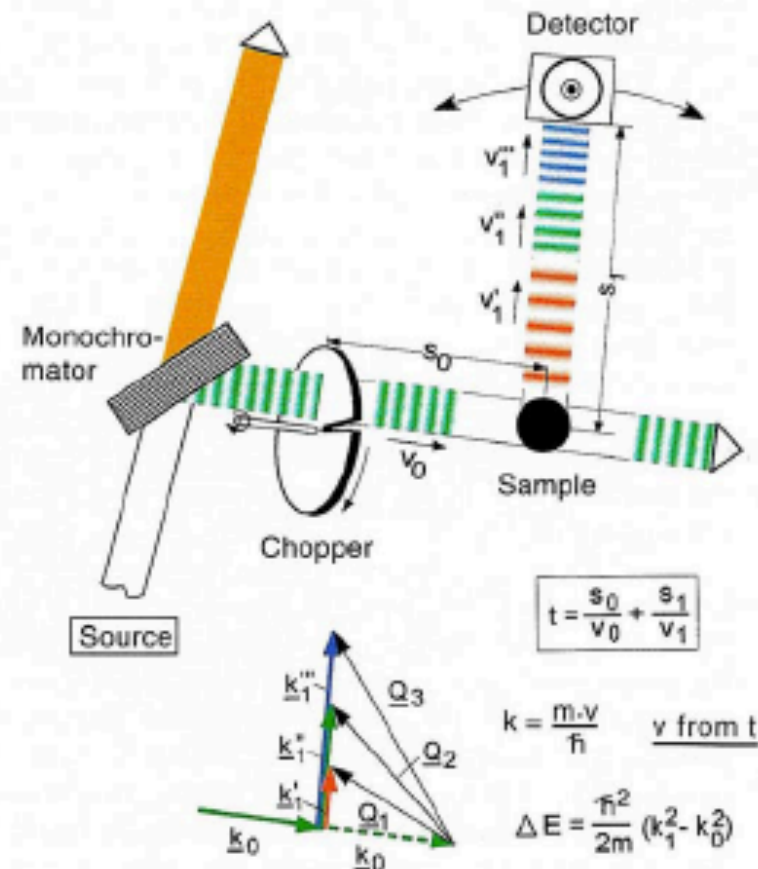
INELASTIC NEUTRON SCATTERING EXPERIMENT



NEUTRON SPECTROSCOPY ON A CONTINUOUS SOURCE

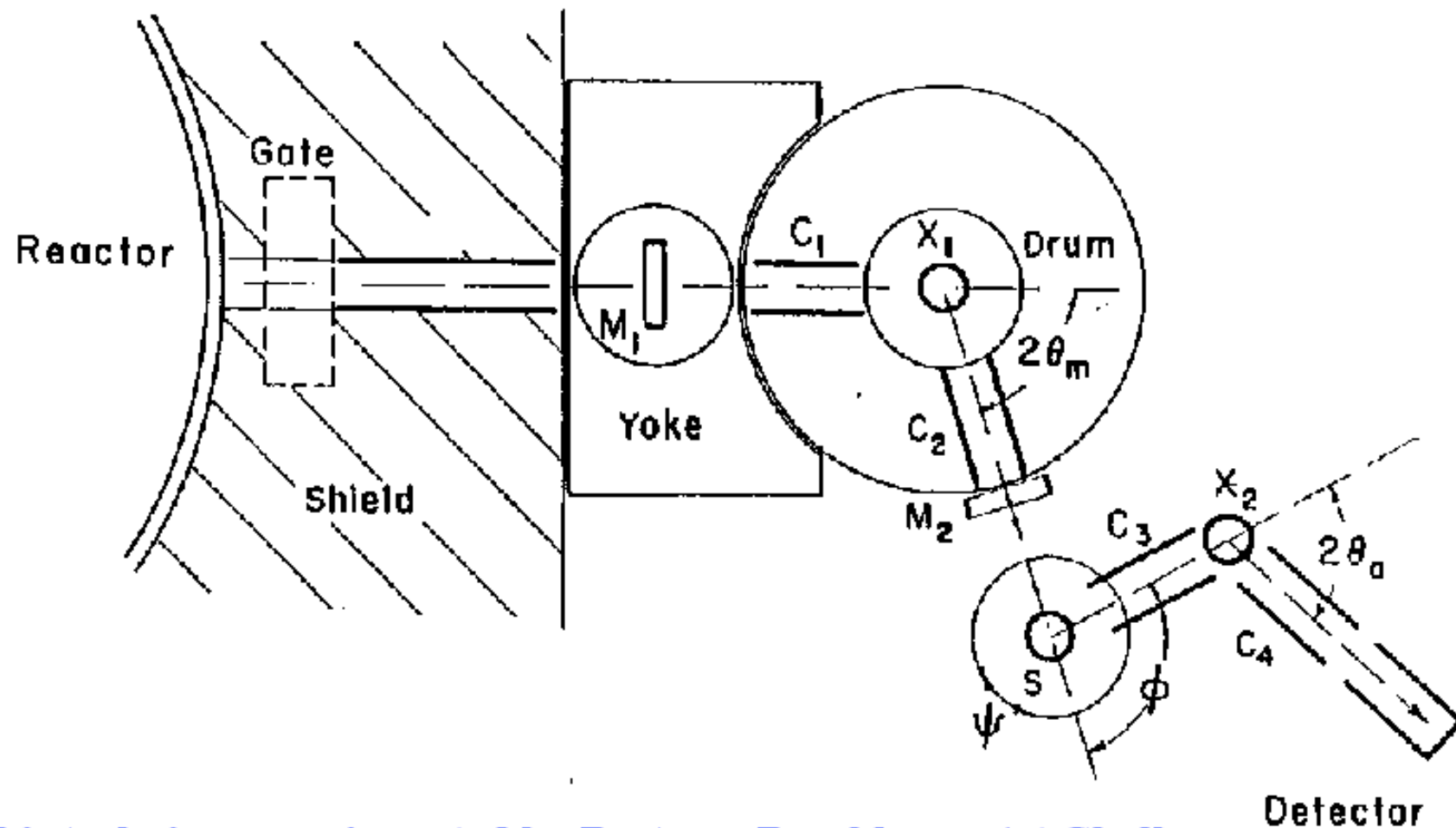


The triple axis spectrometer measures one Q-value at a time. For a scan two of the three axes must be activated



The time of flight spectrometer allows to measure several Q values at each detector setting simultaneously. Multi-detectors are generally used.

Triple Axis Spectrometer

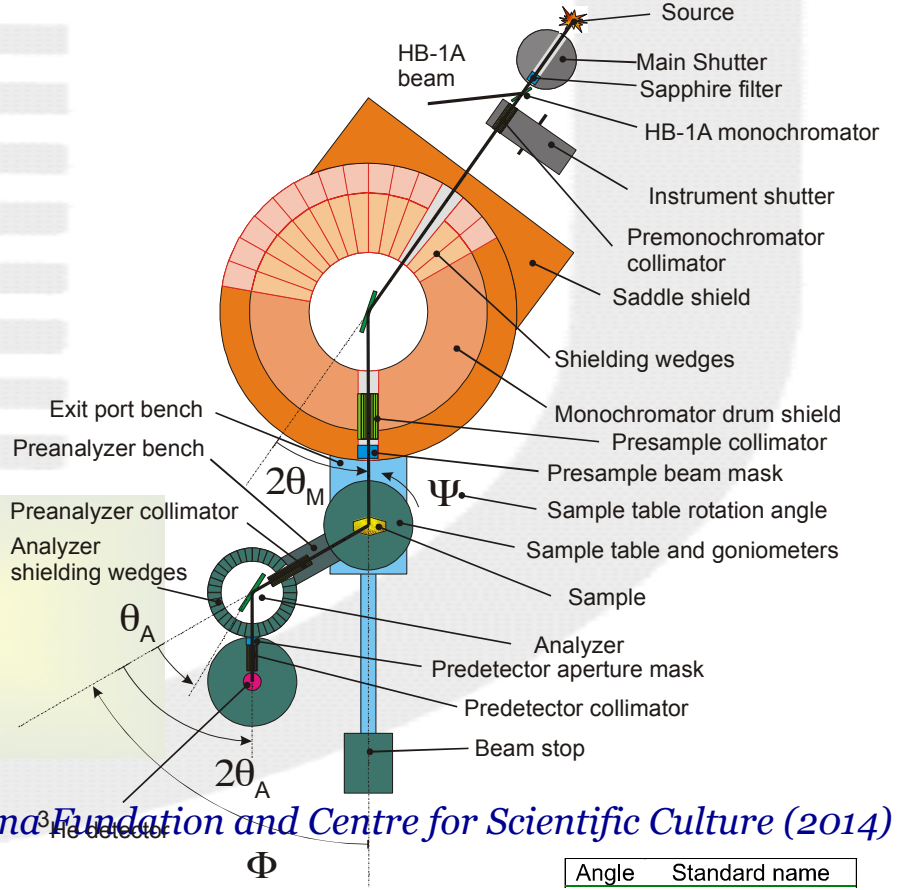


This technique was invented by Bertram Brockhouse (at Chalk River National Lab in Canada) [Nobel Laureate in 1994] in the 1950s and although the instruments have improved beyond recognition, the basic technique has not changed.

TRIPLE AXIS SPECTROMETER

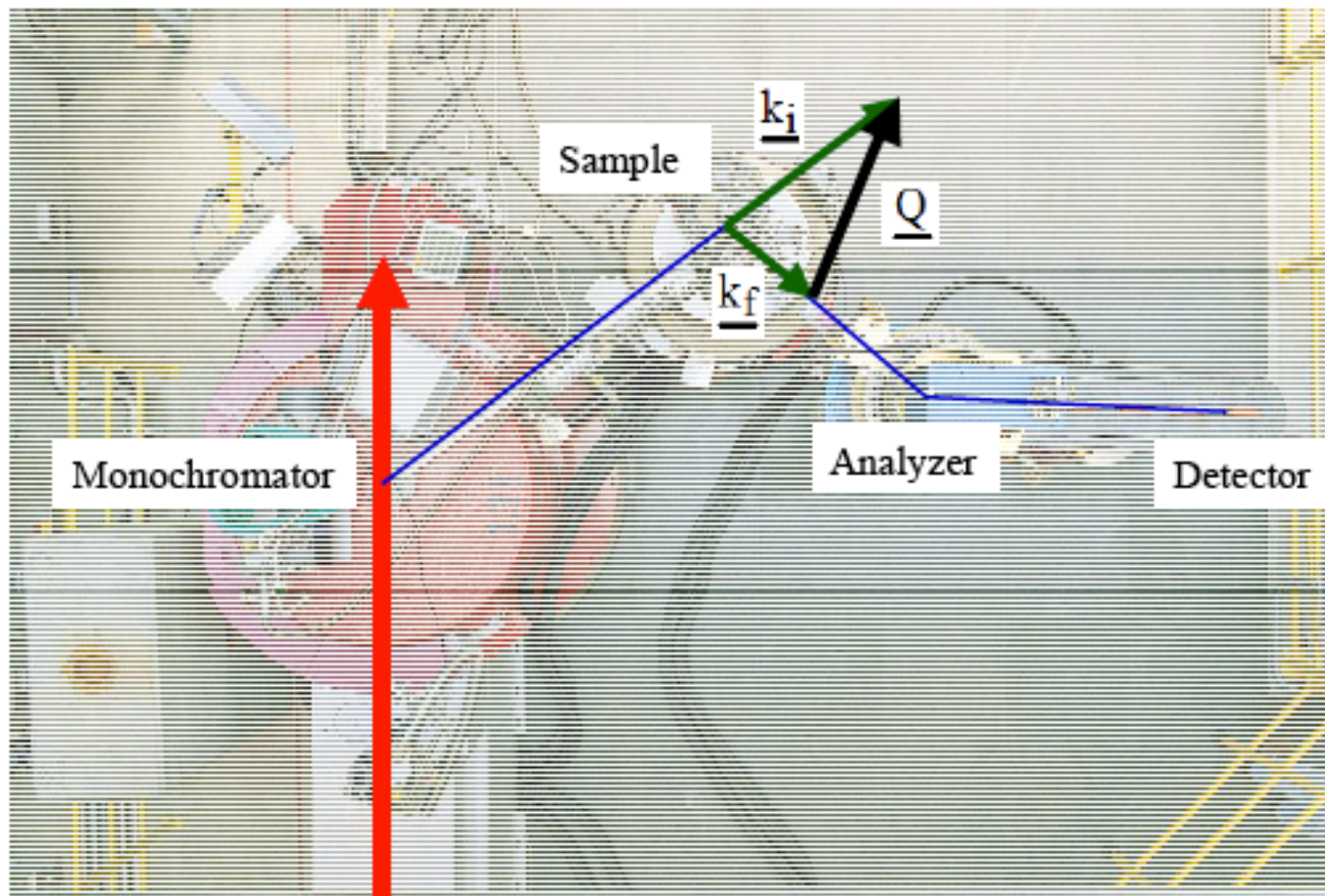
“old” HB3, HFIR

“new” HB1, HFIR



na Foundation and Centre for Scientific Culture (2014)

TRIPLE AXIS SPECTROMETER



$$|\underline{k}| = \frac{2 \cdot \pi}{\lambda}$$

$$E_i = \frac{\hbar^2 \cdot \underline{k}_i^2}{2 \cdot m}$$

$$E_f = \frac{\hbar^2 \cdot \underline{k}_f^2}{2 \cdot m}$$

$$E = \hbar \omega$$

Energy transfer to sample

$$E = \frac{\hbar^2}{2 \cdot m} [\underline{k}_i^2 - \underline{k}_f^2]$$

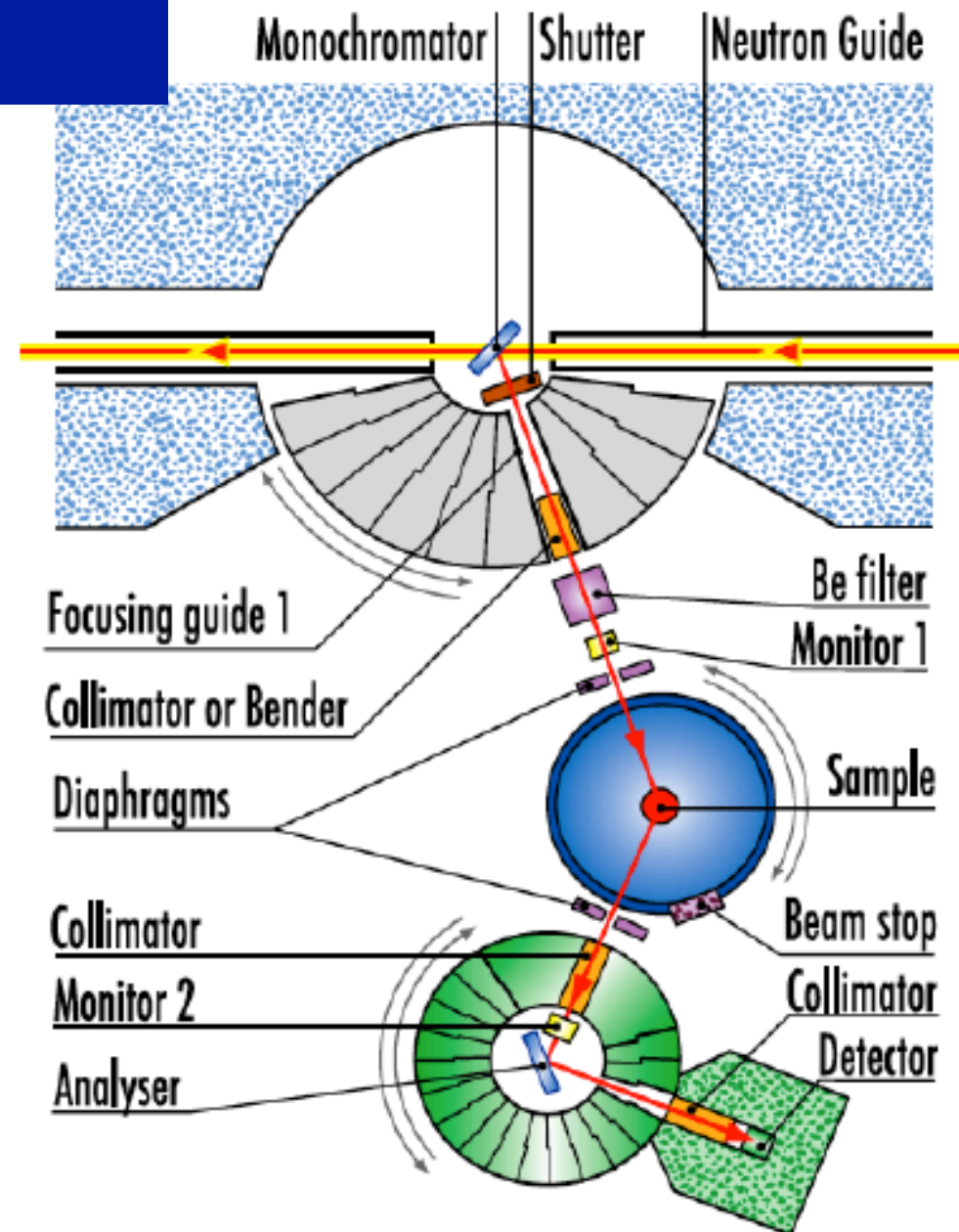
Wave vector transfer to sample

$$\underline{Q} = \underline{k}_i - \underline{k}_f$$

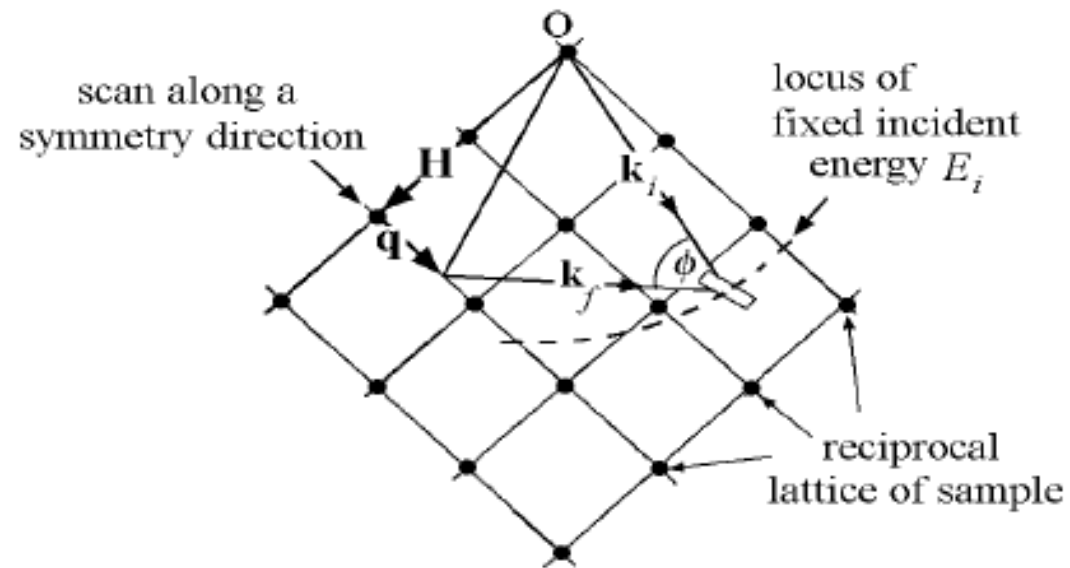
SCHEMATIC OF A TRIPLE AXIS SPECTROMETER

Can be operated in constant- Q or constant E mode.

