

Detectors for neutron scattering experiments

M. Tardocchi

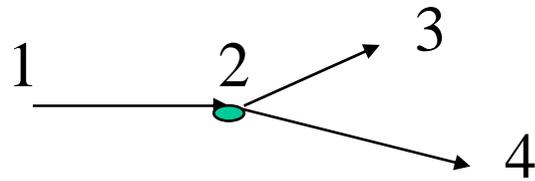
tardocchi@ifp.cnr.it

*Istituto di Fisica del Plasma, EURATOM-ENEA-CNR
Association, Milan, Italy*

Nomenclature

Category	Energy [meV]	Temperature [K]	λ [Å]
Ultra-cold	< 0.1	< 1	< 30
Cold	0.1 – 10	1 – 120	30 – 3
Thermal	10– 25-100	120 – 1000	3 – 1
Epithermal	> 500	> 6000	> 0.4
Fast	> ~ 10⁵	> 10⁶	>0.03

Schematics of nuclear reactions



Energy conservation $M_1+M_2 +T_1= M_3+M_4 +T_3+T_4$ (1)

1 **Elastic collision** (rest) mass and T_{kin} are conserved; $M_1=M_3$; $M_2=M_4$

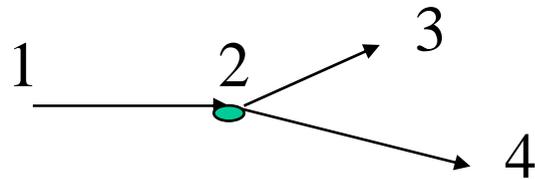
2 **Anelastic collision*** mass is conserved, T_{kin} is not conserved

3 **Nuclear reactions** mass and T_{kin} are not conserved

$Q= (M_1+M_2) - (M_3+M_4)$ (2) **Q-value definition**

• Condensed matter scientist use the term **anelastic** to mean reaction of type 1 when $E_n \text{ initial} \neq E_n \text{ final}$ (exchange of $E = \hbar\omega$ and $Q = \hbar q$)

Schematics of nuclear reactions



Energy conservation $M_1+M_2 +T_1= M_3+M_4 +T_3+T_4$ (1)

eq.1 + eq.2

$$(T_3+ T_4) - T_1= Q$$

$$M_1= M_3, M_2= M_4 \text{ and } Q =0$$

Elastic reactions

$$M_1= M_3, M_2 \rightarrow M_4^* \text{ and } Q \neq 0$$

(M_4 is left in an excited state)

Anelastic reactions

$Q < 0$ Endhotermic Mass is created from kinetic energy

$Q > 0$ Esothermic Mass is trasnformed in kinetic energy

$$M_1 \neq M_3 \text{ or } M_2 \neq M_4 \text{ and } Q \neq 0$$

Nuclear reactions ←

Used for detection of thermal neutrons

Reactions for slow neutron detection

- $n + {}^{10}\text{B}$ (a.i. 20%) $\rightarrow {}^7\text{Li}^* + \alpha \rightarrow {}^7\text{Li} + \alpha + 0.48 \text{ MeV } \gamma + 2.3 \text{ MeV}$ (93%) 3840 barns
 $\rightarrow {}^7\text{Li} + {}^4\text{He} + 2.8 \text{ MeV}$ (7%)
- $n + {}^3\text{He} \rightarrow {}^3\text{H} + {}^1\text{H} + 0.764 \text{ MeV}$ (expensive and scarce resource) **5330 barns**
- $n + {}^6\text{Li}$ (a.i. 7%) $\rightarrow {}^4\text{He} + {}^3\text{H} + 4.79 \text{ MeV}$ (resonance at 100keV) 940 barns

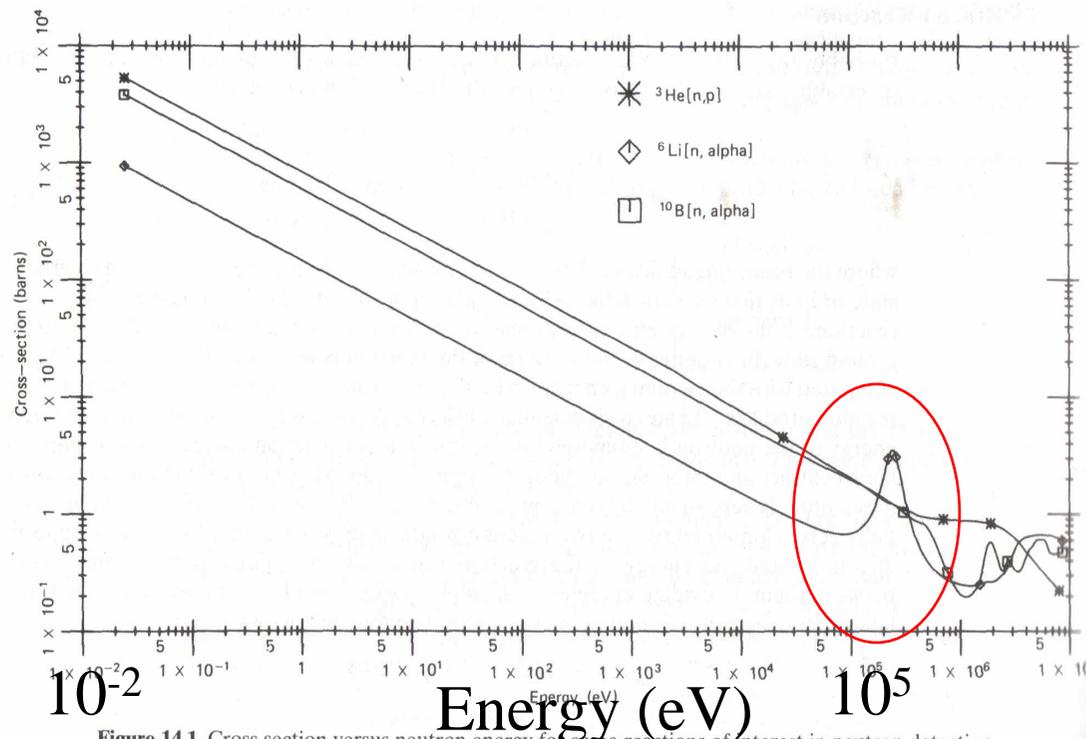


Figure 14.1 Cross section versus neutron energy for some reactions of interest in neutron detection.

Detectors for slow neutrons

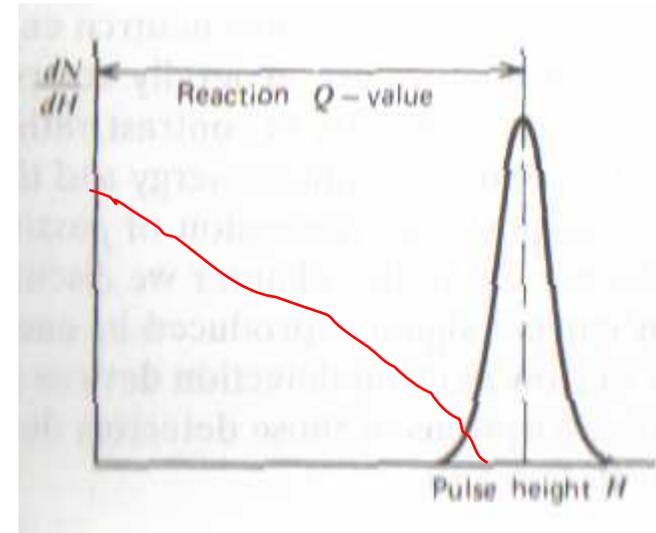
Ideal detector:

High detection efficiency (cross section)

Large Q values (only $Q > 0$)

Stop reaction products

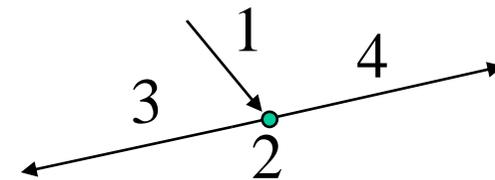
Immune to background (often γ rays)