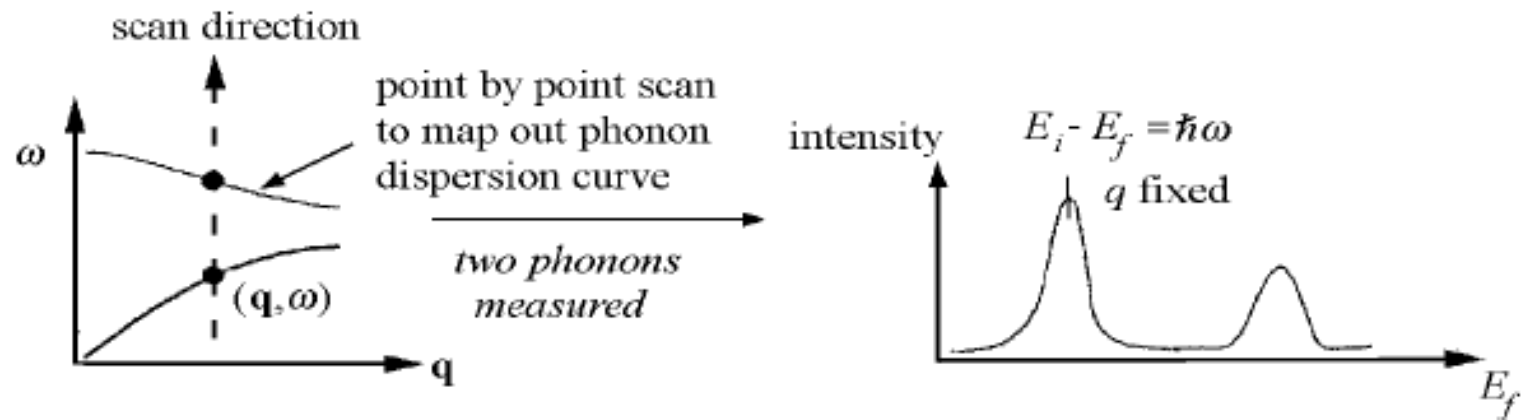
Note that the Be filter, which is used with cold neutrons, can either be before the sample (constant k_i) or after the sample (constant k_f)



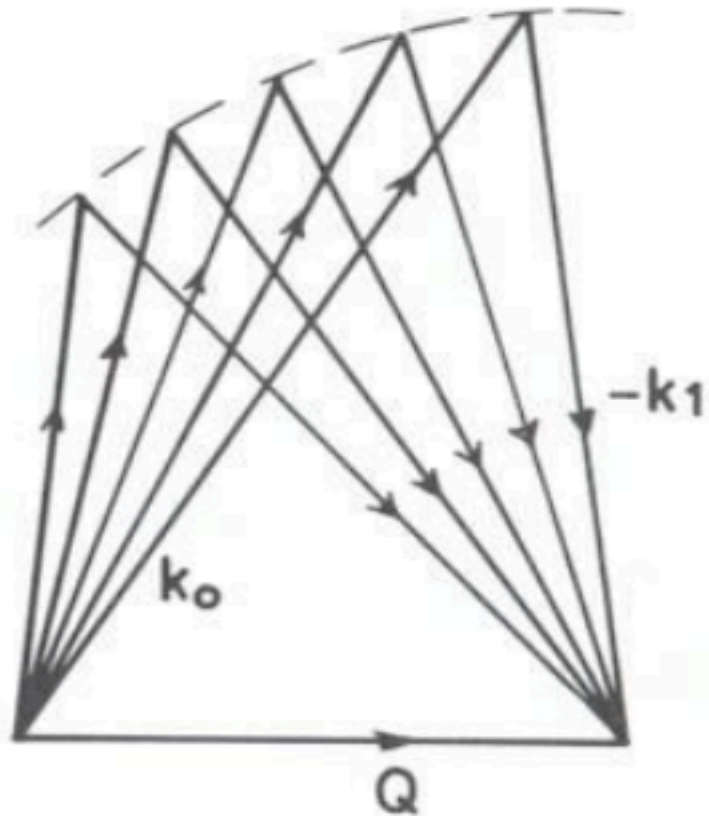
TRIPLE AXIS SPECTROMETER



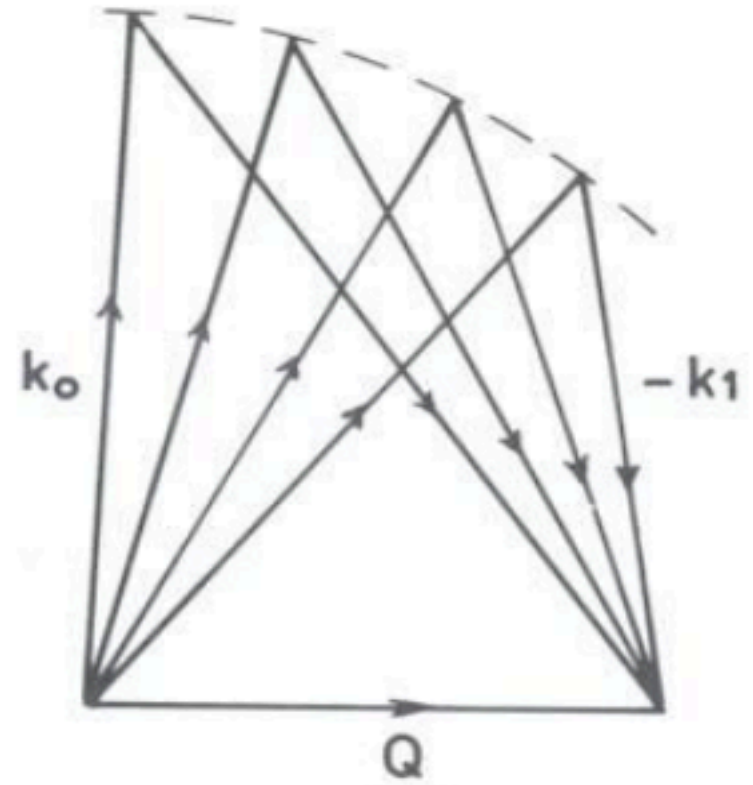
Focusing



TWO DIFFERENT WAYS OF PERFORMING CONSTANT-Q SCAN



Constant k_f (ii)

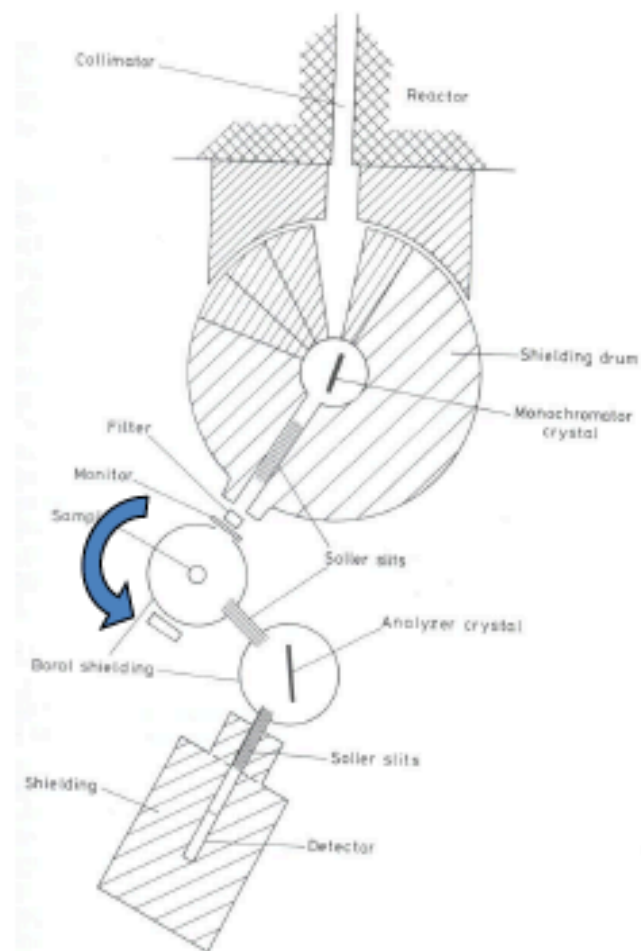
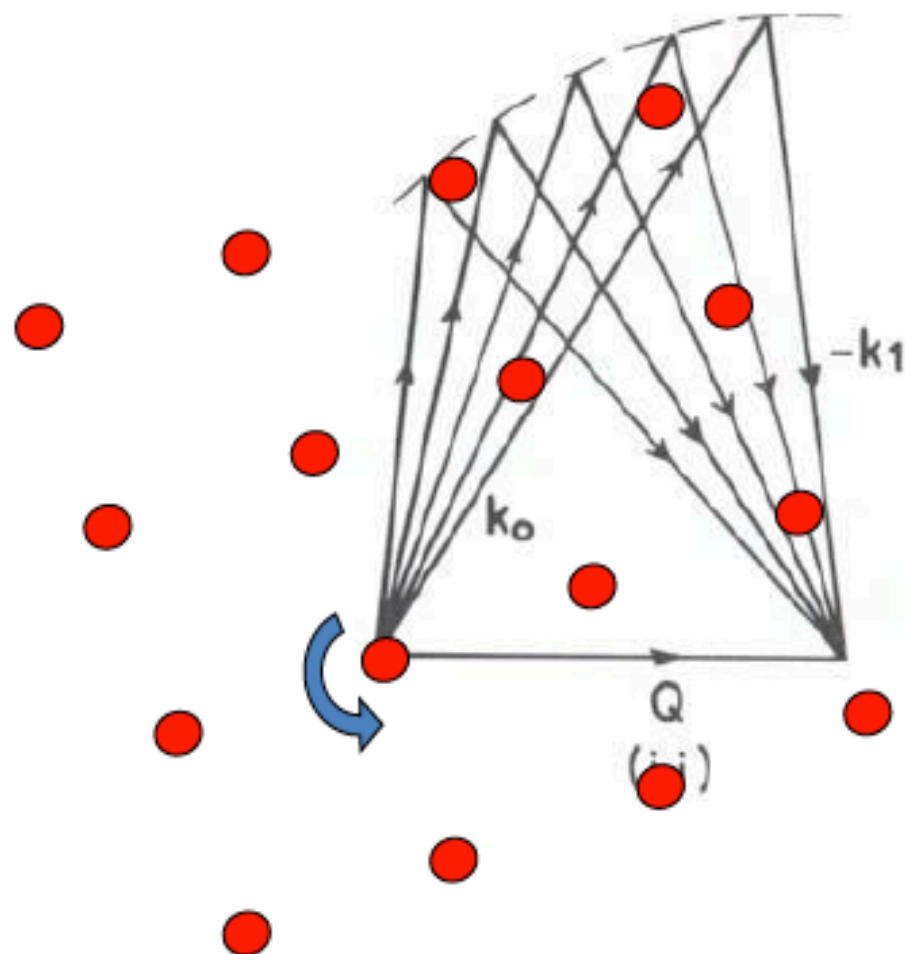
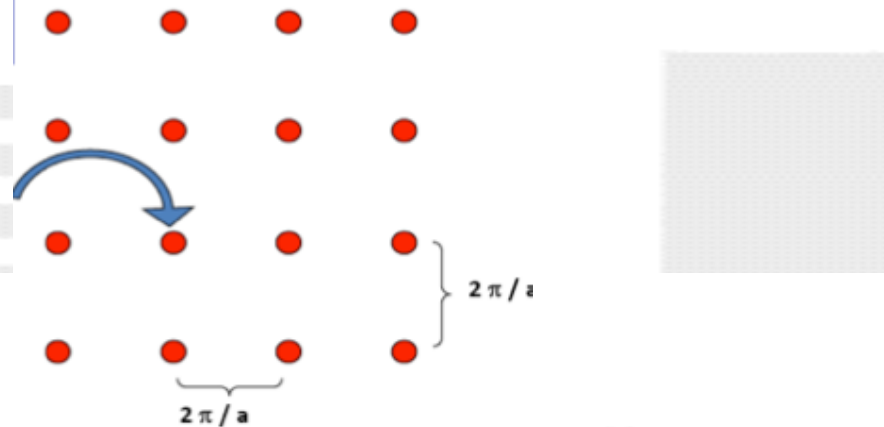


Constant k_i (i)

MAPPING MOMENTUM-ENERGY (Q-W) SPACE

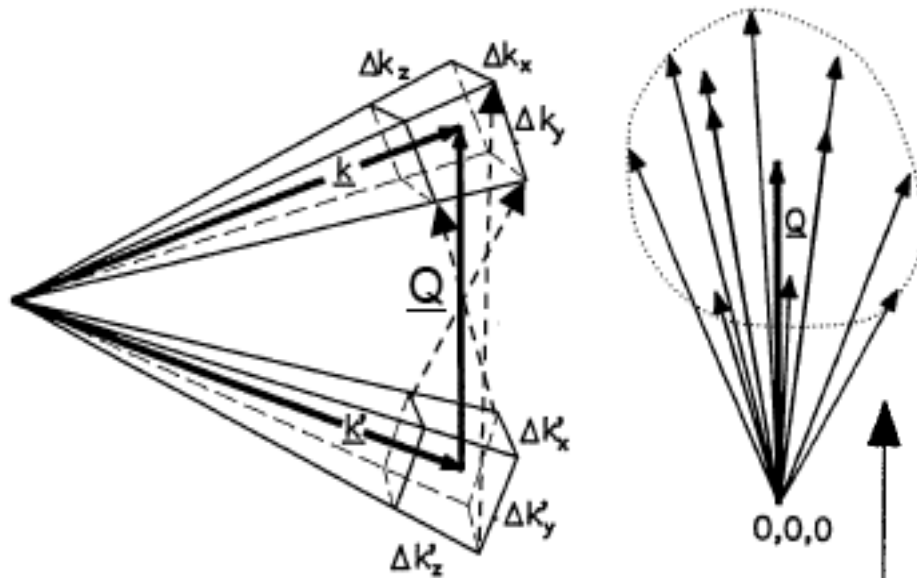
Origin of reciprocal space;

Remains fixed for any sample rotation



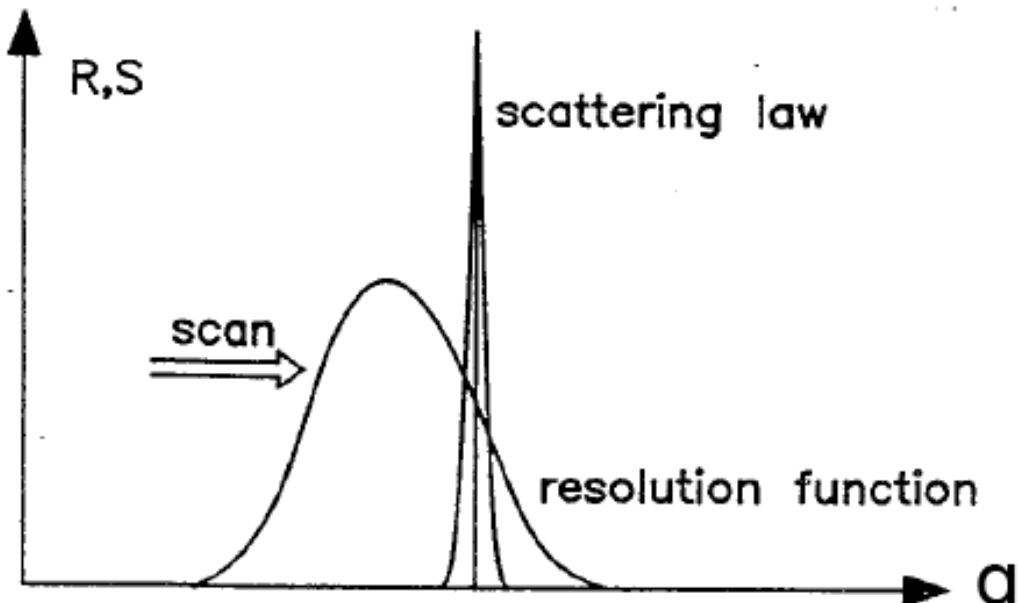
RESOLUTION FUNCTION

The next important question after what can be measured is, **how precisely it can be measured.**



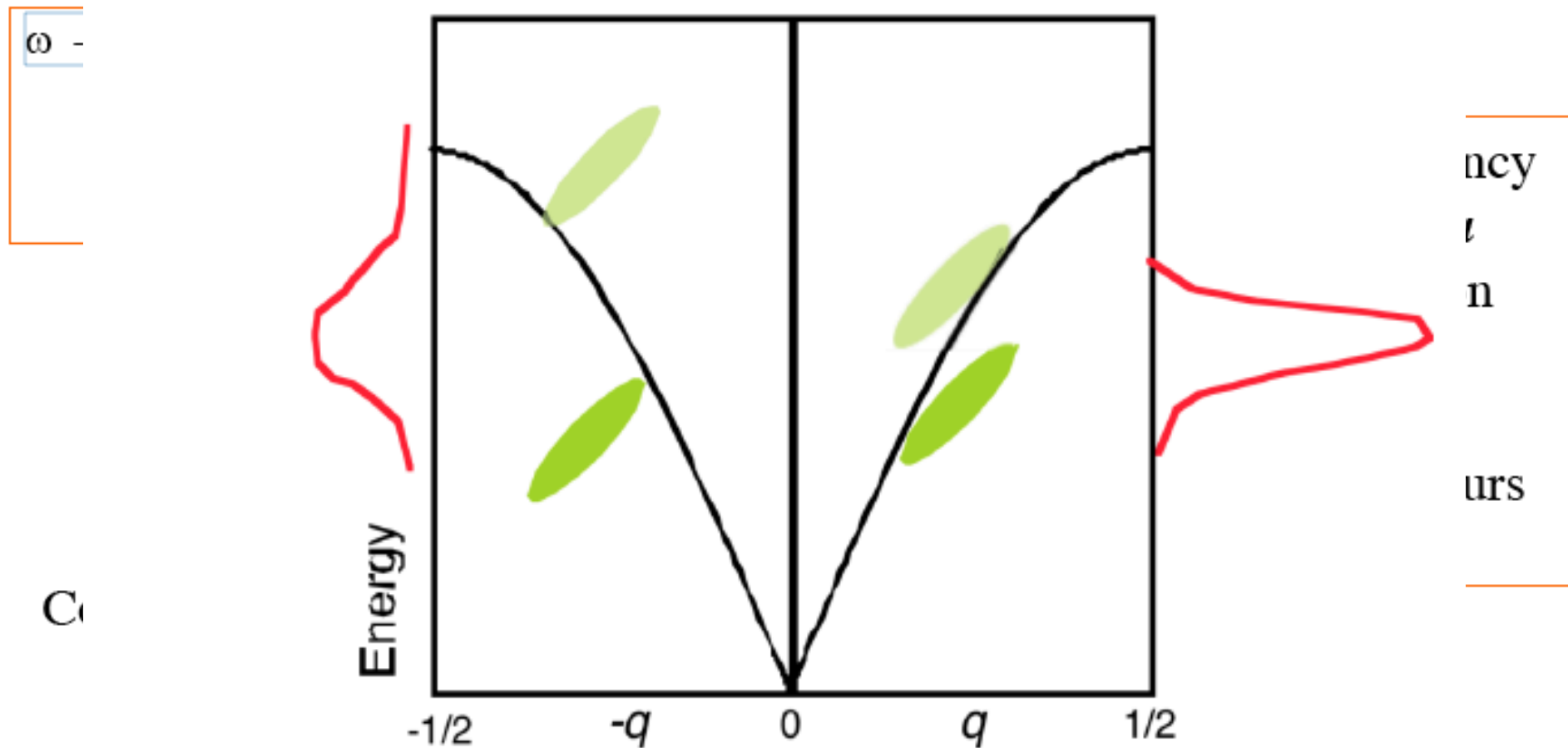
There is always a degree of uncertainty, as to how well K_i and K_f can be defined by experimental means.