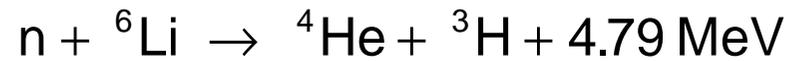
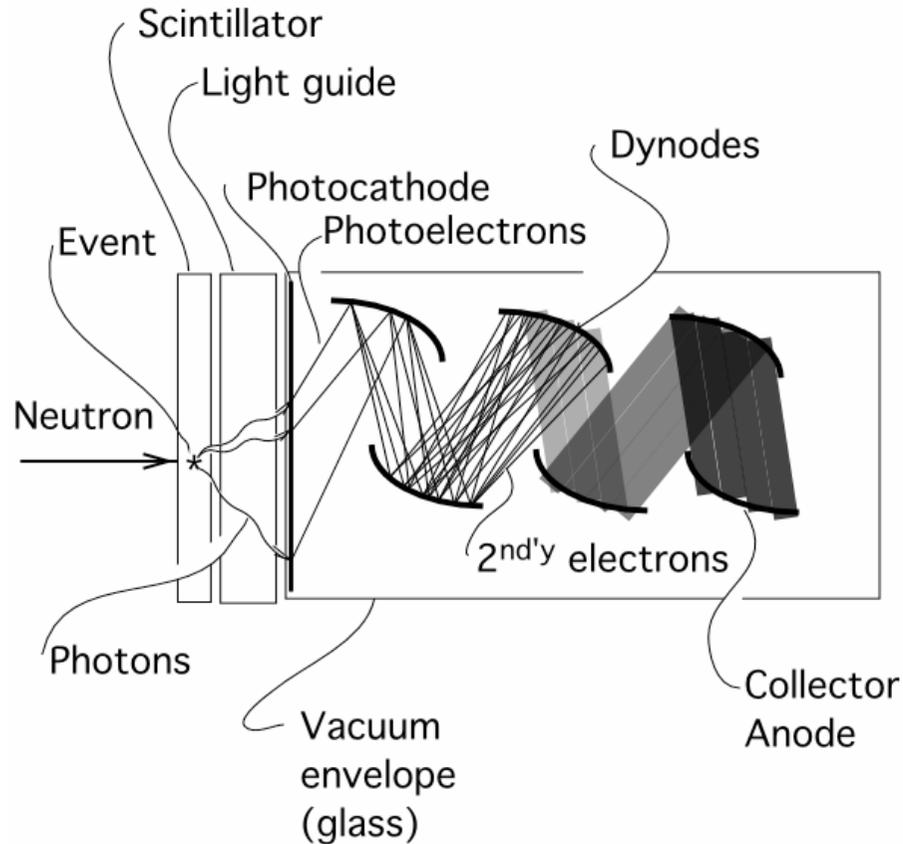
$$E_{\text{kinetics}} \text{ of the products} = Q + \cancel{E_n} = Q$$

Products emitted back-to-back ($P_{\text{init}} \sim 0$) in the lab. ref. system

$$\begin{cases} E_3 + E_4 = Q \\ m_3 v_3 + m_4 v_4 = 0 \end{cases} \quad \dots \quad E_{3,4} = \frac{m_{4,3}}{m_3 + m_4} \cdot Q$$



Scintillation Detectors



$$\sigma = 940 \frac{\lambda}{1.8} \text{ barns}$$

Some Common Scintillators for Neutron Detectors

Intrinsic scintillators contain small concentrations of ions ("wave shifters") that shift the wavelength of the originally emitted light to the longer wavelength region easily sensed by photomultipliers.

ZnS(Ag) is the brightest scintillator known, an intrinsic scintillator that is mixed heterogeneously with converter material, usually Li^6F in the "Stedman" recipe, to form scintillating composites. These are only semitransparent. But it is somewhat slow, decaying with $\sim 10 \mu\text{sec}$ halftime.

GS-20 (glass, Ce^{3+}) is mixed with a high concentration of Li_2O in the melt to form a material transparent to light.

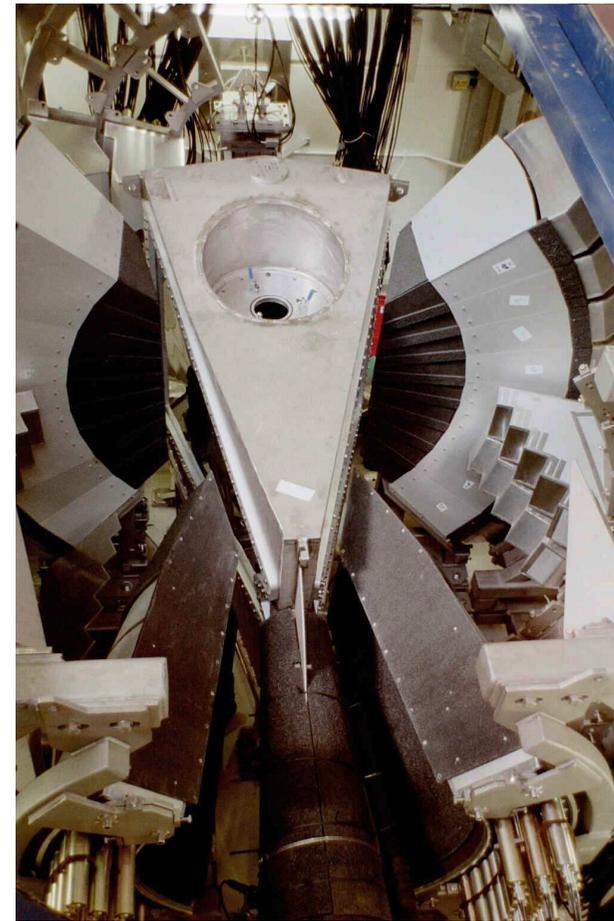
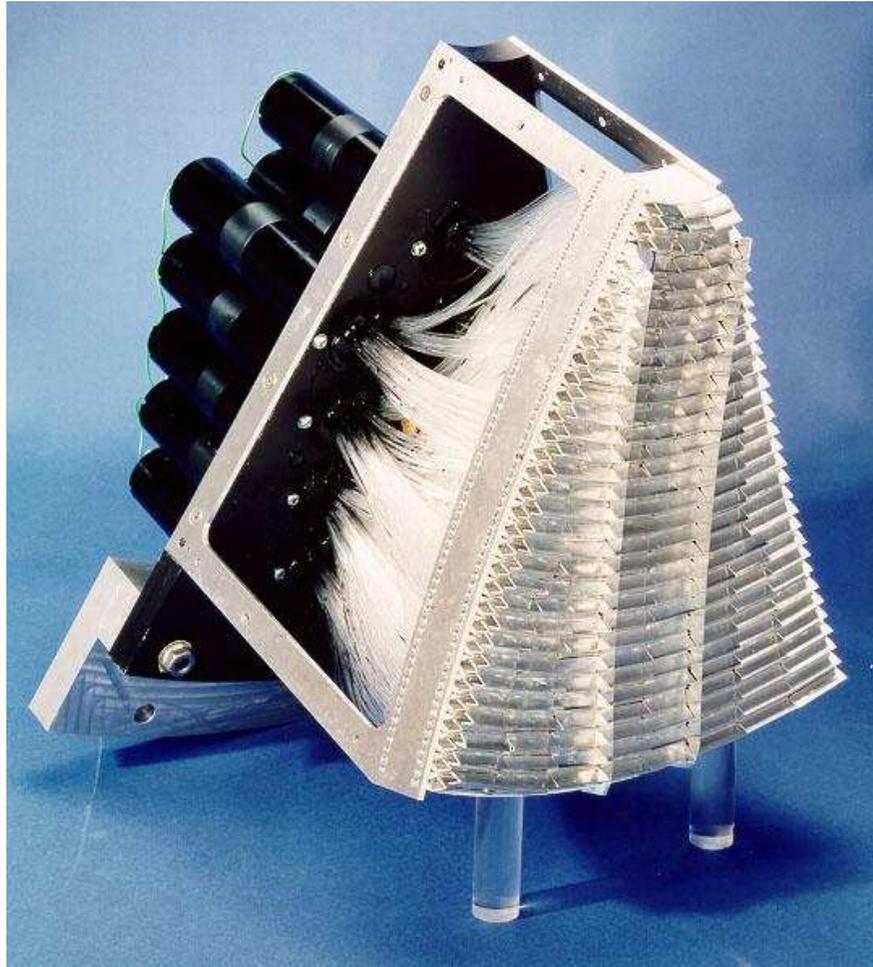
$\text{Li}_6\text{Gd}(\text{BO}_3)_3 (\text{Ce}^{3+})$ (including ^{158}Gd and ^{160}Gd , ^6Li , and ^{11}B), and $^6\text{LiF}(\text{Eu})$ are intrinsic scintillators that contain high proportions of converter material and are typically transparent.

An efficient gamma ray detector with little sensitivity to neutrons, used in conjunction with neutron capture gamma-ray converters, is **YAP** (yttrium aluminum perovskite, $\text{YAl}_2\text{O}_3(\text{Ce}^{3+})$).