Scanning a wide resolution function across a narrow scattering law essentially reproduces the resolution function



TRIPLE AXIS SPECTROMETER

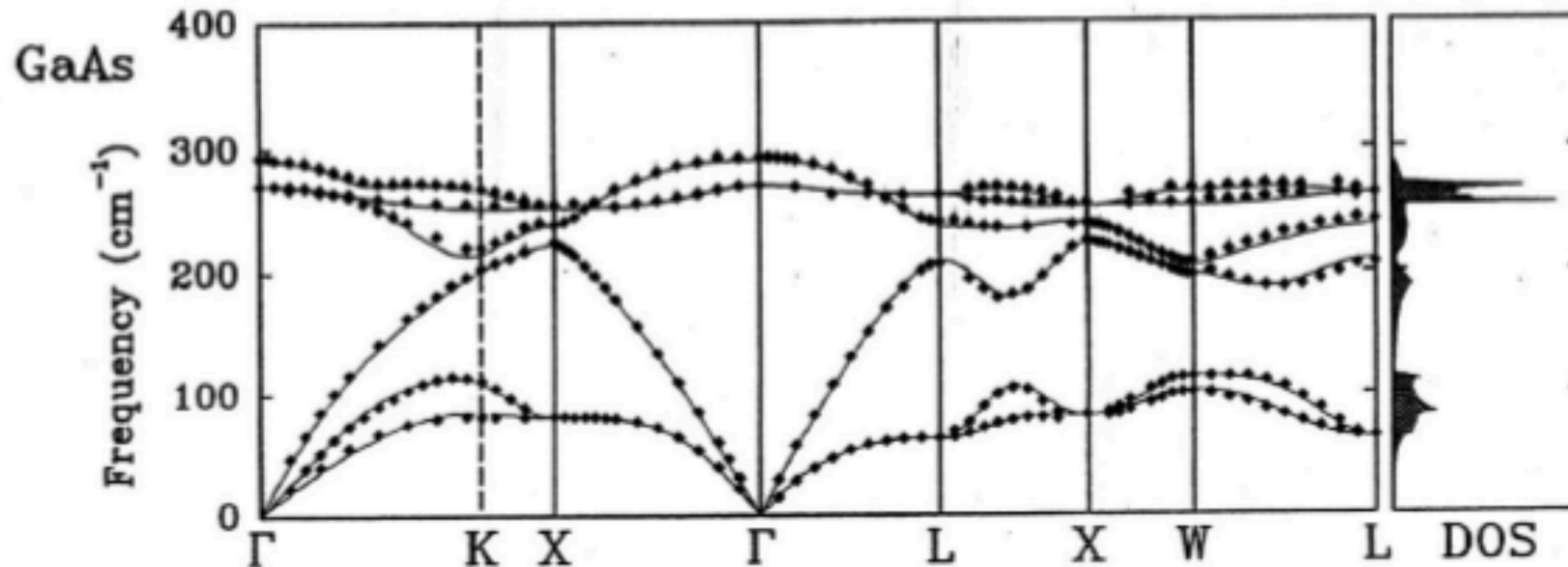
- Optimizing peak intensity
- Match slope of resolution to dispersion



The IN14 triple-axis spectrometer at the ILL

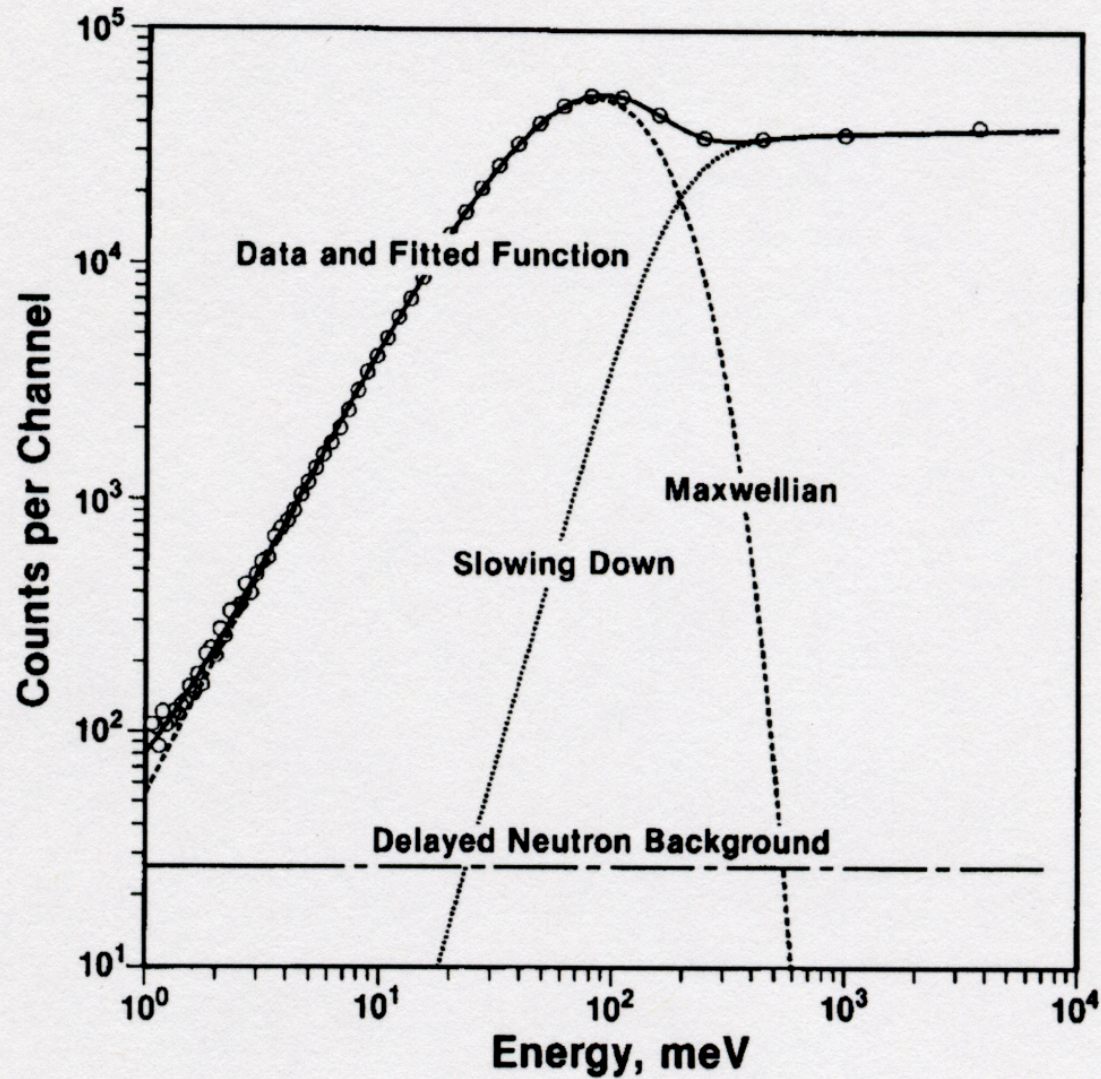


Phonons in Gallium Arsenide, a III/V semiconductor, is used in the manufacture of devices such as microwave frequency integrated circuits

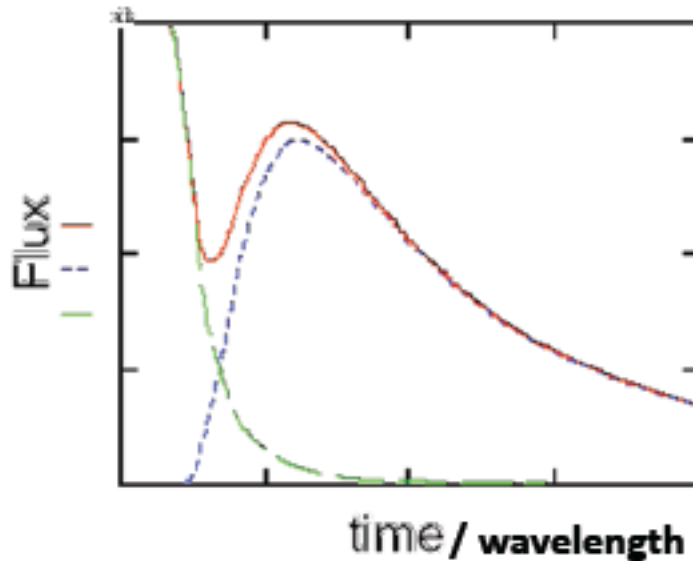


Phonon dispersion curves as measured for GaAs. Strauch & Dornier (1990). The lines give the result of *ab initio* calculations and show that the forces between the atoms are well understood. The letters below give the notation for the symmetry directions or points.

INELASTIC NEUTRON SCATTERING ON A PULSED SOURCE

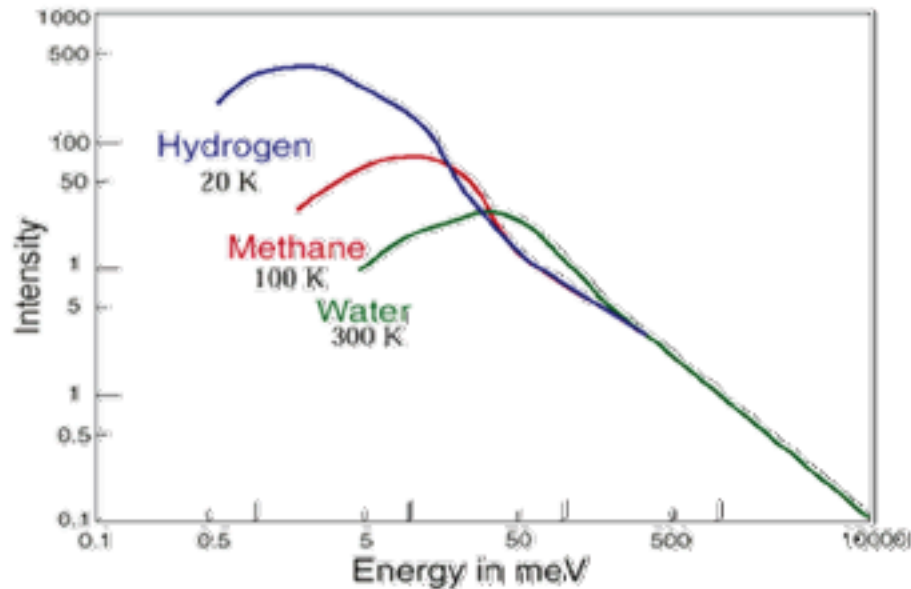


NEUTRON SPECTROSCOPY ON A PULSED SOURCE



Spallation source moderators under-moderate the neutrons to produce a high flux of epi-thermal neutrons

- Moderated Neutrons
- Under-moderated neutrons



The energy (or wavelength) range of the produced neutrons can be changed by controlling the temperature of the moderator

TIME OF FLIGHT (TOF) TECHNIQUES

- Velocity of thermal neutrons $\sim \text{km s}^{-1}$.
- Hence their energy can be determined by measuring time-of-flight over a distance of a few metres.
- All spectrometers at pulsed sources use time-of-flight. On steady state sources pulsing devices such as choppers are required.

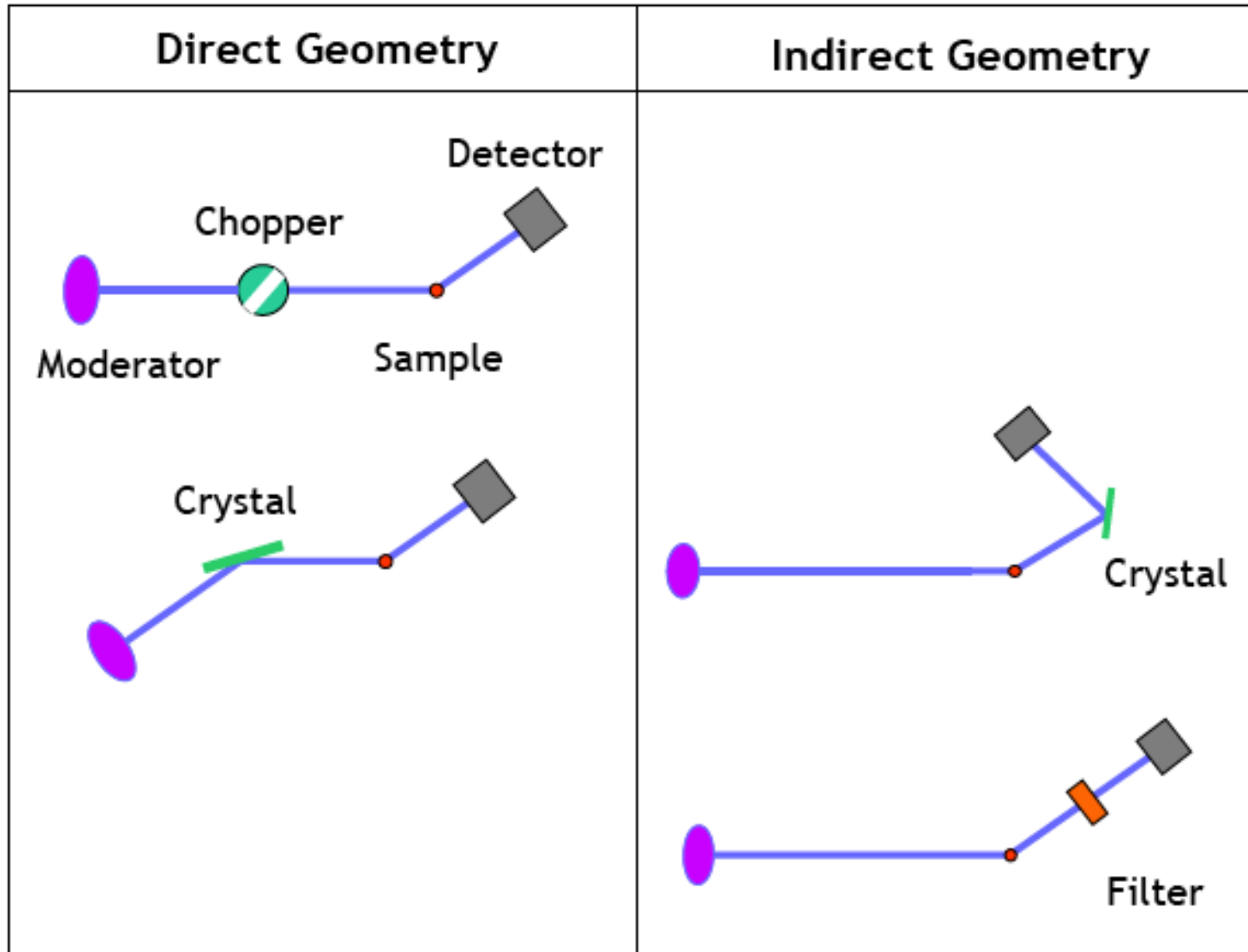
$$Q = k_i - k_f$$

$$\hbar\omega = E_i - E_f$$

TIME OF FLIGHT (TOF) SPECTROMETERS

- Time-of-flight spectrometers may be divided into two classes:-
 - **Direct geometry spectrometers:** in which the incident energy, E_i , is defined by a device such as a crystal or a chopper and the final energy, E_f , is determined by time of flight.
 - **Indirect (inverted) geometry spectrometers:** In which the sample is illuminated by a white incident beam and E_f is defined by a crystal or a filter and E_i is determined by time of flight

TIME OF FLIGHT (TOF) SPECTROMETERS



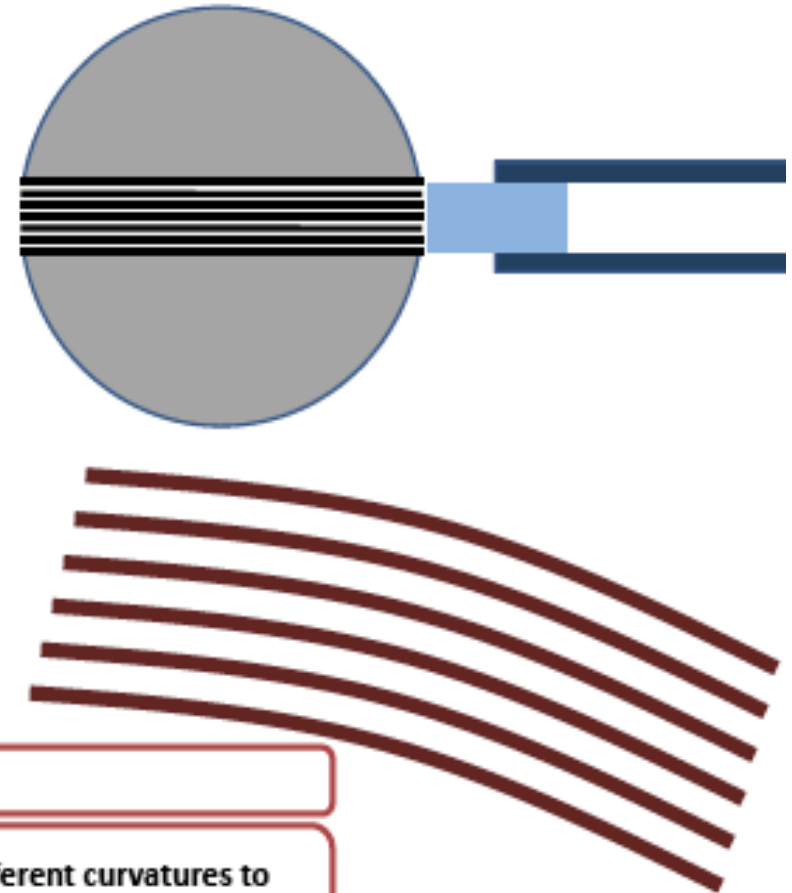
CHOPPERS

Fermi Chopper

The chopper is made of a series of curved slits

The curvature is chosen to match the energy range of interest

Resolution is determined by slit width and rotation speed



Advantages

By phasing the opening with the pulse from the target it is possible to select any neutron velocity

Rotation speeds up to 600Hz are possible

Usable over a wide energy range

Disadvantages

Need several rotors with different curvatures to cover the full energy range

Poor Transmission for low energy neutrons

Small area – only suitable for primary flight path

Filters

- The Be filter
 - Scatters neutrons above the Bragg cut-off (5.2meV)
 - Absorbing slits absorb the scattered neutrons



- Absorbs neutrons with energy above 5.2meV
- Cooling increases the filters effectiveness by reducing energy gain scattering from thermally excited phonons.