Some Common Scintillators for Neutron Detectors

Material	Density of ${}^6\text{Li}$ atoms (cm^{-3})	Photon wavelength (nm)	Photons per neutron
Li glass (Ce)	1.75×10^{22}	395 nm	~7,000
LiI (Eu)	1.83×10^{22}	470	~51,000
ZnS (Ag) - LiF	1.18×10^{22}	450	~160,000
$\text{Li}_6\text{Gd}(\text{BO}_3)_3$ (Ce),	3.3×10^{22}	~ 400	~40,000
YAP	--	350	~18,000 per MeV gamma

GEM Detector Module (at ISIS)



Hamamatsu Multicathode Photomultiplier

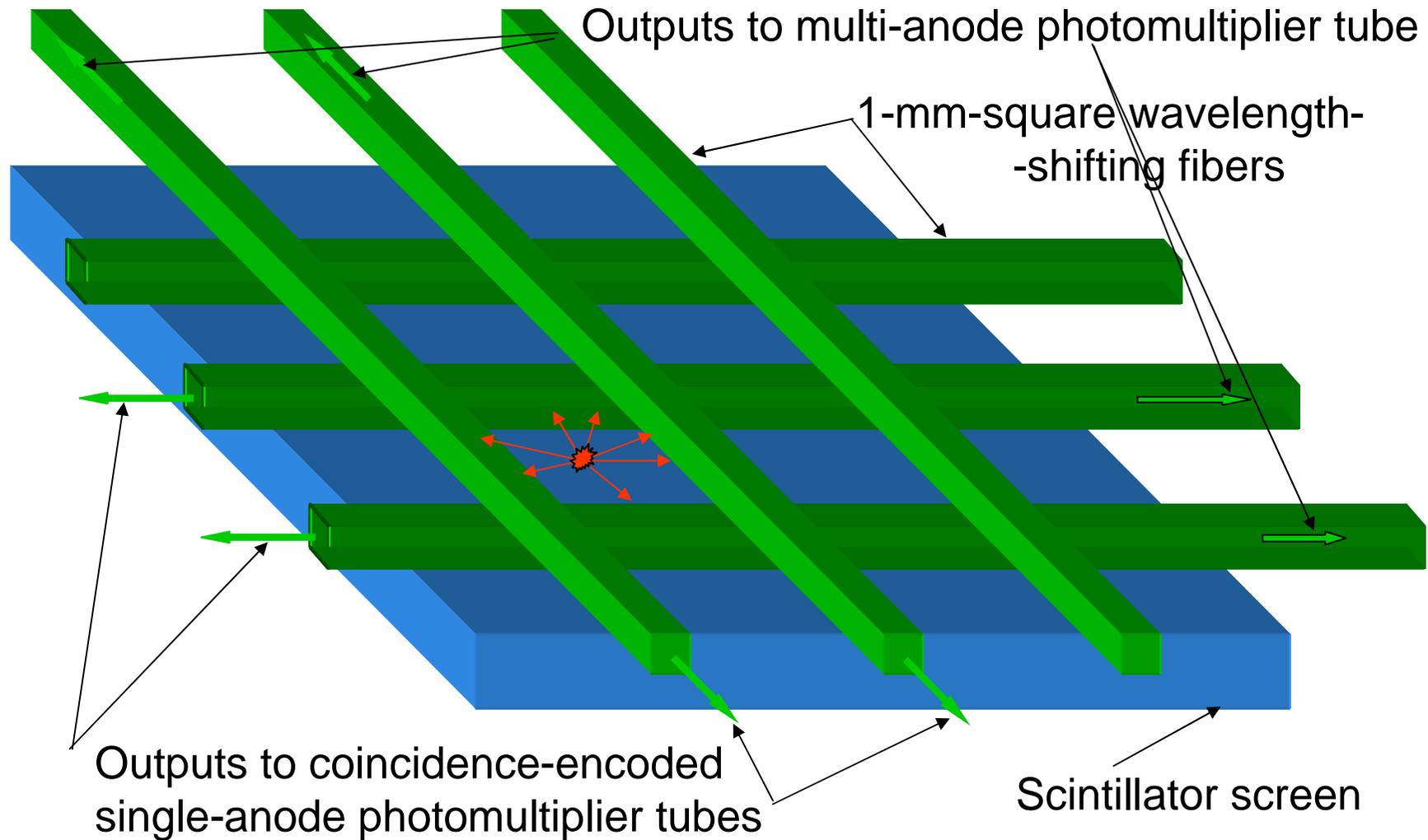
Compact photomultipliers are essential components of scintillation area detectors. The figure shows a recently developed multicathode photomultiplier, Hamamatsu model 8500.