How to select neutrons in the range 1-100 eV ?

- ❖ Mechanical analysers (“Choppers”):

 - Neutrons too fast

- ❖ Crystal analysers

 - λ too short

Then...

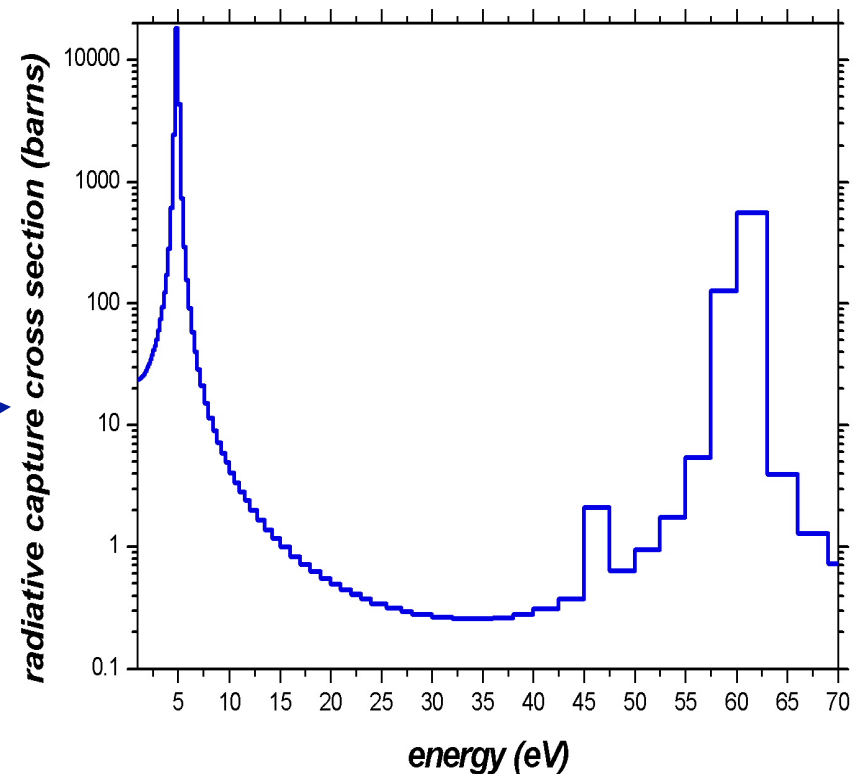
- ❖ Nuclear resonances

- ❖ Two configurations possible:

 - ❖ RFS (Resonance Filter Spectrometer)

 - ❖ RDS (Resonance Detector Spectrometer)

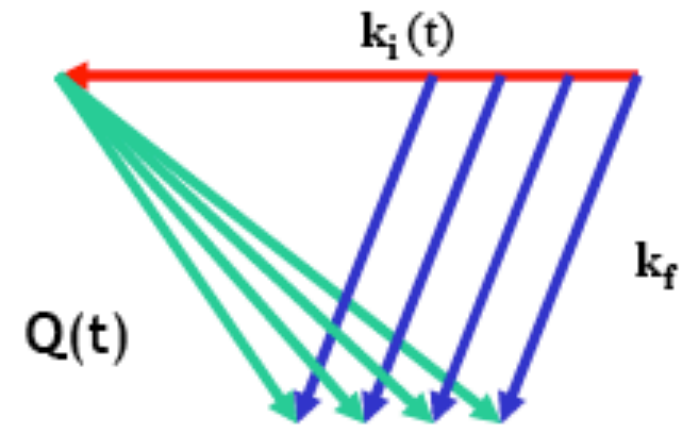
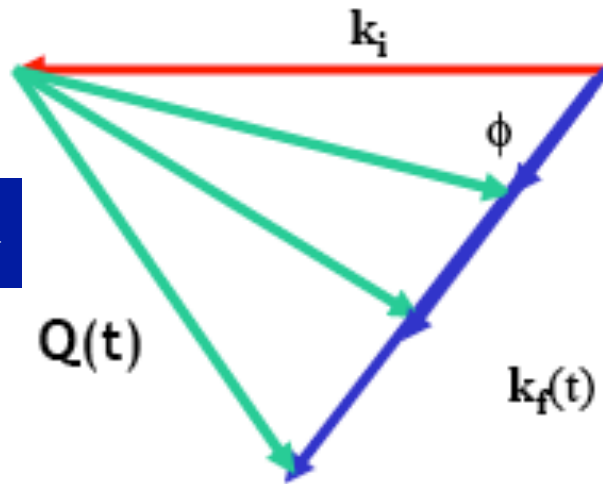
radiative capture cross section of ^{197}Au



ACCESSIBLE REGION OF (Q, ω) SPACE

- To understand the trajectories traced out through (Q, ω) consider the scattering triangles.

**DIRECT
GEOMETRY**



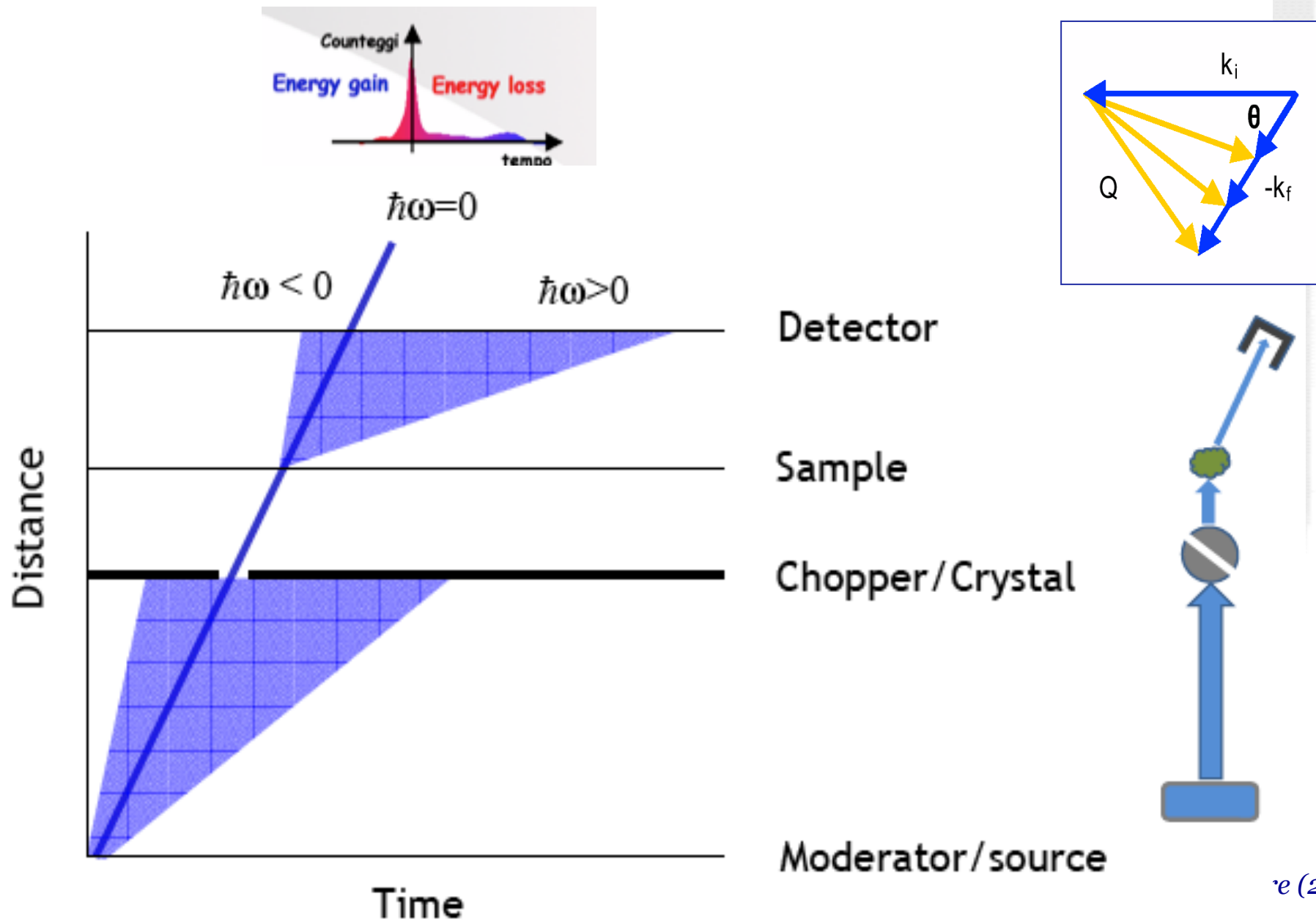
**INDIRECT
GEOMETRY**

- Applying the cosine rule

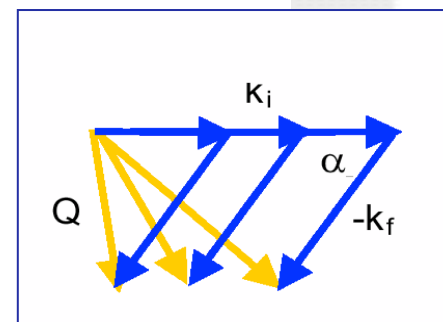
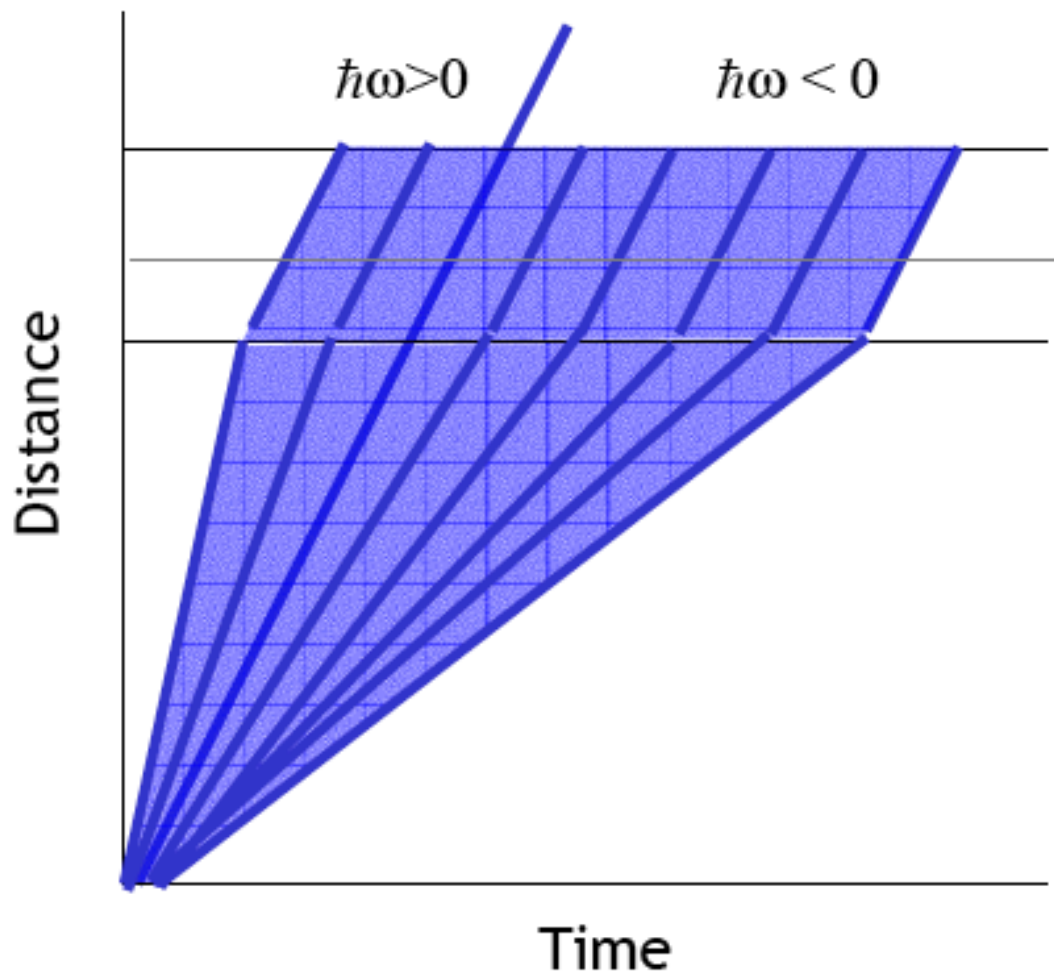
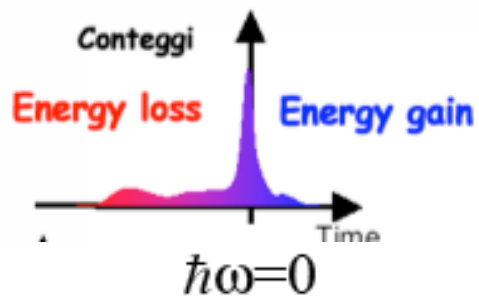
$$Q^2 = k_i^2 + k_f^2 - 2k_i k_f \cos \phi$$

$$\frac{\hbar^2 Q^2}{2m} = E_i + E_f - 2(E_i E_f)^{1/2} \cos \phi$$

SPECTROMETERS OPERATING IN DIRECT GEOMETRY



TIME OF FLIGHT (TOF) SPECTROMETER → INDIRECT GEOMETRY



Detector

Crystal / Filter
Sample

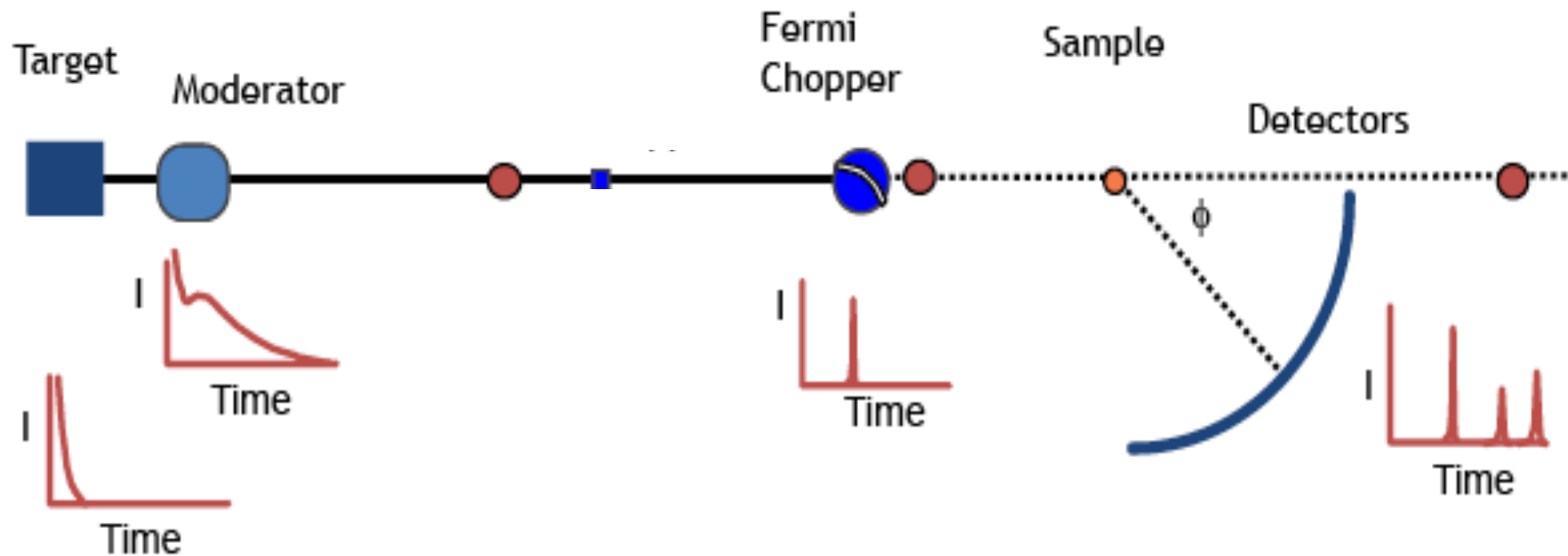
Moderator/source



SPECTROMETERS OPERATING IN DIRECT GEOMETRY

A chopper spectrometer on a pulsed source

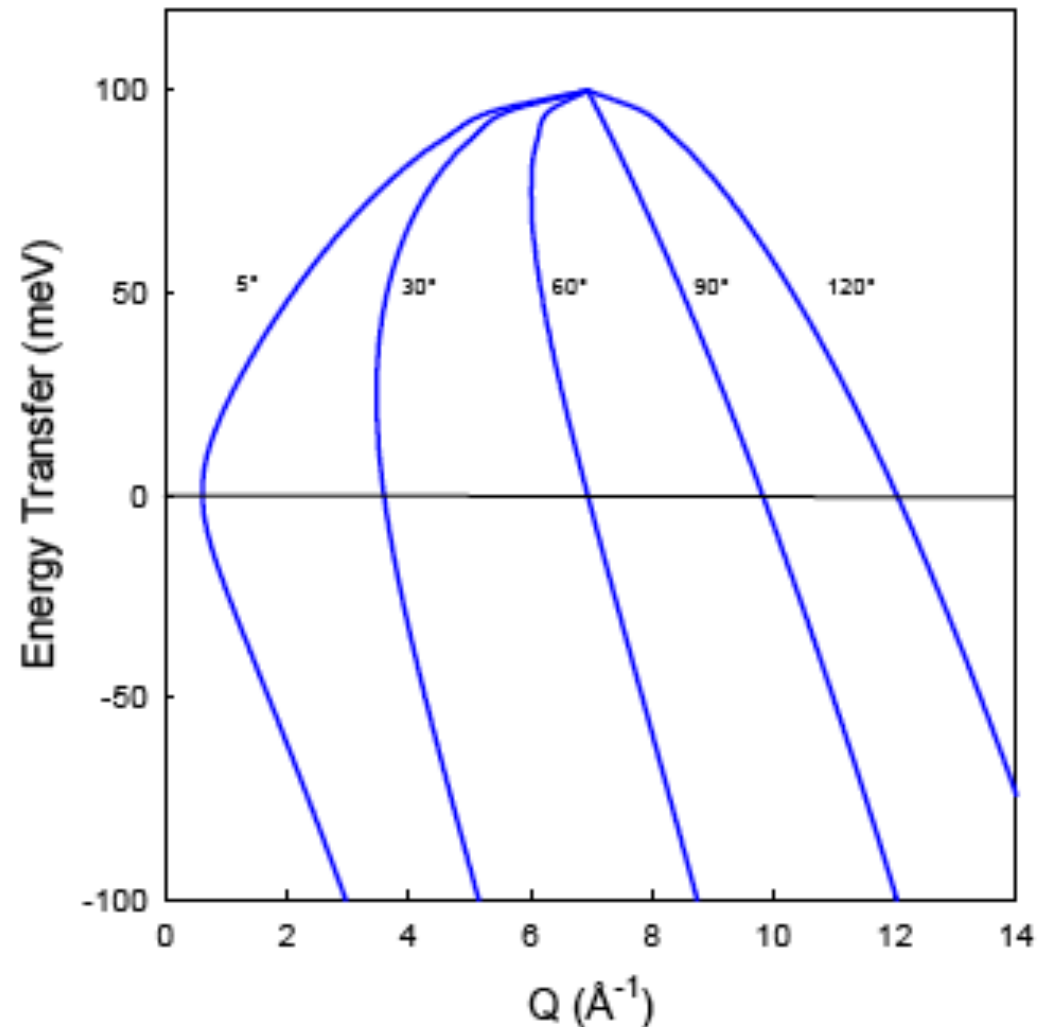
- The source operates at 50Hz.
- The neutron beam is under-moderated to preserve a high flux of epithermal neutrons and a short pulse width.