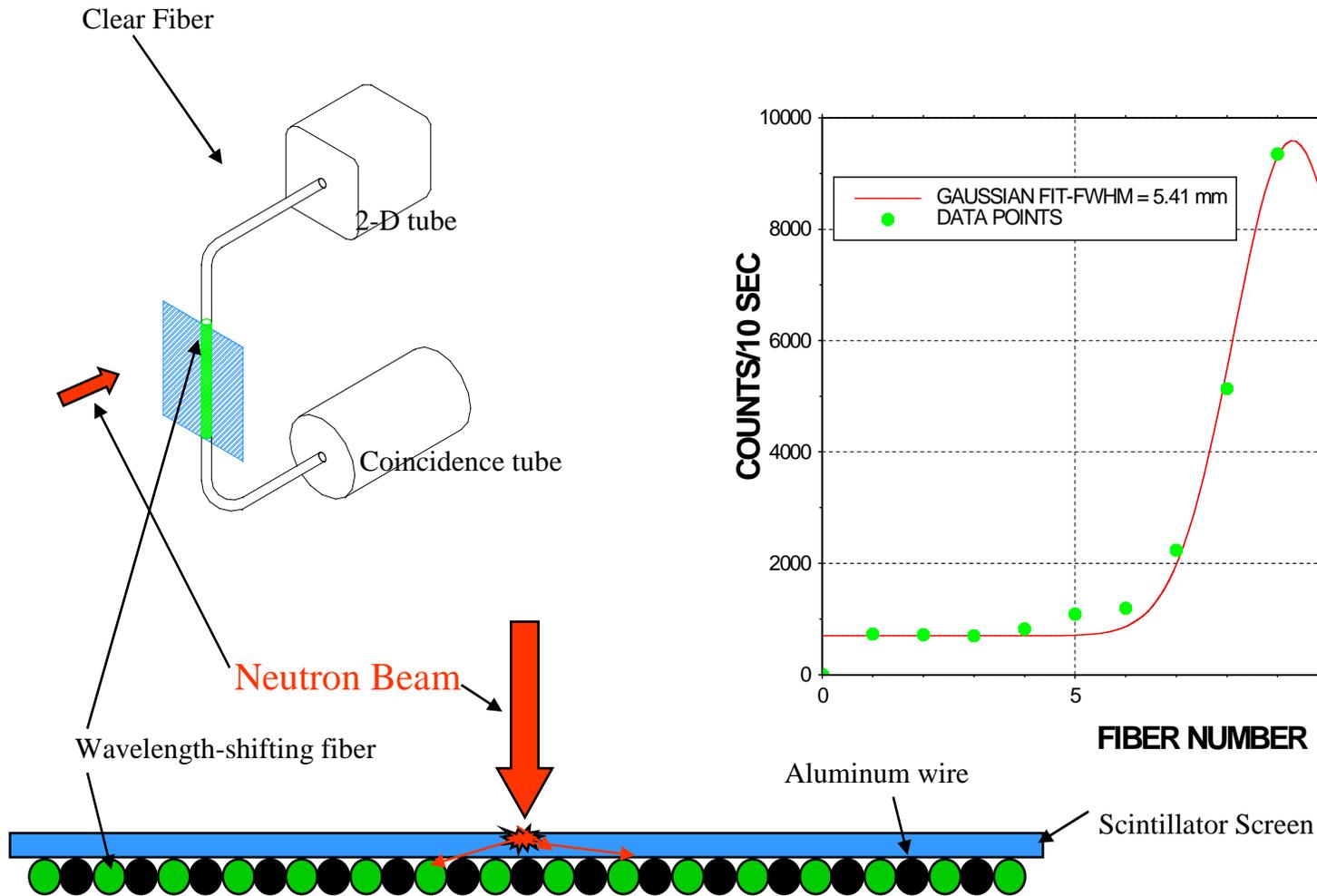
256 ch Focusing Type

64 ch Focusing Type

Principle of Crossed-Fiber Position-Sensitive Scintillation Detector



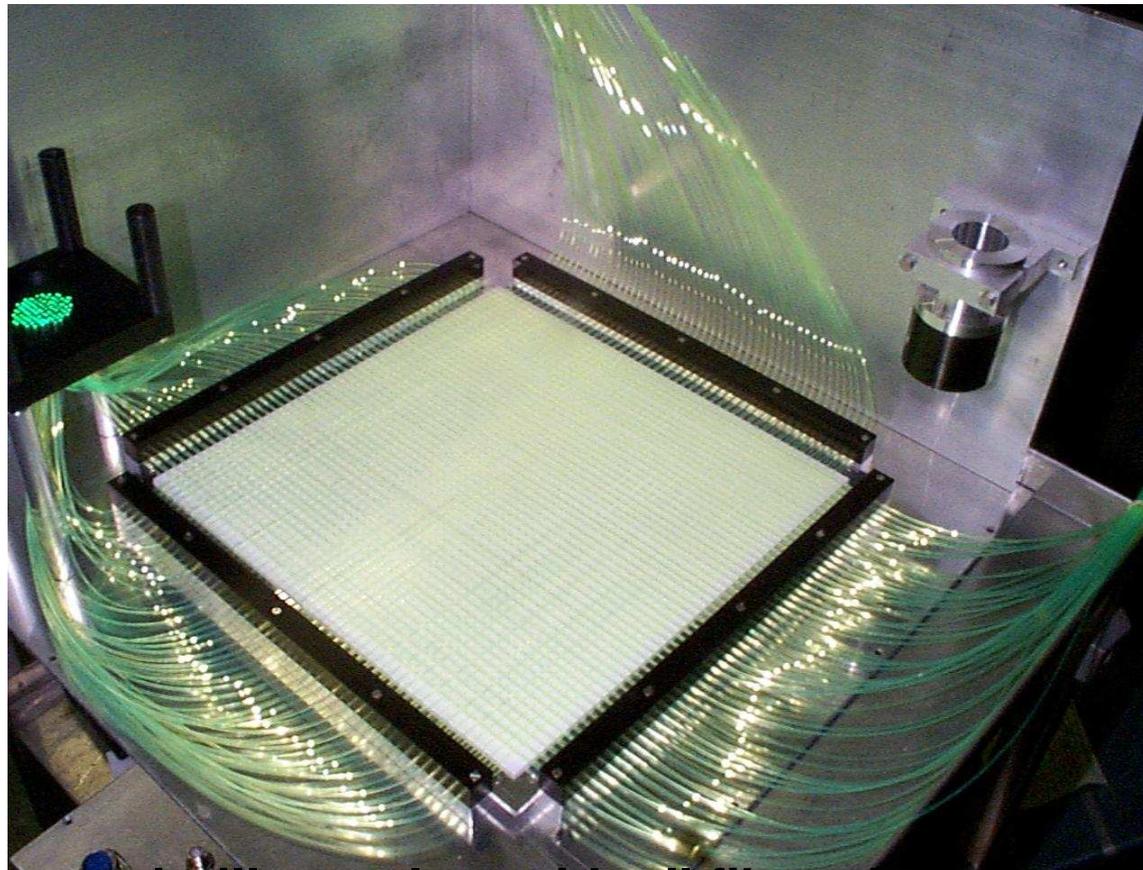
16-element WAND Prototype Schematic and Results



Crossed-Fiber Scintillation Detector Design Parameters (ORNL I&C)

- Size: 25-cm x 25-cm.
- Thickness: 2-mm.
- Number of fibers: 48 for each axis.
- Multi-anode photomultiplier tube: Phillips XP1704.
- Coincidence tube: Hamamastu 1924.
- Resolution: < 5 mm.
- Shaping time: 300 nsec.
- Counting-rate capability: ~ 1 MHz.
- Time-of-flight resolution: 1 μ sec.

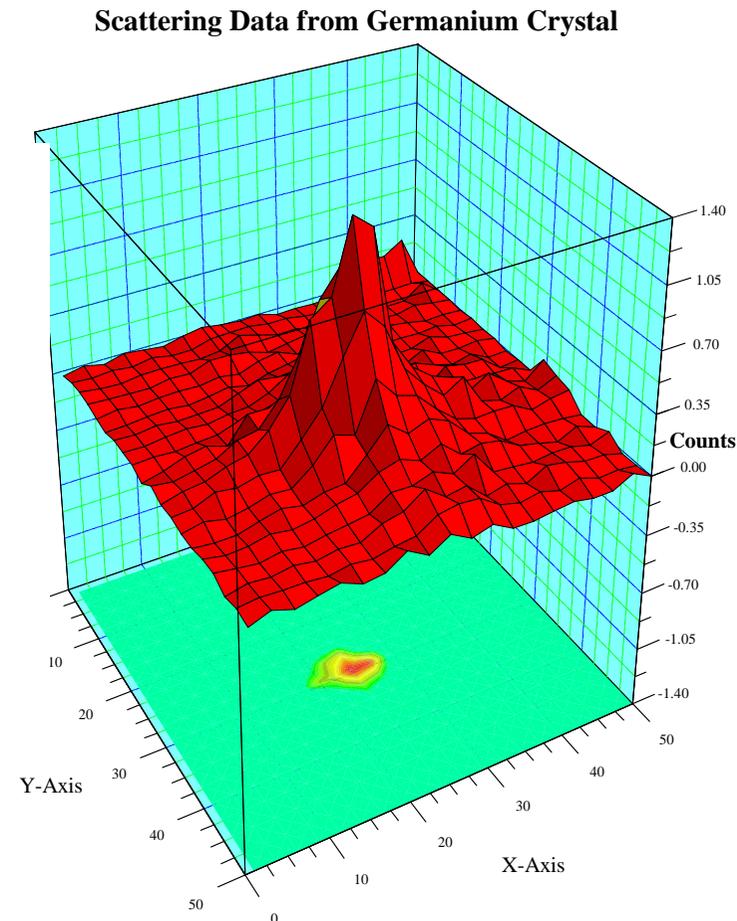
SNS 2-D Scintillation Detector Module



Shows scintillator plate with all fibers installed and connected to multi-anode photomultiplier mount.

Neutron Scattering from Germanium Crystal Using Crossed-Fiber Detector

- Normalized scattering from 1-cm-high germanium crystal.
- $E_n \sim 0.056$ eV.
- Detector 50 cm from crystal.



Neutron Detector Screen Design

The scintillator screen for this 2-D detector consists of a mixture of ${}^6\text{LiF}$ and silver-activated ZnS powder in an optical grade epoxy binder. Neutrons incident on the screen react with the ${}^6\text{Li}$ to produce a triton and an alpha particle. These charged particles passing through the ZnS(Ag) cause it to emit light at a wavelength of approximately 450 nm. The 450-nm photons are absorbed in the wavelength-shifting fibers where they convert to 520-nm photons, some of which travel toward the ends of the fibers guided by critical internal reflection. The optimum mass ratio of ${}^6\text{LiF}:\text{ZnS}(\text{Ag})$ is about 1:3.

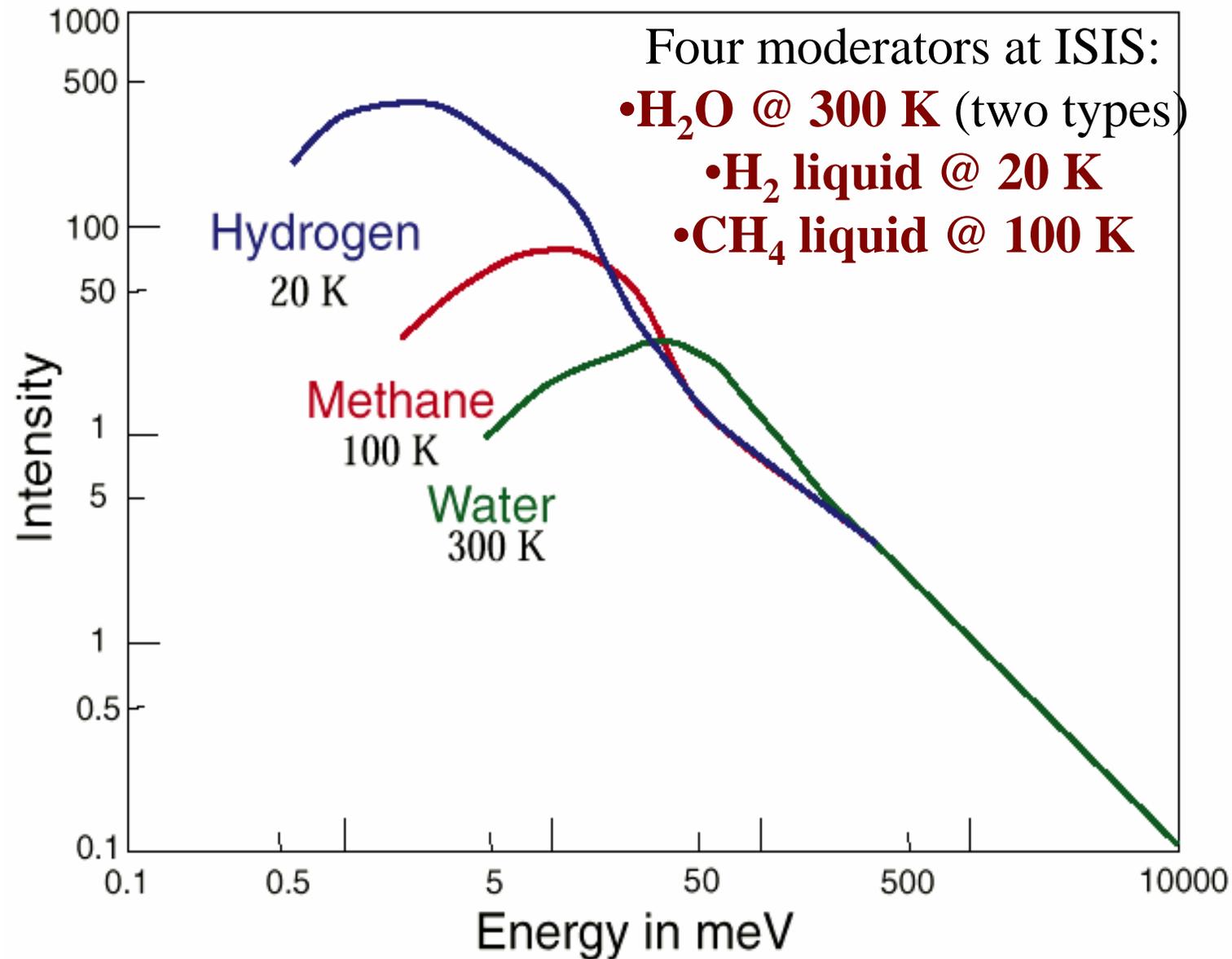
The screen is made by mixing the powders with uncured epoxy and pouring the mix into a mold. The powder settles to the bottom of the mold before the binder cures. The clear epoxy above the settled powder mix is machined away. The mixture of 40 mg/cm² of ${}^6\text{LiF}$ and 120 mg/cm² of ZnS(Ag) used in this screen provides a measured neutron conversion efficiency of over 90% for 1.8 Å neutrons.