TRAJECTORIES IN (Q, ω) SPACE \rightarrow DIRECT GEOMETRY

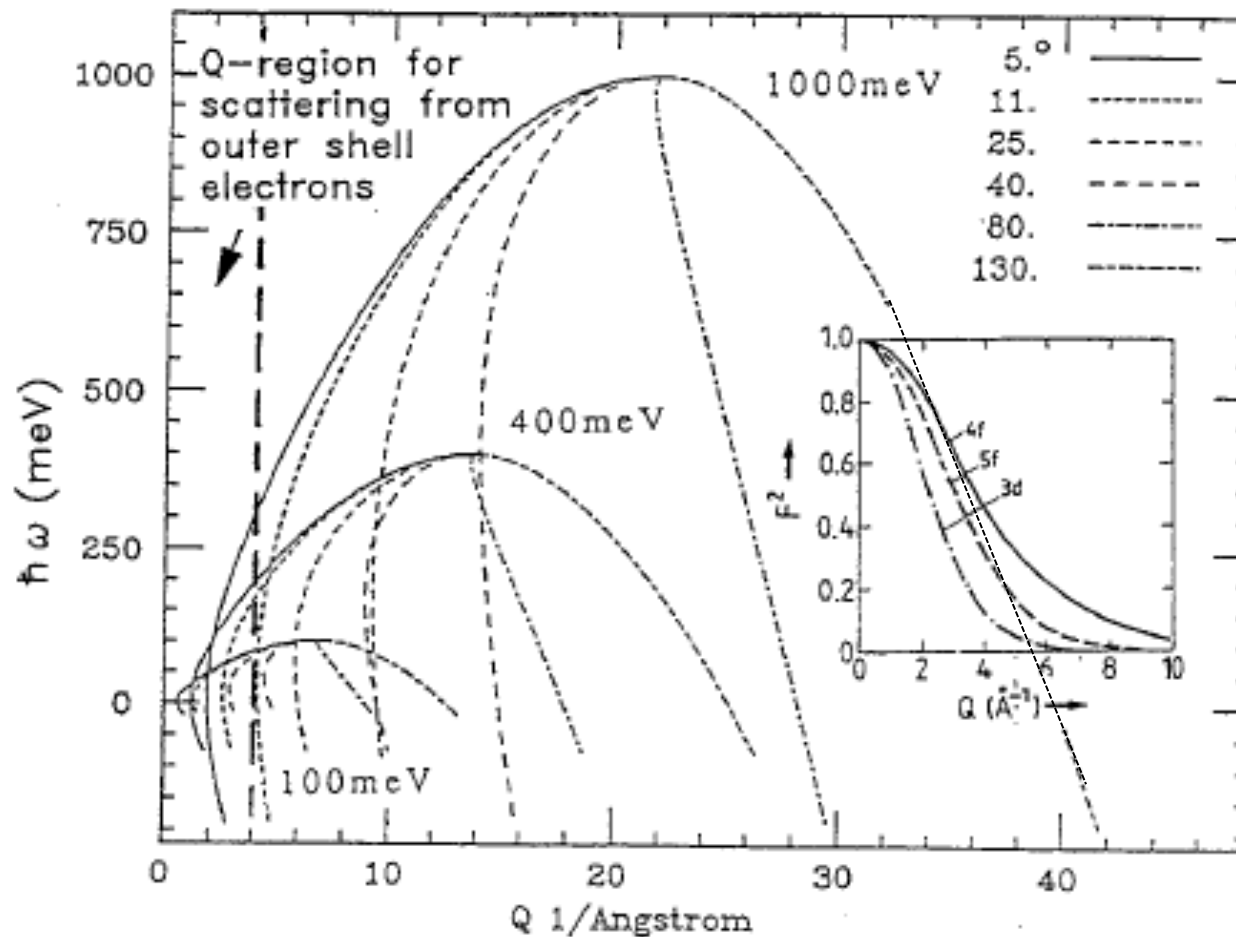
- For direct geometry eliminate E_f .
- Each detector has a parabolic trajectory through (Q, ω) space.
- To make optimum use of direct geometry spectrometers they feature large detector arrays, giving simultaneous access to a large area of (Q, ω) space.

$$\frac{\hbar^2 \mathbf{Q}^2}{2m} = 2E_i - \hbar\omega - 2 \cos \phi [E_i (E_i - \hbar\omega)]^{1/2}$$



KINEMATICS RANGE → DIRECT GEOMETRY

Kinematic range for different incident energies (from 100 to 1000 meV) and an angular detector range from 5° to 130°.



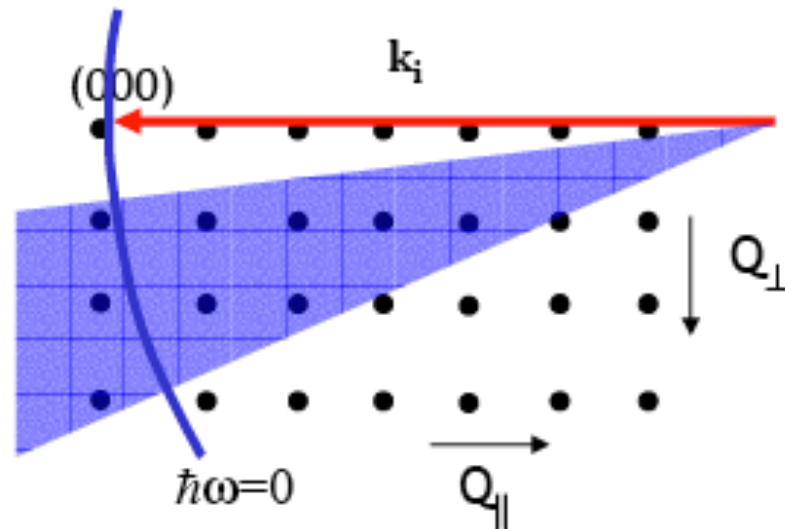
The size of the kinematic surface grows as the incident energy increases.

The insert shows the form factor of the outer shell electrons responsible for magnetic scattering

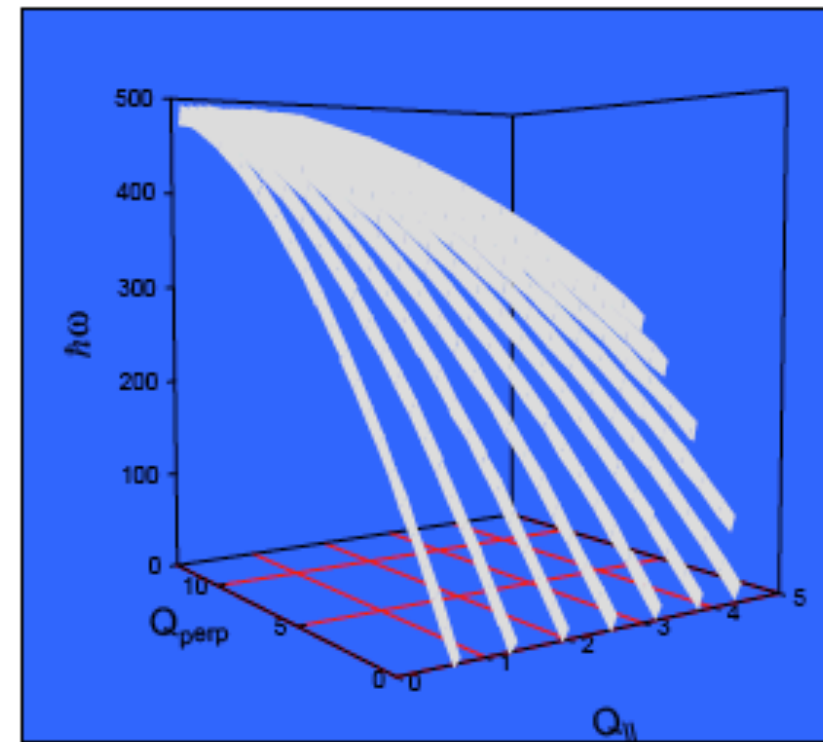
It is very difficult to measure high energy transfers at very small Q !

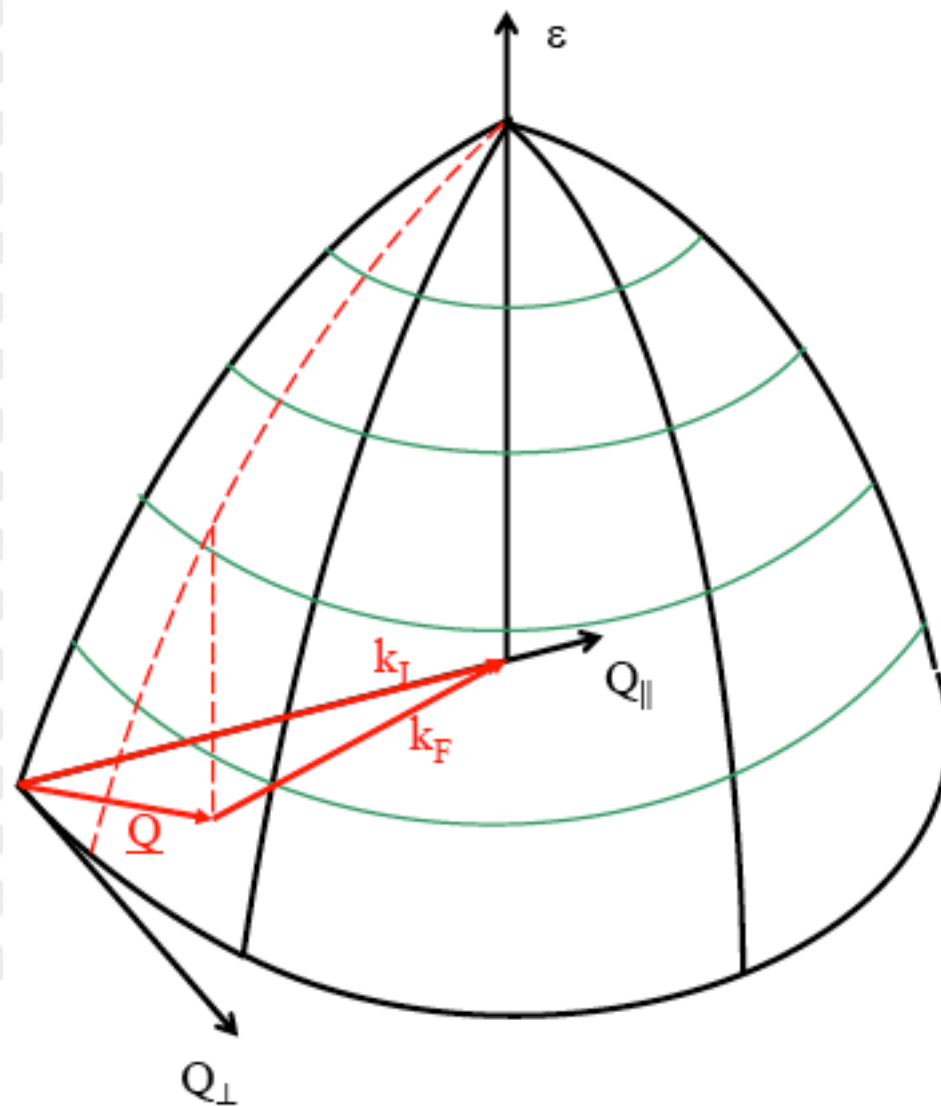
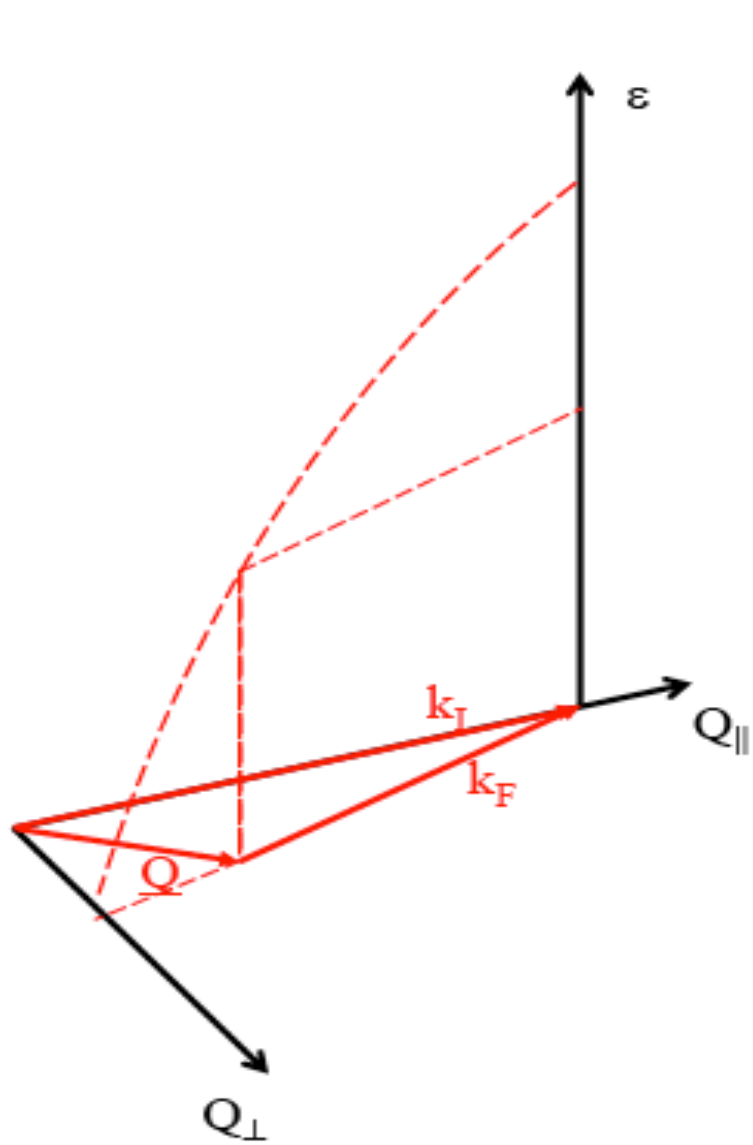
Single Crystal Experiments on a Chopper Spectrometer

- From the scattering triangle we can see that an array of detectors will trace out a sector in reciprocal space

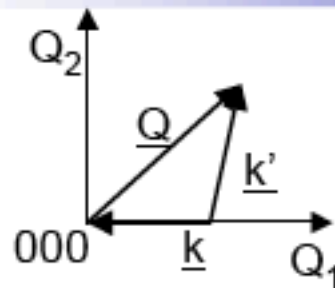


- Each detector has a parabolic trajectory through (Q, ω) space.
- The detector array produces a surface in $(Q_{\parallel}, Q_{\perp}, \omega)$ space





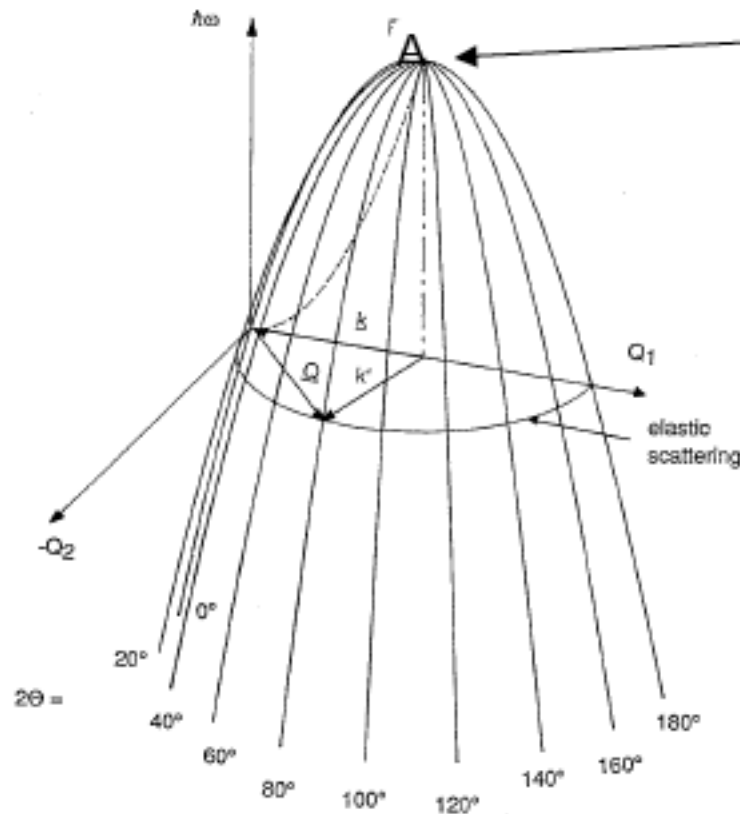
KINEMATICS RANGE → DIRECT GEOMETRY



$$Q_1 = k - k' \cdot \cos(2\theta)$$

$$Q_2 = k' \cdot \sin(2\theta)$$

$$k' = \sqrt{k^2 - \frac{2m}{\hbar^2} \cdot \hbar\omega}$$



Energy of the incident neutron:

$E = \hbar\omega = \hbar^2 k^2 / 2m \Rightarrow$ Parabola in the $\hbar\omega, k$ plane

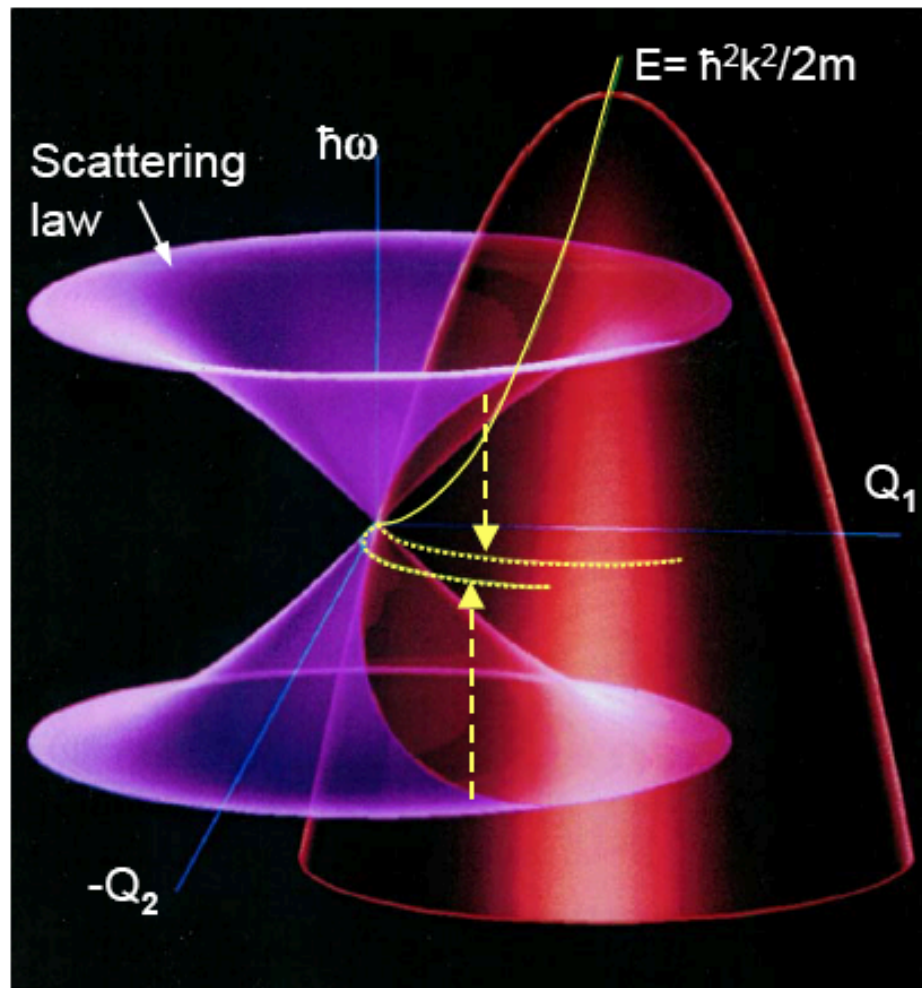
Similarly, the loci for all scattered neutrons are parabolae in the $\hbar\omega, k'$ planes with apex point A.

The paraboloid spanned by all scattering angles is therefore the locus for all possible combinations of \underline{Q} and $\hbar\omega$ that can be measured with neutrons of incident wave vector \underline{k} (kinematic scattering surface).

Elastic scattering occurs for $\hbar\omega = 0$; $\hbar\omega > 0$ means neutron energy loss, $\hbar\omega < 0$ means neutron energy gain.

KINEMATICS RANGE → DIRECT GEOMETRY

Back to 3 dimensions



In order for a neutron to be scattered its kinematic surface must intersect with the scattering law of the sample.

With a triple axis spectrometer, which does point wise scans, it is possible to follow the scattering law along symmetry directions in the reciprocal crystal lattice.

In a multidetector time of flight scan with fixed incident neutron energy the loci for the Q-vectors measured are curved. The scattering law along symmetry directions must be constructed from many scans at different orientations of the sample.

ENERGY RESOLUTION OF A CHOPPER SPECTROMETER

Energy Resolution

Moderator term

$$\Delta t_{md} = \frac{L_2 + L_3}{L_1} \Delta t_{mm}$$

Chopper term

$$\Delta t_{cd} = \frac{L_2 + L_1 + L_3}{L_1} \Delta t_{cc}$$

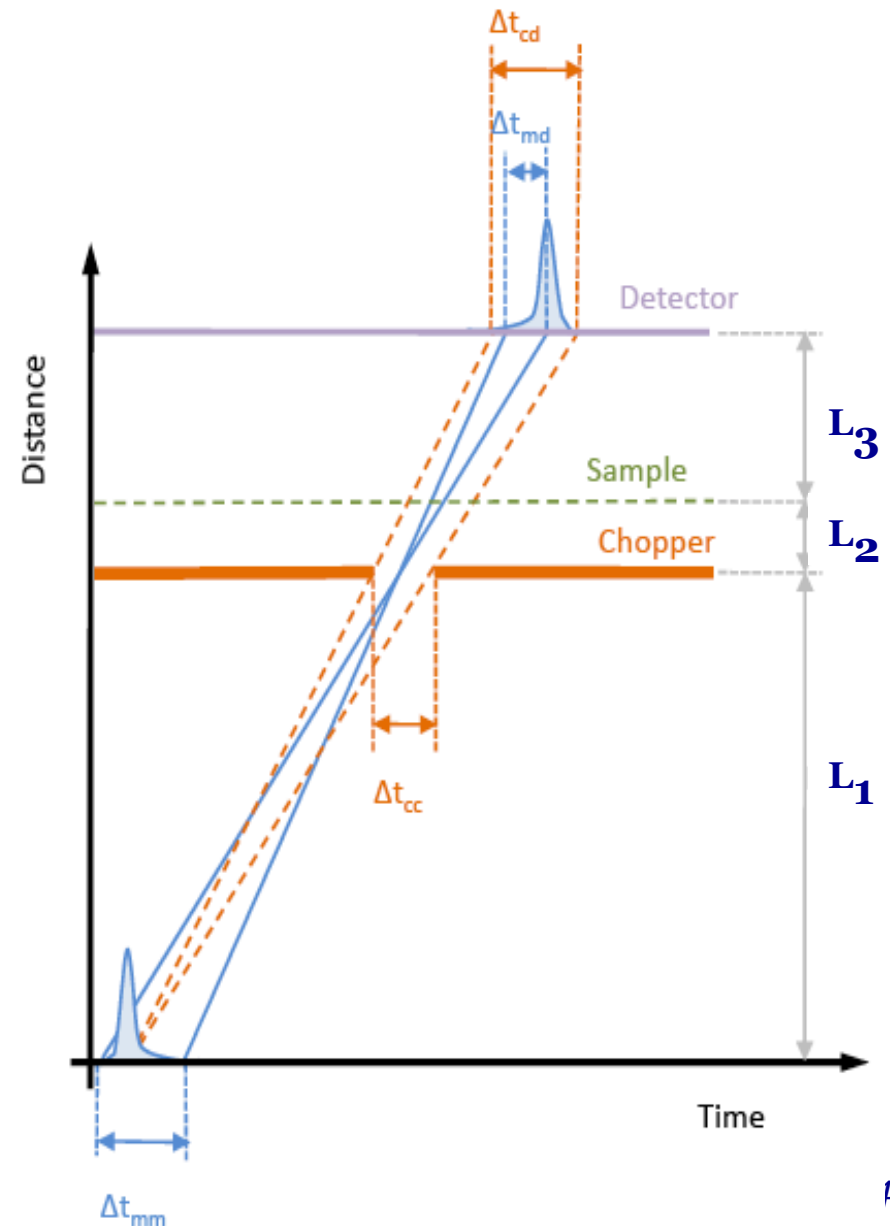
These are added in quadrature and converted to energy

$$\frac{\Delta \varepsilon}{E_i} = 8.7478 \times 10^{-4} \frac{\sqrt{E_i (meV)}}{L_2 (m)} \Delta t (\mu s)$$

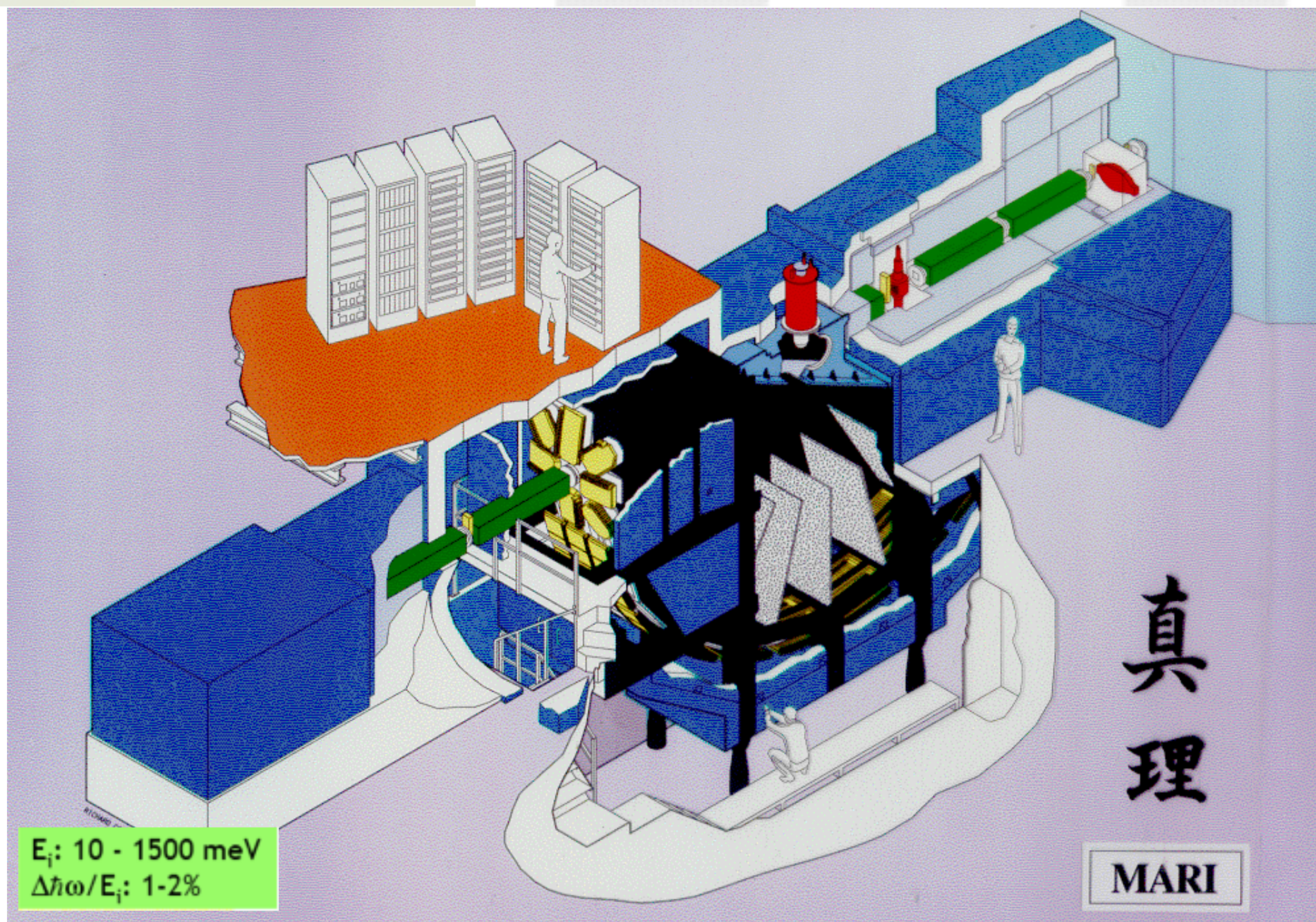
Resolution increases as L_2 lengthens

Flux

$$\Phi = \left(\frac{p}{d} \right) \left(\frac{WmHm}{(L_1 + L_3)^2} \right) \left(\frac{\Delta t_{cc}}{L_1} \right)$$

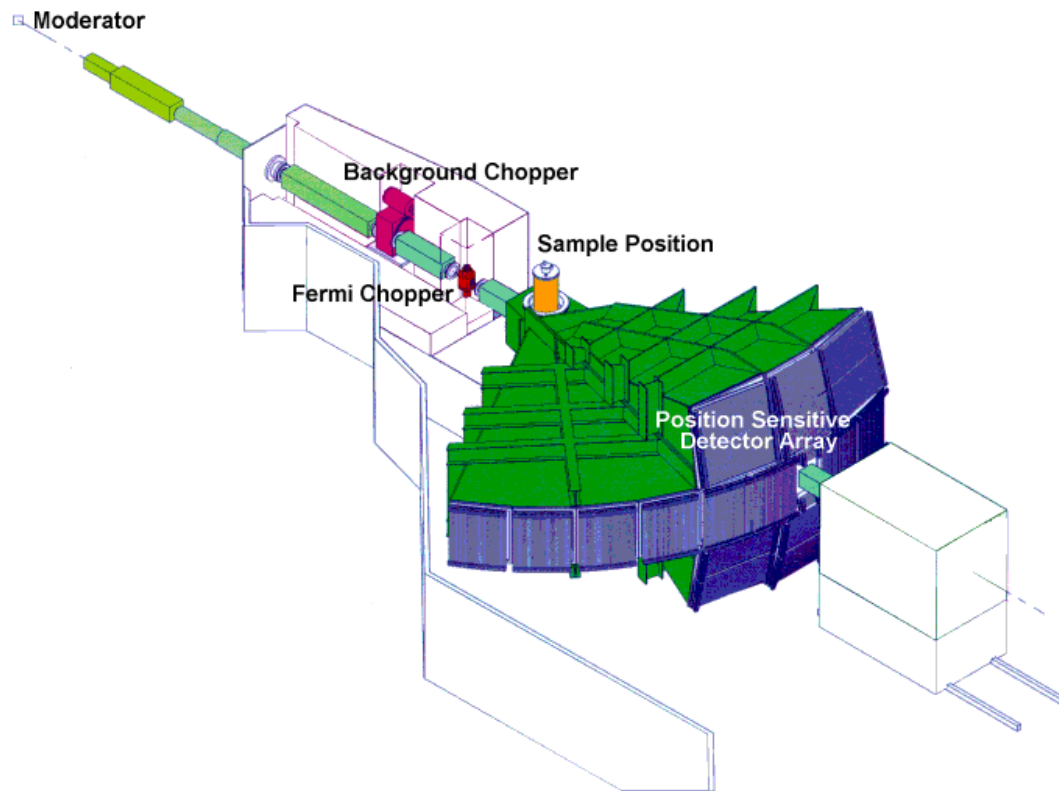


ENERGY RESOLUTION OF A CHOPPER SPECTROMETER



MAPS

- 16 m² of PSDs
- 600 1m detectors
- 70 000 pixels

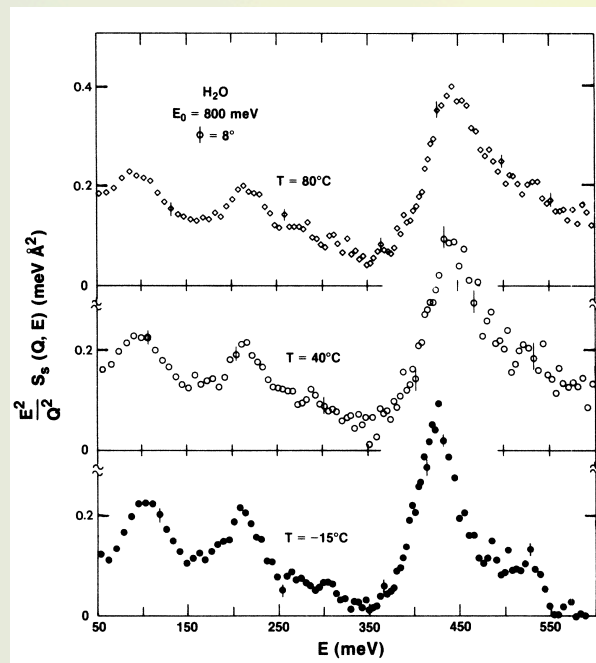


EXPERIMENTS ON DIRECT GEOMETRY SPECTROMETER

- Direct geometry spectrometers offer simultaneous coverage of a wide area of (Q, ω) space with a wide range in incident energies (IN6~ 2.3 meV, HET~ 2eV)
 - **Vibrational motions in molecular systems**
 - Magnetic excitations in polycrystalline samples
 - Amorphous materials
 - Quantum fluids
 - Excitations in single crystal
 - Momentum distributions

Inelastic Neutron Scattering: a probe for H vibrations

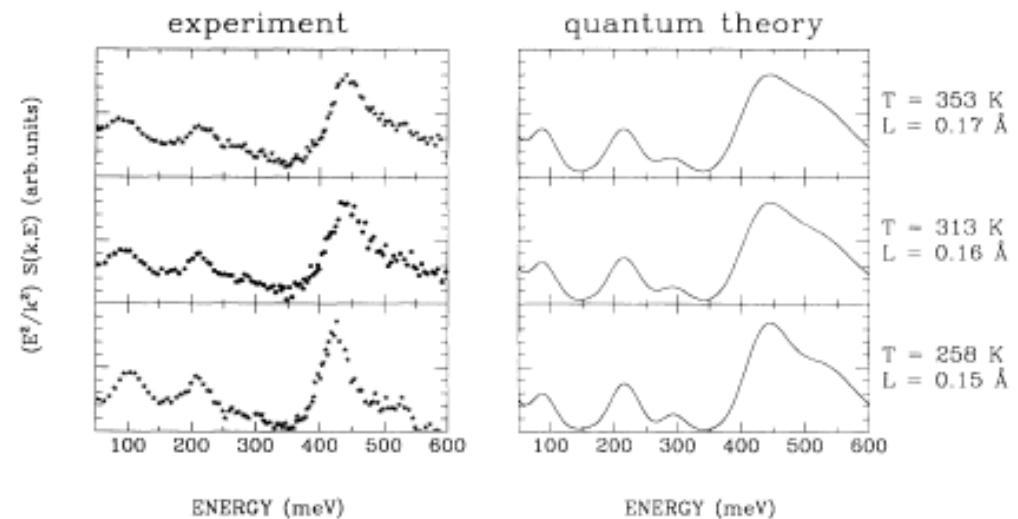
Inelastic Incoherent Neutron Scattering measurements of water at IPNS