Spatial Resolution of Area Scintillation Detectors

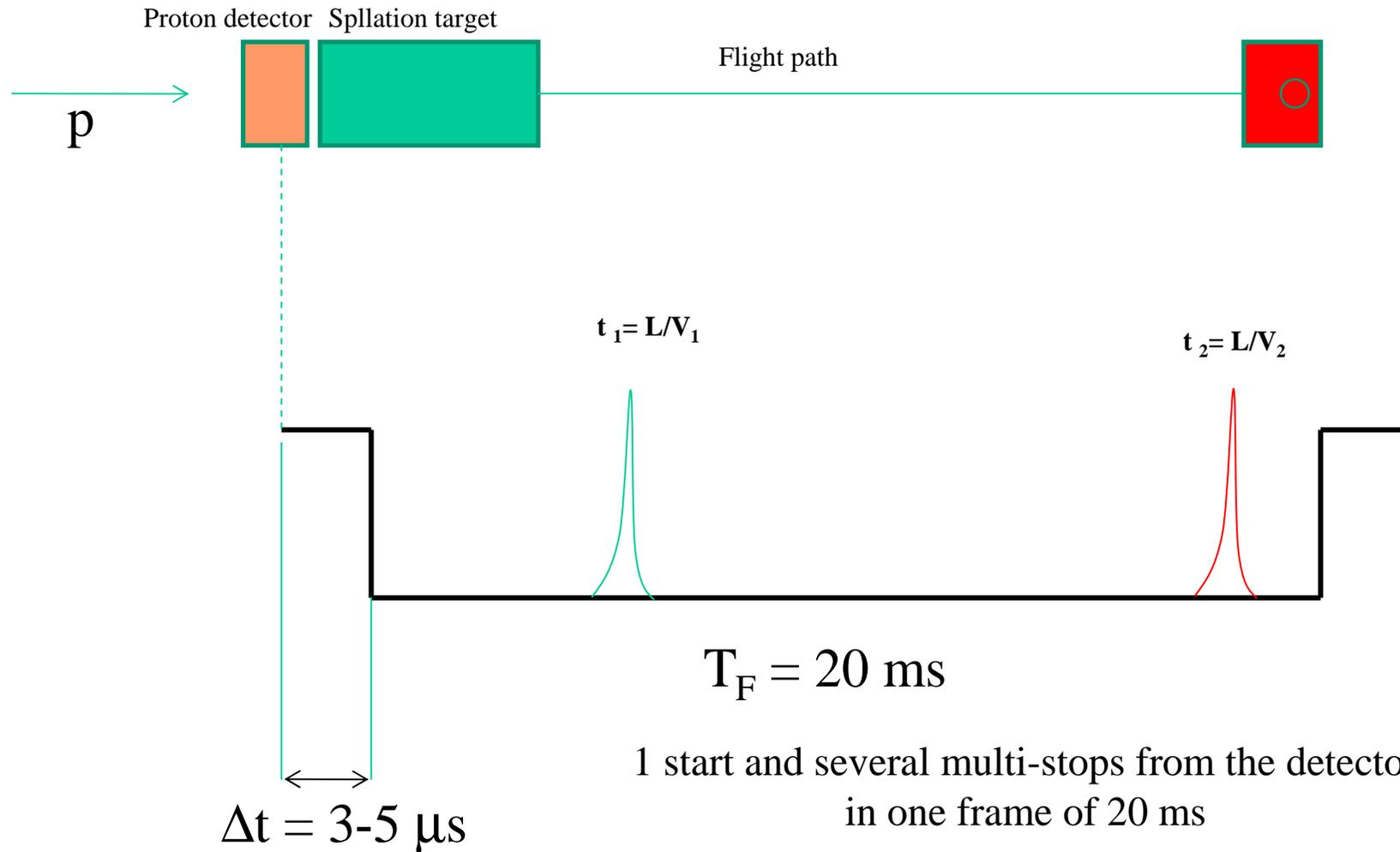
The spatial resolution accomplishable in SDs is typically better than in gas detectors. The range of neutrons is less. The range of ionizing particles is less in solid materials than in gases.

However, the localization of the light source (an optical process) imposes the limit on position resolution. This in turn depends statistically on the number of photons produced in the scintillator (more is better, of course).

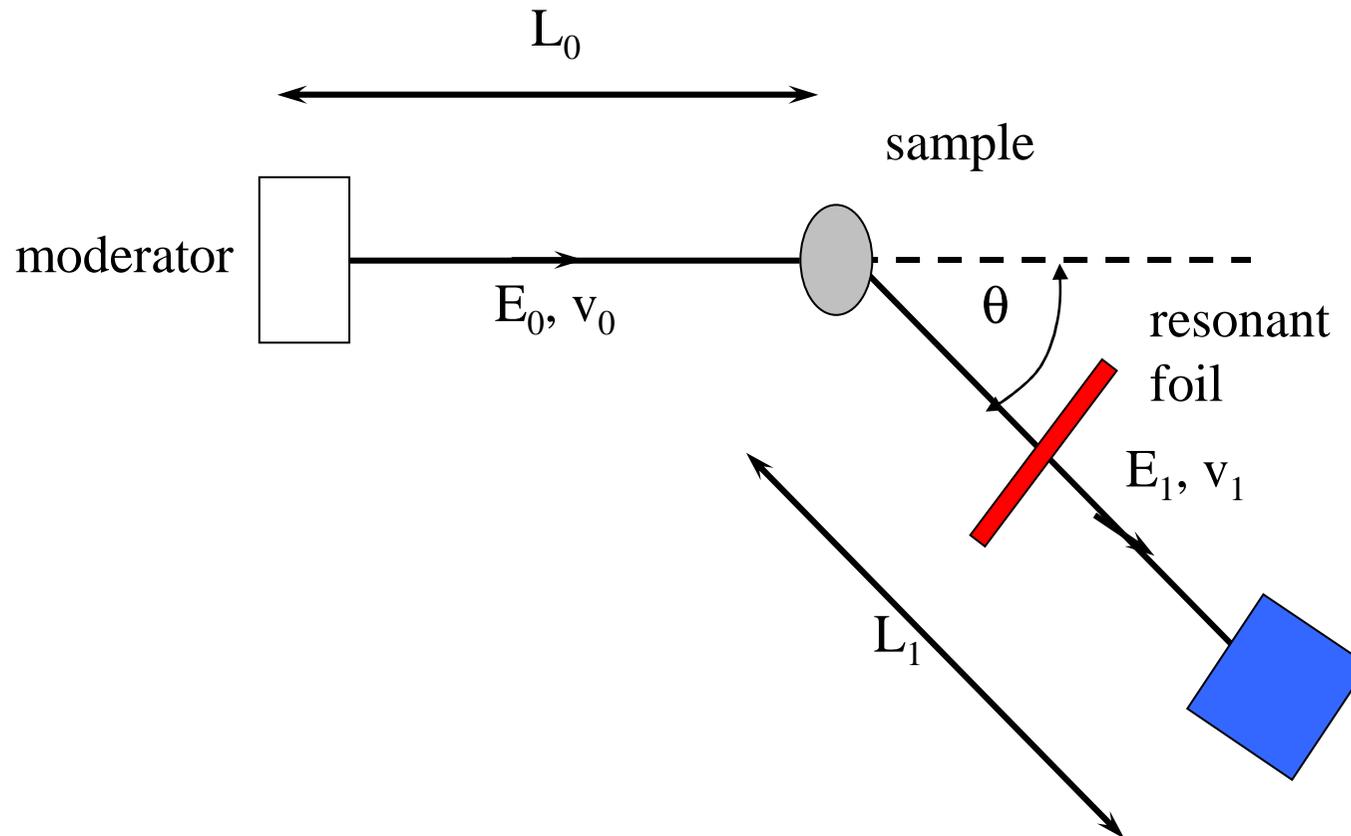
Epithermal neutrons from spallation sources



The Time of flight technique



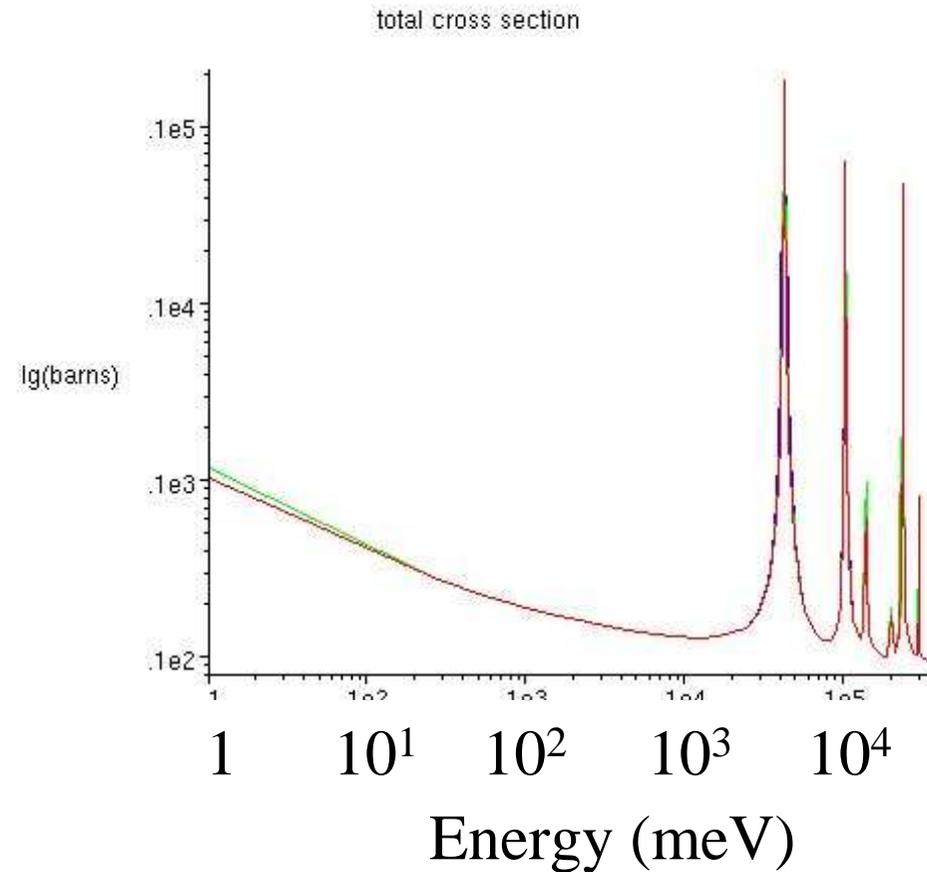
Time of flight neutron spectroscopy in inverted geometry



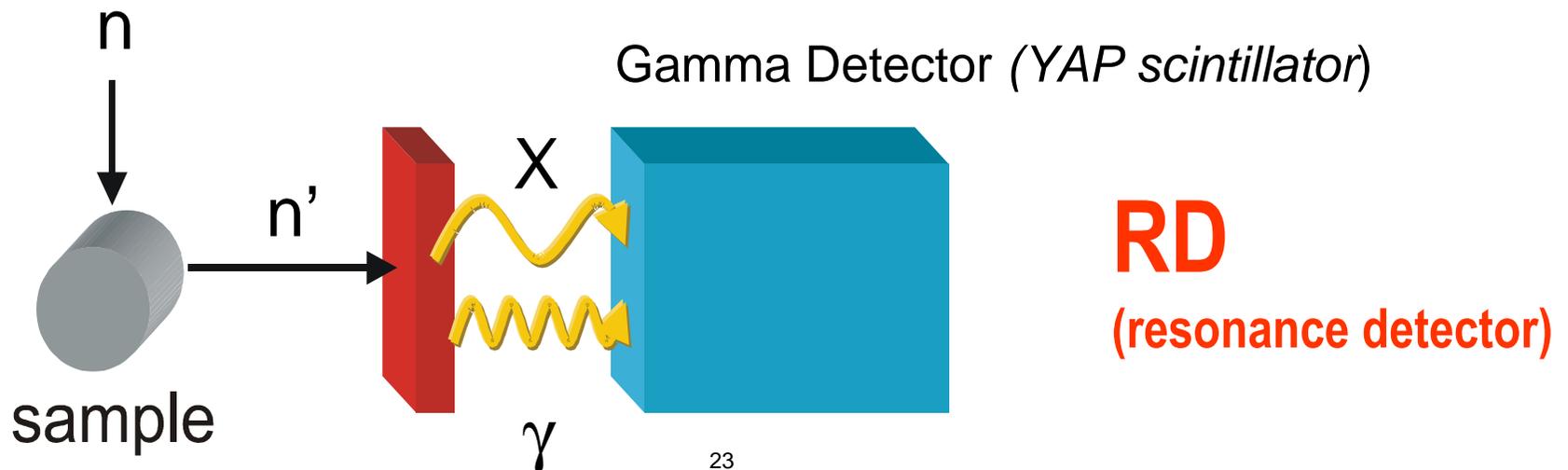
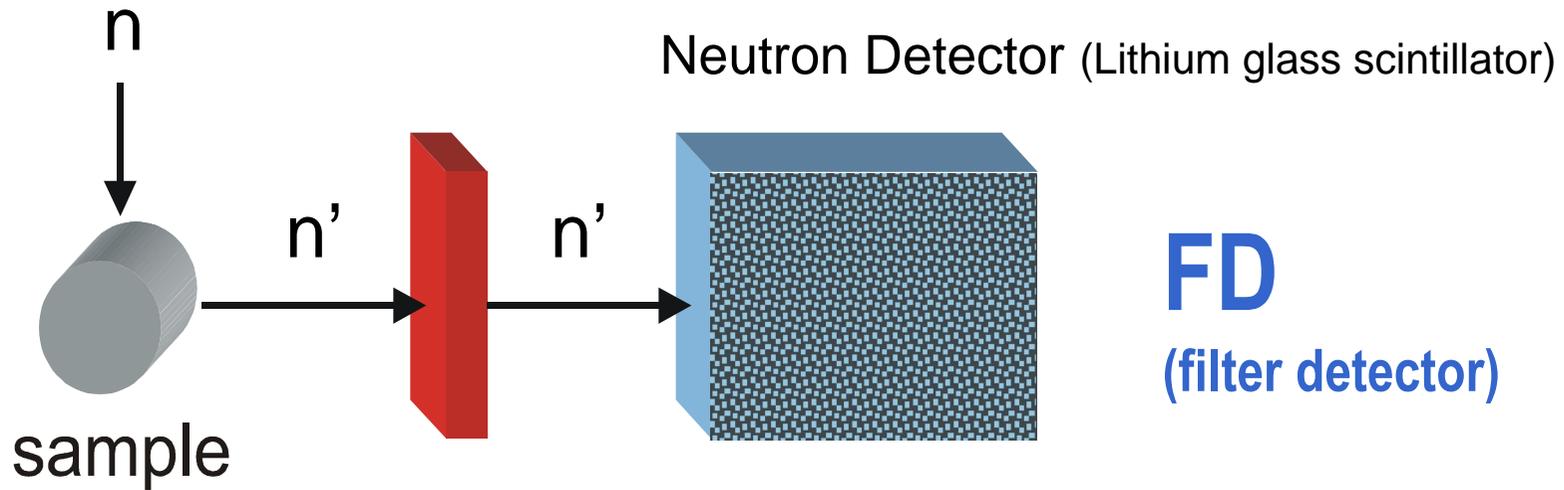
$$t = t_0 + \frac{L_0}{v_0} + \frac{L_1}{v_1}$$