Toukan et al., PRB (1988)

Incoherent inelastic neutron scattering from liquid water: A theoretical investigation

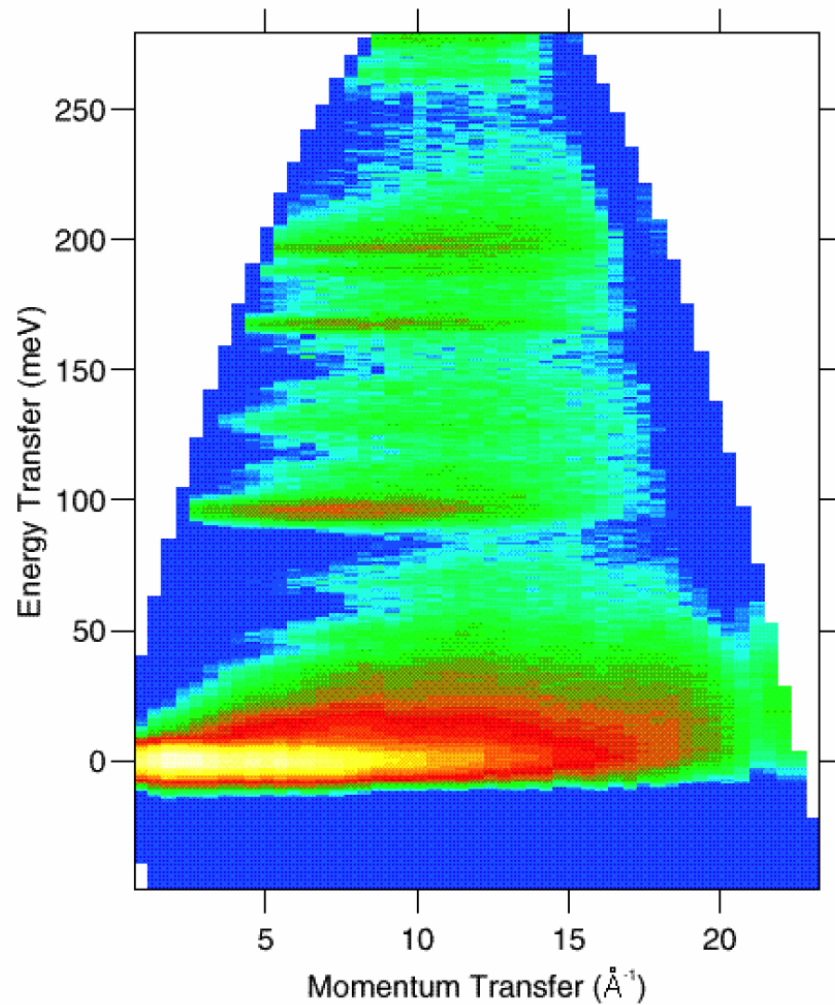
S. Bratos, M. Diraison, G. Tarjus, and J-Cl. Leicknam



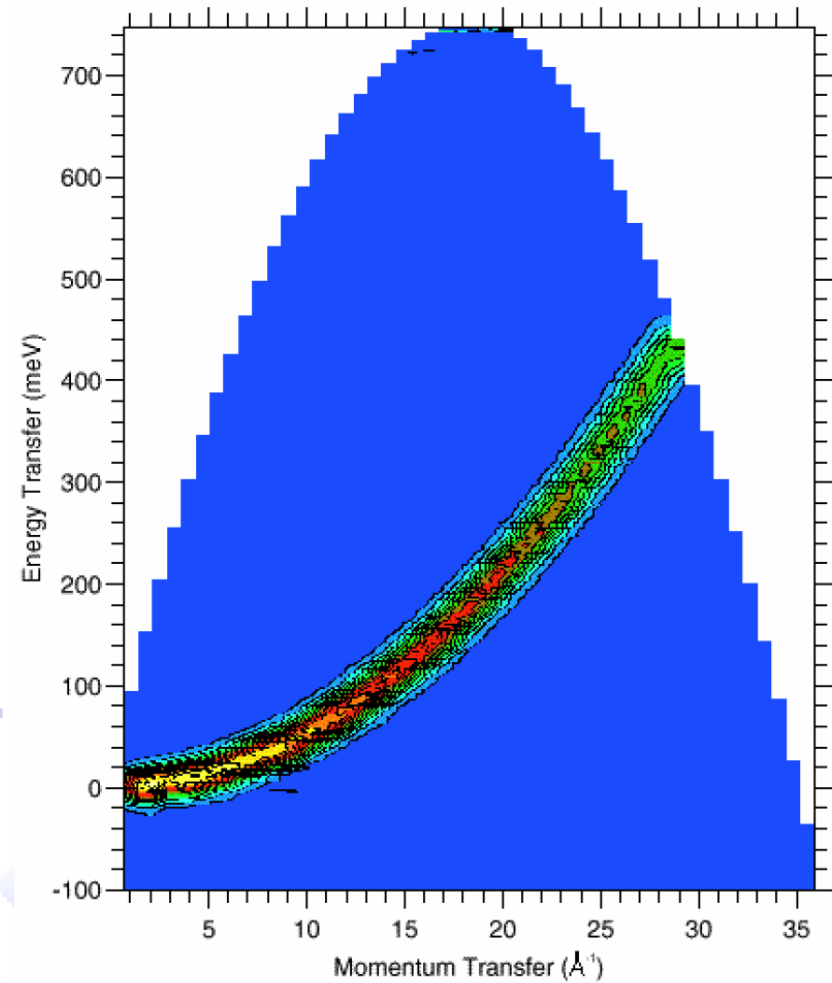
Bratos et al., PRA (1992)

@ HET → EXAMPLE OF EXPERIMENT

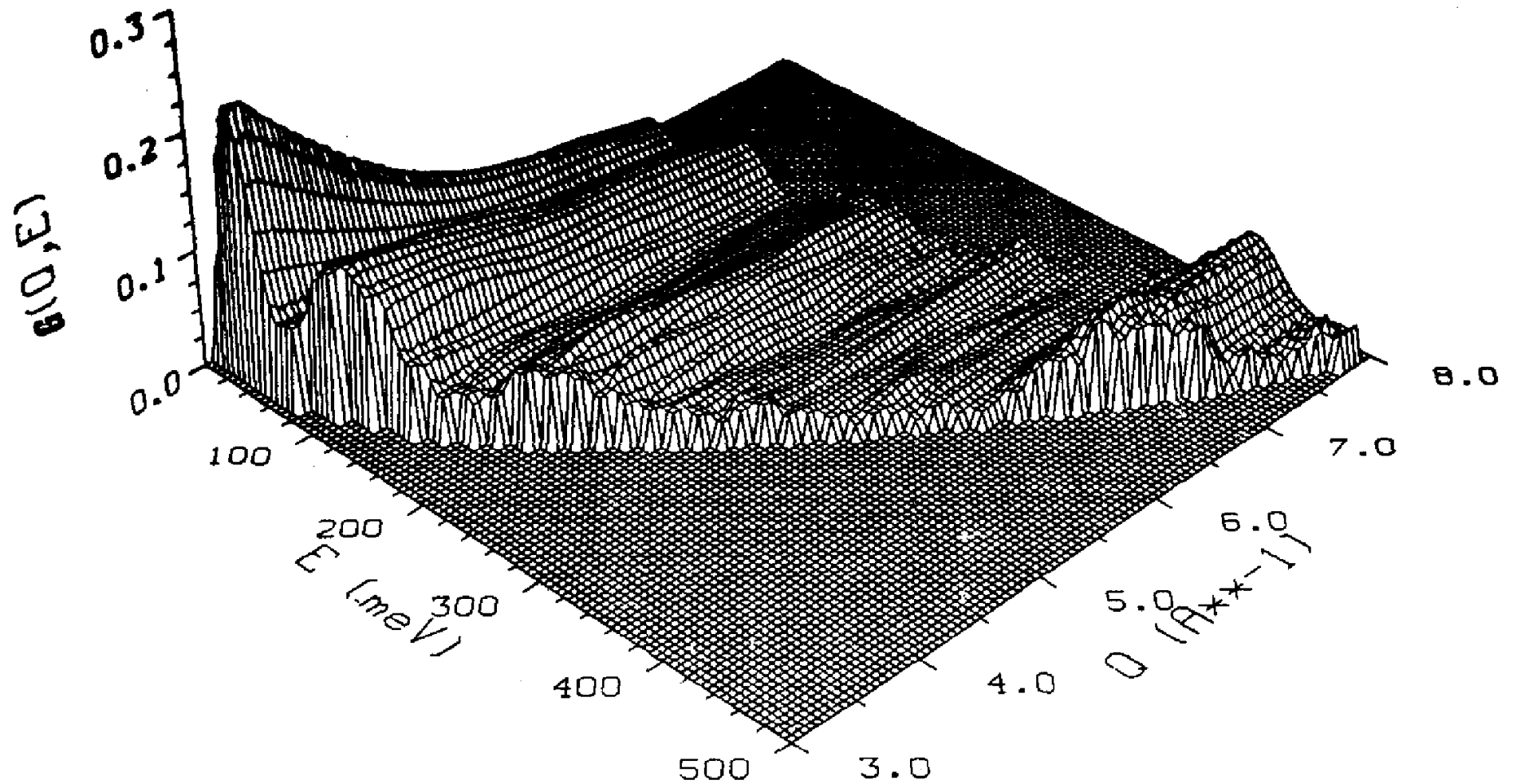
n-methyl acetamide (Fillaux)



^4He (Stirling et al)

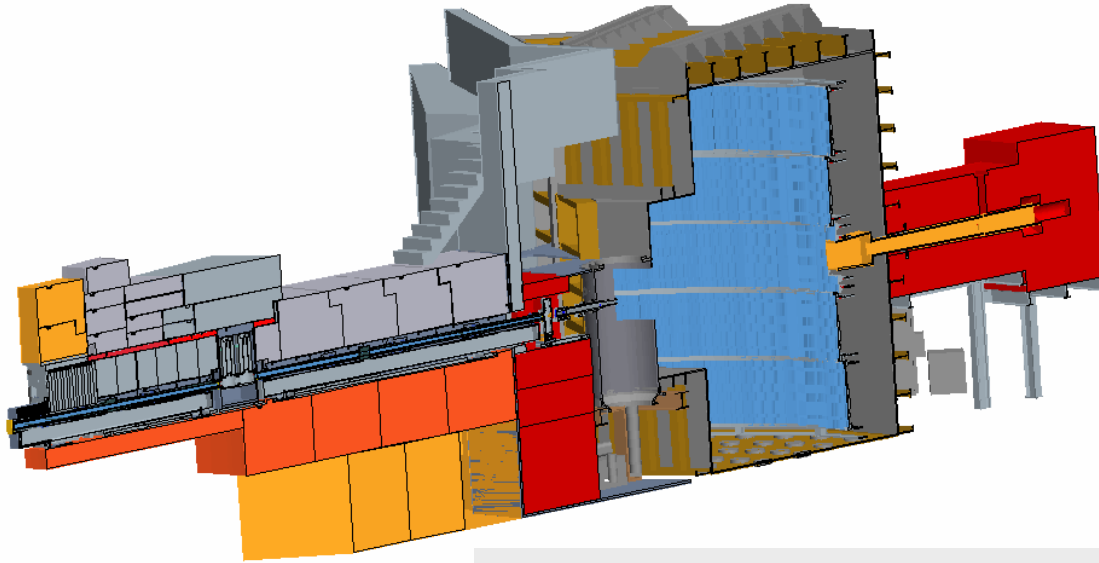


$S(q, \omega)$ for ice Ih, @ IPNS



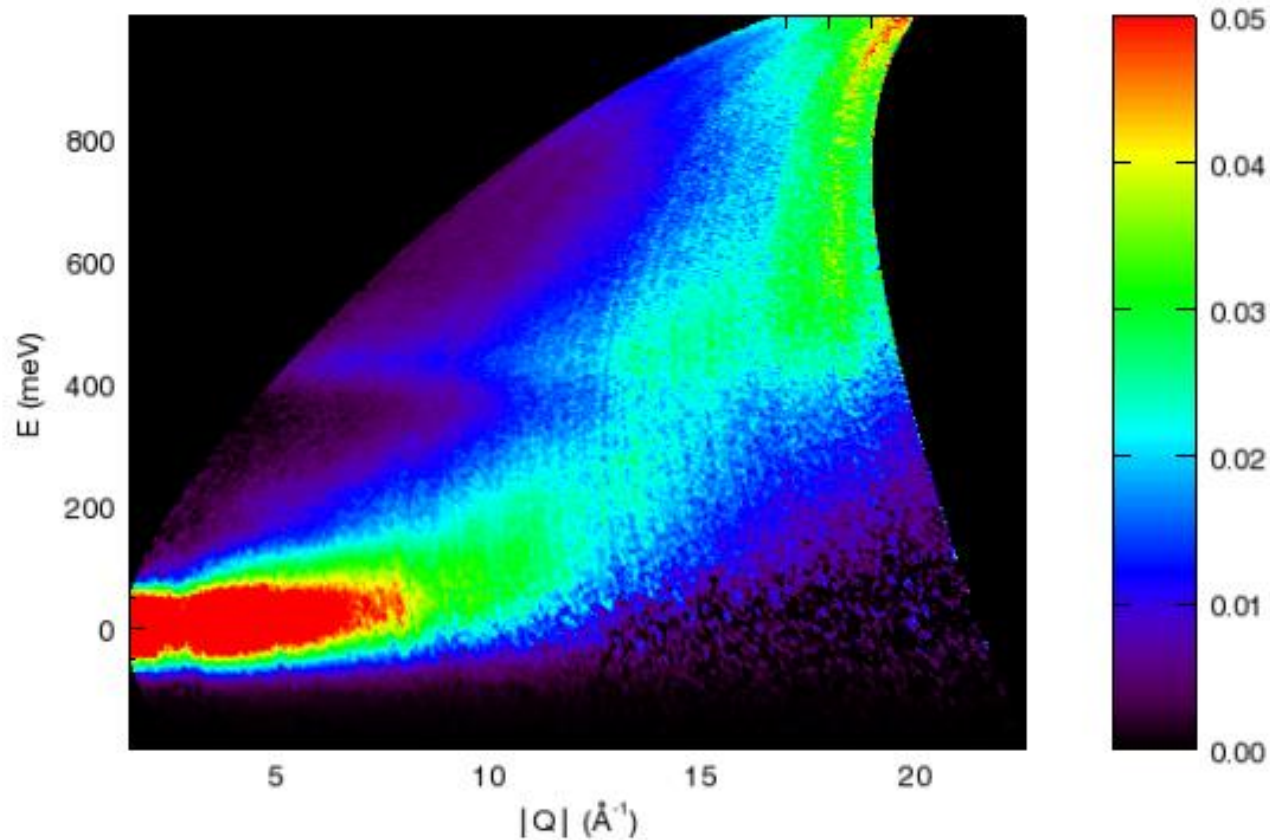
SEQUOIA @ SNS

<http://neutrons.ornl.gov/sequoia/>

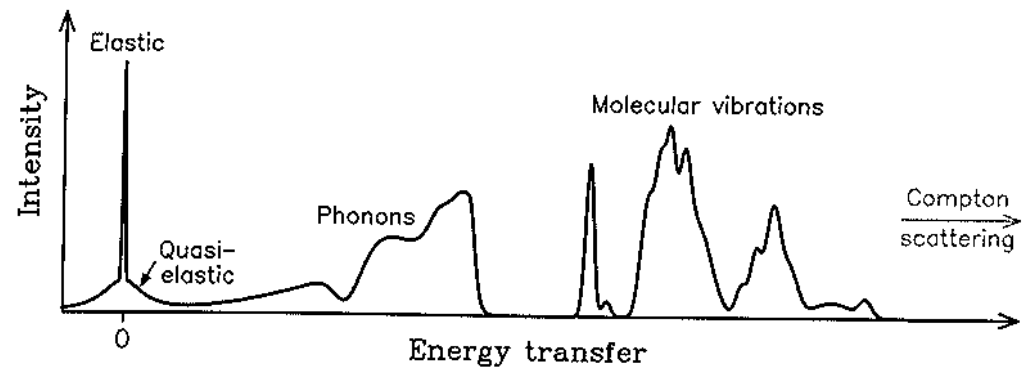
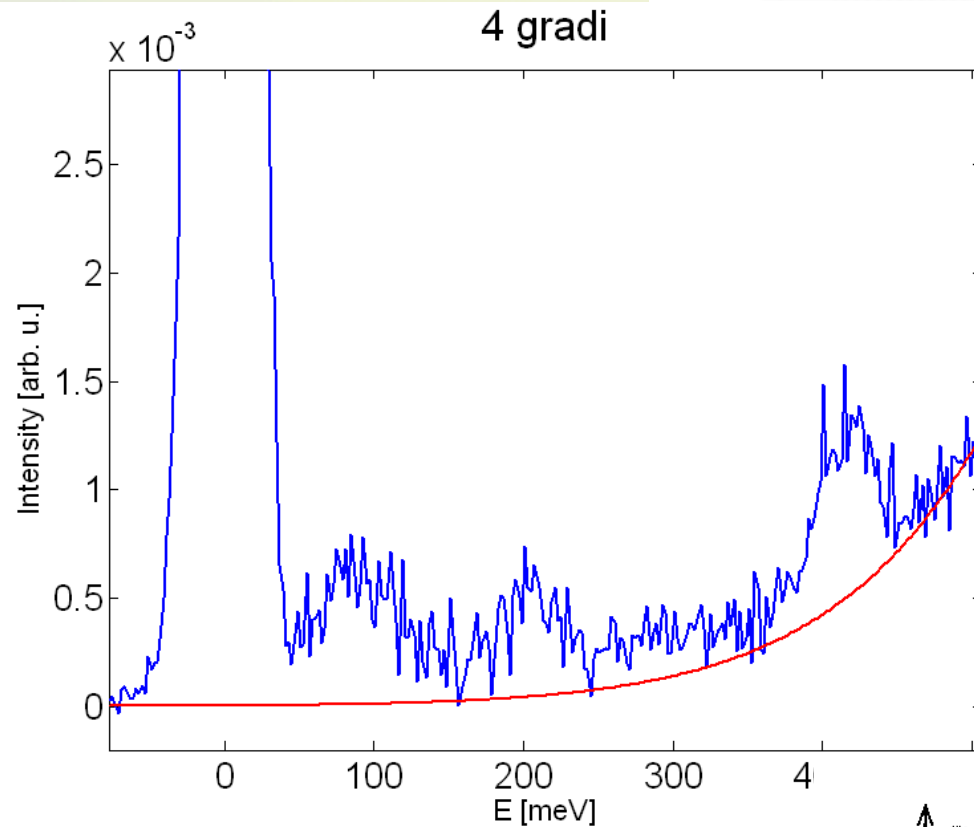