Total Cross Section of Tantalum

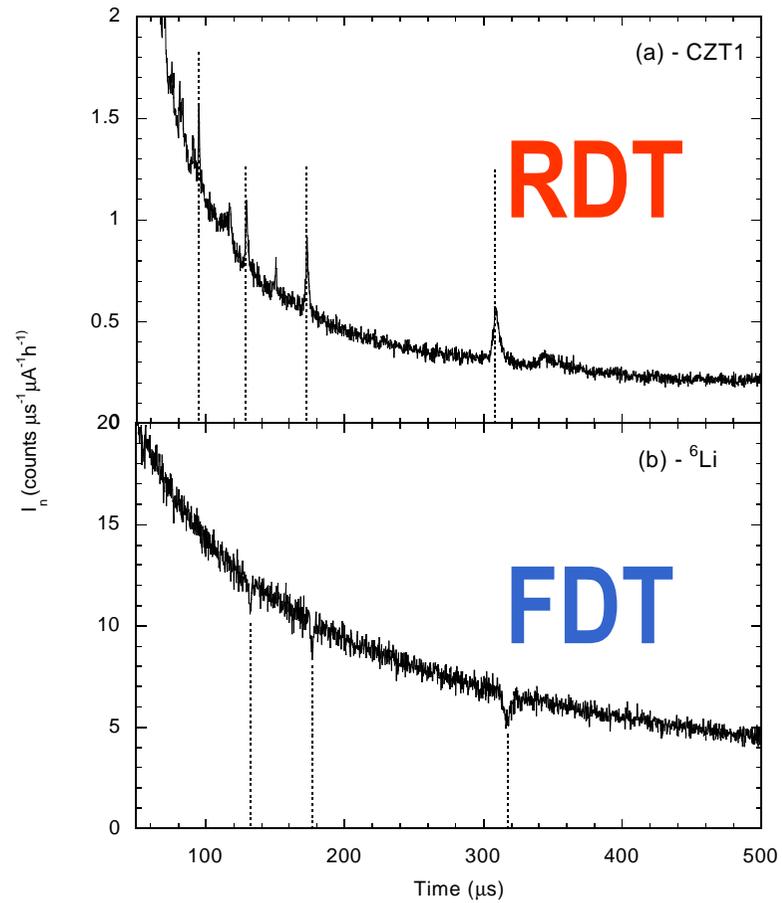
Tantalum is essentially
 monoisotopic ^{181}Ta and is
 used as a neutron converter
 sensitive to energies near
 4.28 eV.



Neutron detection techniques for epithermal neutrons



TOF spectra for FD (neutron detector) and RD (γ detector)



Pb sample
U foil

Resonance Detector principles

Two step process

- i) Neutron absorption and conversion through **(n, γ) reaction**
- ii) **Detection** of the emitted prompt gammas

RD properties

- The absorbing resonance fixes the final neutron energy
- Neutron conversion controls mainly the **energy resolution ($\Delta E/E$)**
- Ability to detect gammas controls the **signal to background ratio (S/B)**

Advantages over Li-glass detectors

- no principle need for background subtraction
- no $1/v$ decrease of the detection efficiency:
 - i) detection efficiency not dependent on the neutron energy
 - ii) high efficiency for epithermal neutrons (10-100 eV)

Choice of the resonant element

Neutron absorbing resonance

$$\sigma(E) = \frac{\sigma_0}{1 + 4(E - E_0)^2 / \Gamma^2}$$

(Breit Wigner)

E_0 neutron resonance energy

σ_0 peak cross section (at $E=E_0$)

Γ (FWHM) intrinsic resonance width

Criteria for resonant elements

- Resonance energy: **$E_0=10-100$ eV**
- Isolated resonance (to avoid overlapping with other resonances)
- High cross section (high conversion efficiency) \Rightarrow **$\sigma_0 = 10^4-10^5$ b**

- Narrow resonance (low $\Delta E/E$)

⇒ $\Gamma \sim 100 \text{ meV}$

⇒ $\Delta \sim 100 \text{ meV}$ (Thermal Doppler broadening)

- Foil thickness

Compromise between neutron absorption probability, $P(E_n)$, and $\Delta E/E$

$$P(E_n) = 1 - \exp(-N_d \cdot \sigma) \qquad N_d = \rho \cdot t$$

⇒ $N_d \cdot \sigma_0 = 1$

- High yield of low energy gammas (10-300 keV) ⇒ choice of suitable detector

Other desirable features: *high isotopic abundance, exist as metallic or oxide, low gamma self-absorption*

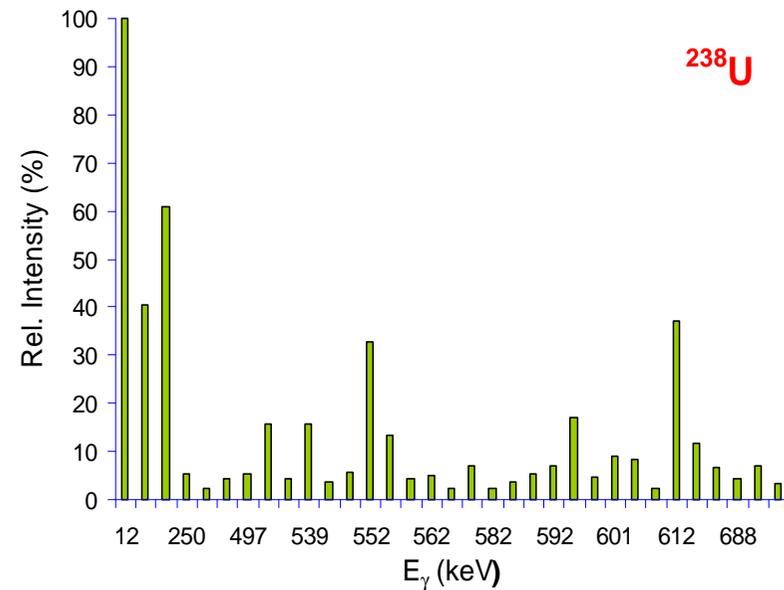
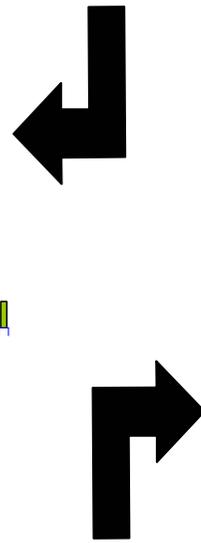
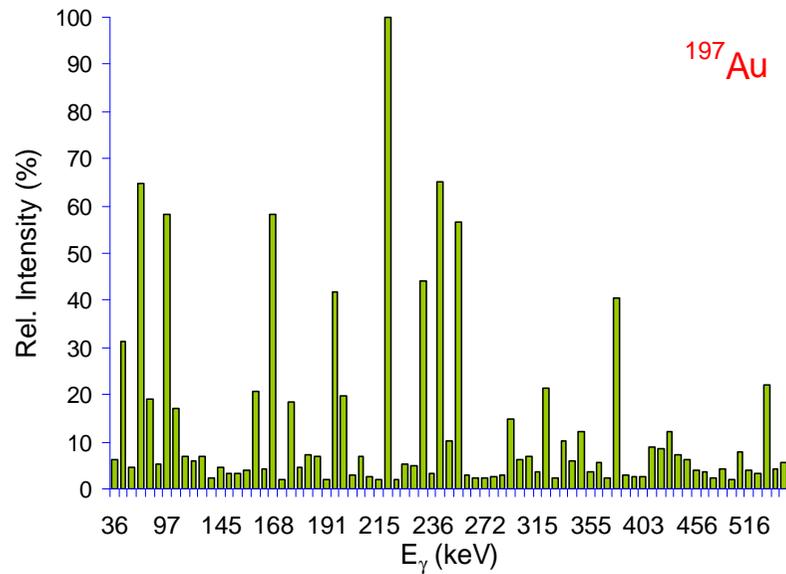
Best Isotopes

	i.a. (%)	density (g/cm ³)	E ₀ (eV)	σ ₀ (barn)	Γ (meV)	Δ (T=295 K) (meV)	Δ (T=75 K) (meV)
¹³⁹ La ₅₇	99.9	6.1	72.2	5969	96.0	231	125
¹⁵⁰ Sm ₆₂	7.4	7.4	20.7	56207	109.0	119	66
²³⁸ U ₉₂	99.3	18.9	6.7	23564	25.0	54	31
	“	“	20.9	37966	34.0	96	55
	“	“	36.7	42228	57.0	127	73
	“	“	66.0	20134	48.0	170	98

Gamma yield (absolute) exists for thermal neutron capture, but incomplete database for resonant neutron capture

¹³⁹ La ₅₇	32 %	218 keV	9 %	²³⁸ U ₉₂	12 keV
	13 %	288 keV	9 %		48 keV
					60 keV
¹⁵⁰ Sm ₆₂	94 %	65 keV	30 %		134 keV
		105 keV	9 %		4060 keV

Energy and relative intensities of γ -rays



Suitable γ detectors for the RDS

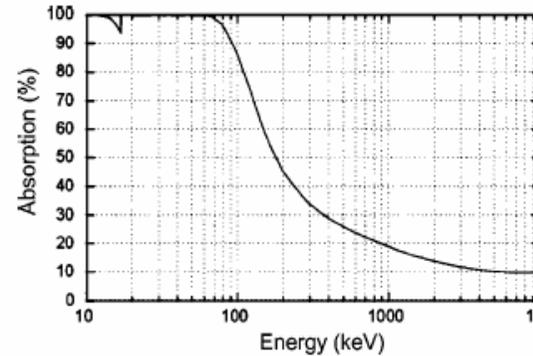
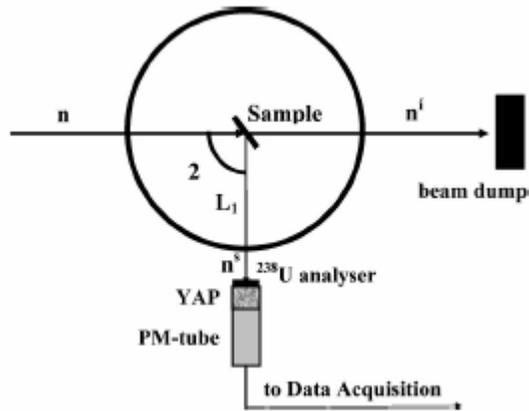
Scintillators