Incident energy	
Range	10–2000 meV
Resolution (elastic)	1–5% E_i
Vertical	
Detector coverage	-18–18°
Horizontal	
Detector coverage	-30–60°
Minimum detector Angle	3°

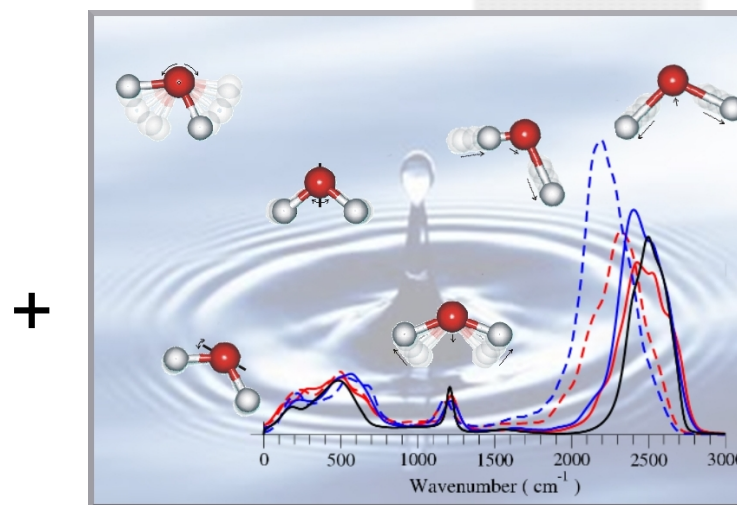
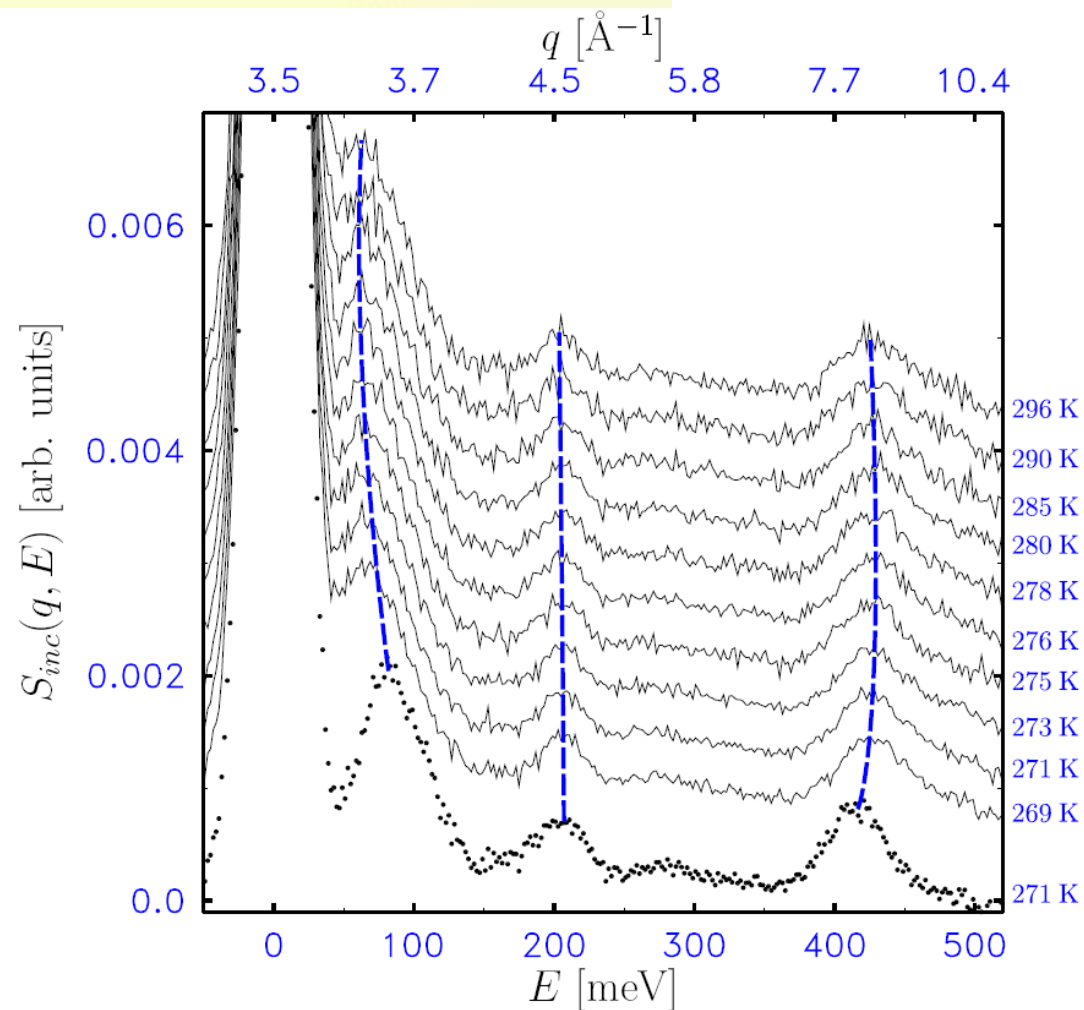
- Put a sample in the beam (ice and/or water)
- Record intensity maps as a function of energy-momentum
- Transfers using data reduction (softwares)



- data correction: multiple scattering, multiphonon, etc, to obtain $S(Q,\omega)$ [from $S(\Theta,\omega)$]



INS and *ab initio* electronic structure calculations on Results from Inelastic neutron scattering on water across the triple point

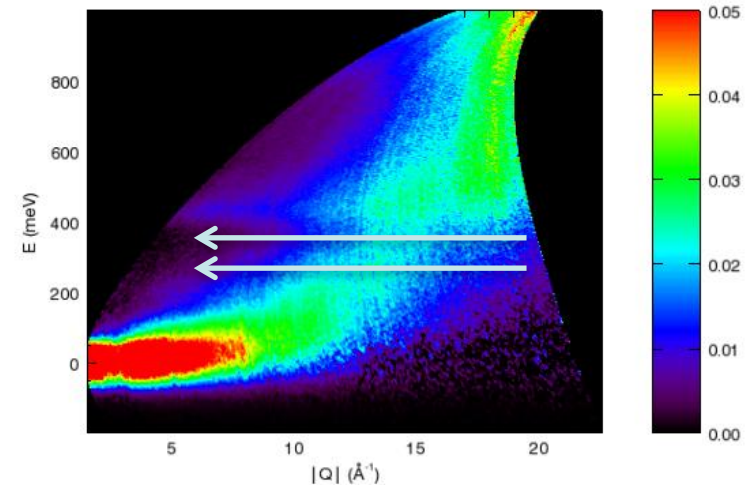
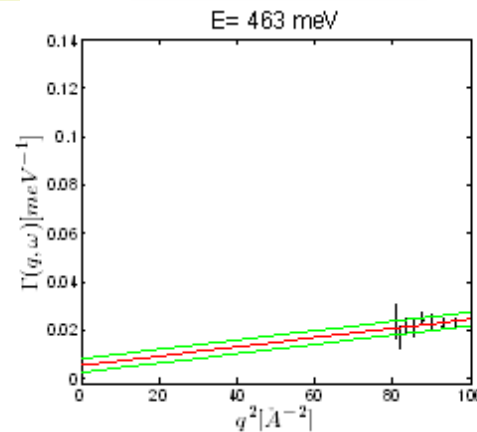
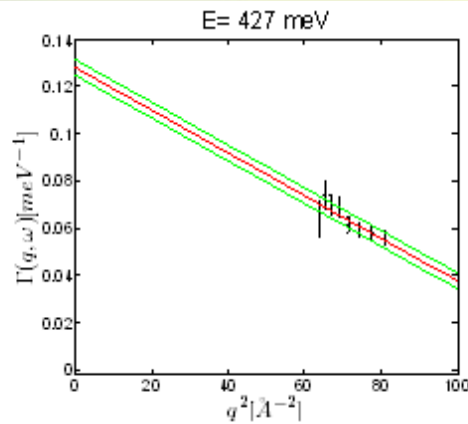
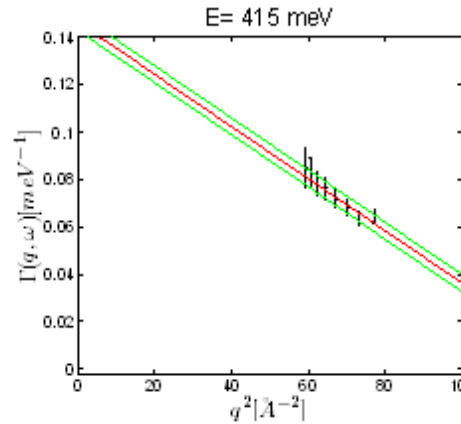
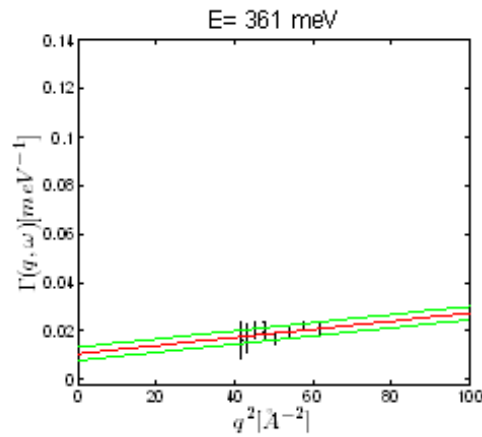


<http://angstrom.ucdavis.edu/>

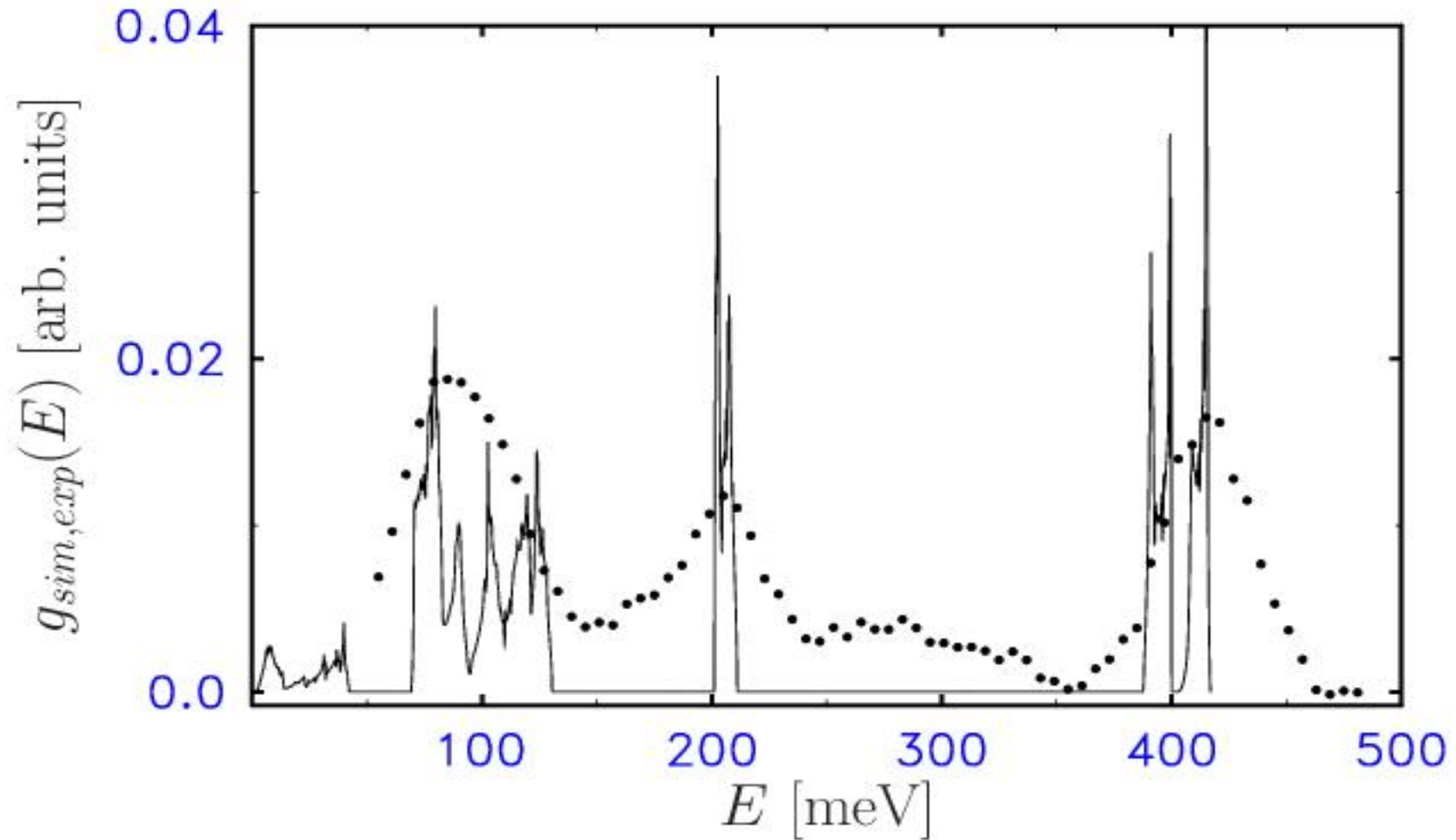
Make a constant ω cut towards $Q=0$!
 i.e. for each ω within the range of
 the stretching band take
 the zero- Q limit by extrapolation.

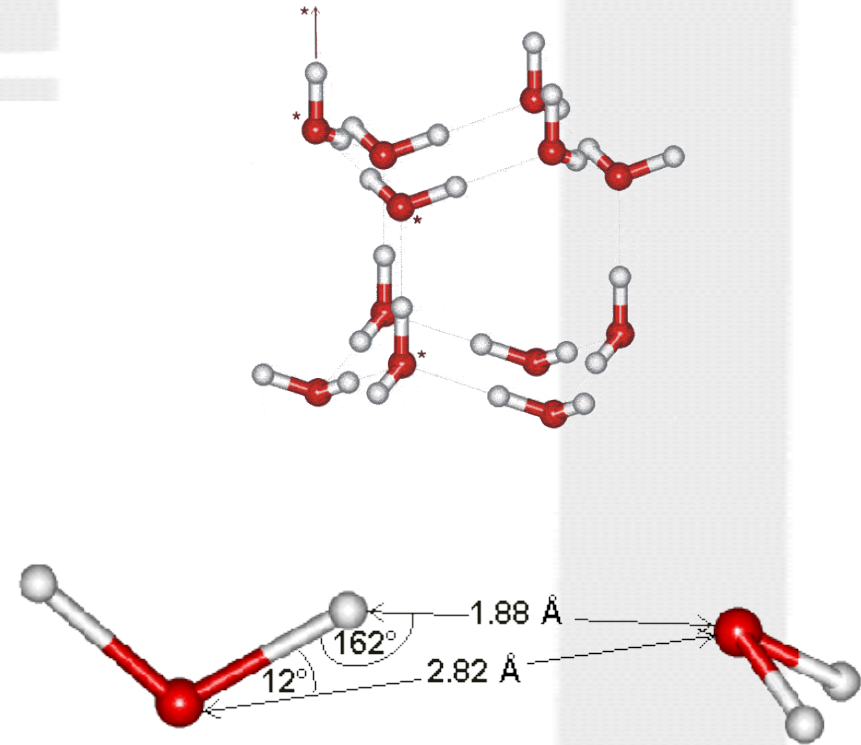
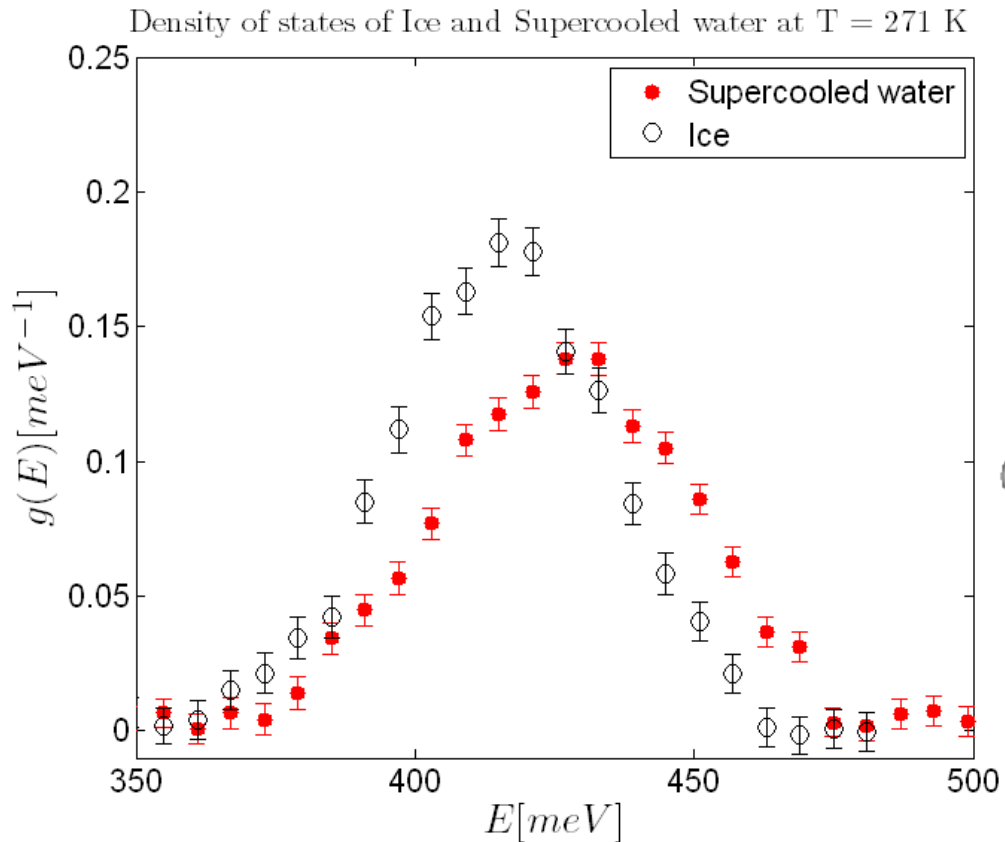
$$\lim_{Q \rightarrow 0} \left[\frac{S_{inc}(Q, \omega)}{Q^2} \frac{2m \hbar \omega}{e^{-2W}} \frac{4\pi}{\sigma_{inc}} \frac{1}{n(\omega) + 1} \right] = \lim_{Q \rightarrow 0} \Gamma(Q, \omega)$$

$$= g(\omega) \text{ or } g(E)$$



$g(E)$: from *ab initio* electronic structure calculations and from INS





Red-shifted in ice due greater association of the proton with the oxygen accepting the hydrogen bond, thereby resulting in the weakening of the covalent bond...



THANK YOU

QUESTIONS?