- 1) NaI(Tl): good efficiency and energy resolution for MeV gamma ray. High sensitivity to neutron background (need of shielding).
- 2) **YAP**: good efficiency and moderate energy resolution. Low sensitivity to neutron background

Solid State

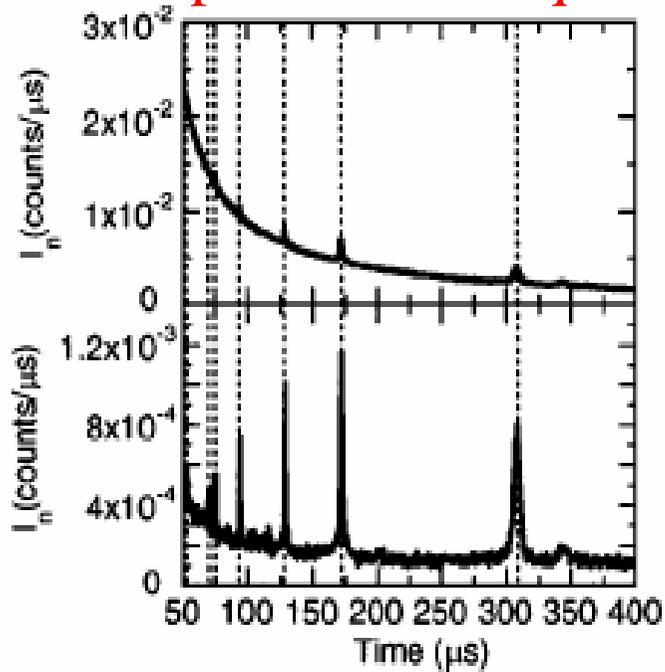
- 1) HpGe: excellent energy resolution, adequate efficiency. Radiation damage. Need to operate at cold temperature (77K)
- 2) Silicon: good for x-ray and low energy γ ray. Requires some cooling. Radiation damage
- 3) CdZnTe (CZT): Good efficiency and high energy resolution. Operate at room temperature. Small areas.

YAP detector



6mm thick

Biparametric acquisition



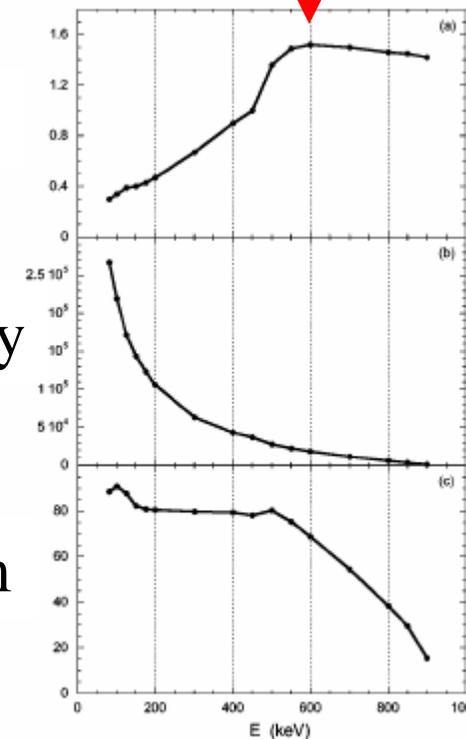
>40keV

>600keV

S/B

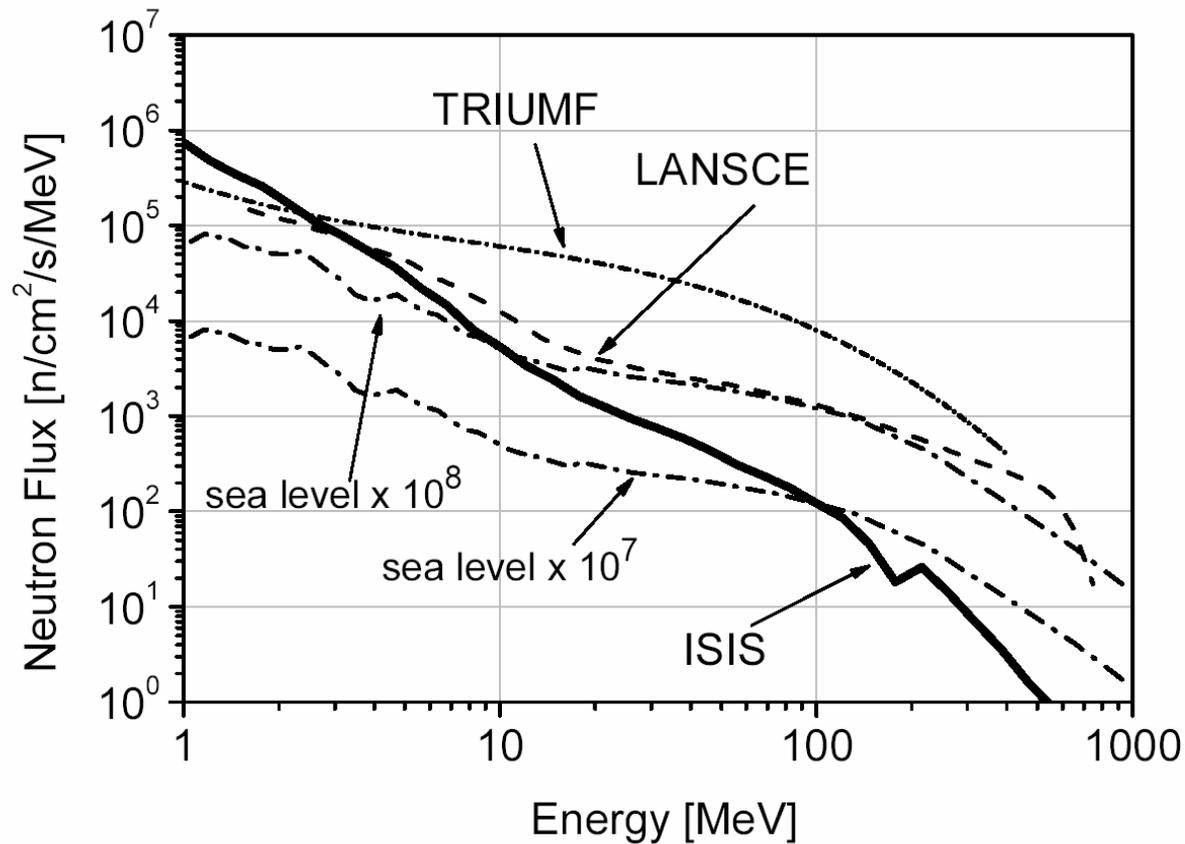
Intensity

fom



Fast neutrons

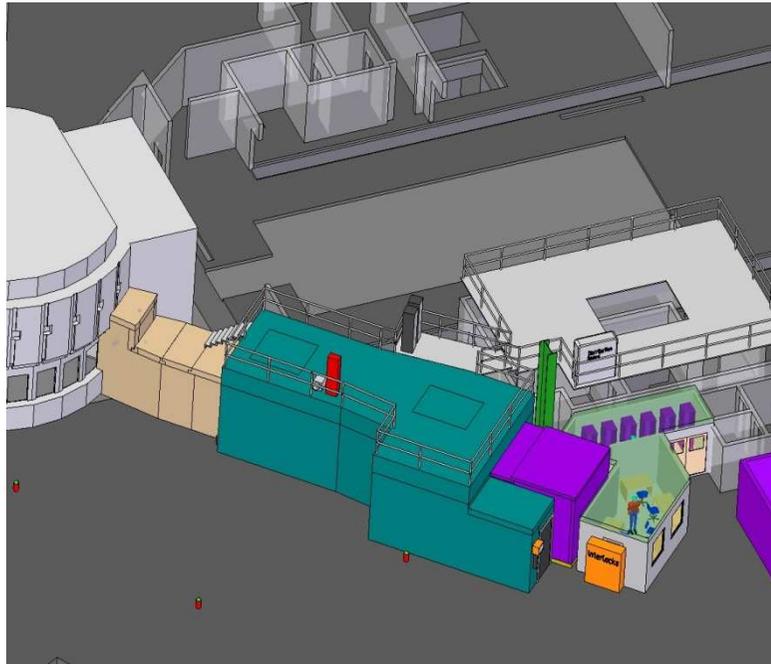
Fast neutron flux of VESUVIO at ISIS



Measured with
activation
targets

Enormous energy range 1-800 MeV!

CHIPIR beam line at ISIS



Beam line dedicated to samples with dimension from a single chip to an entire electronic system.

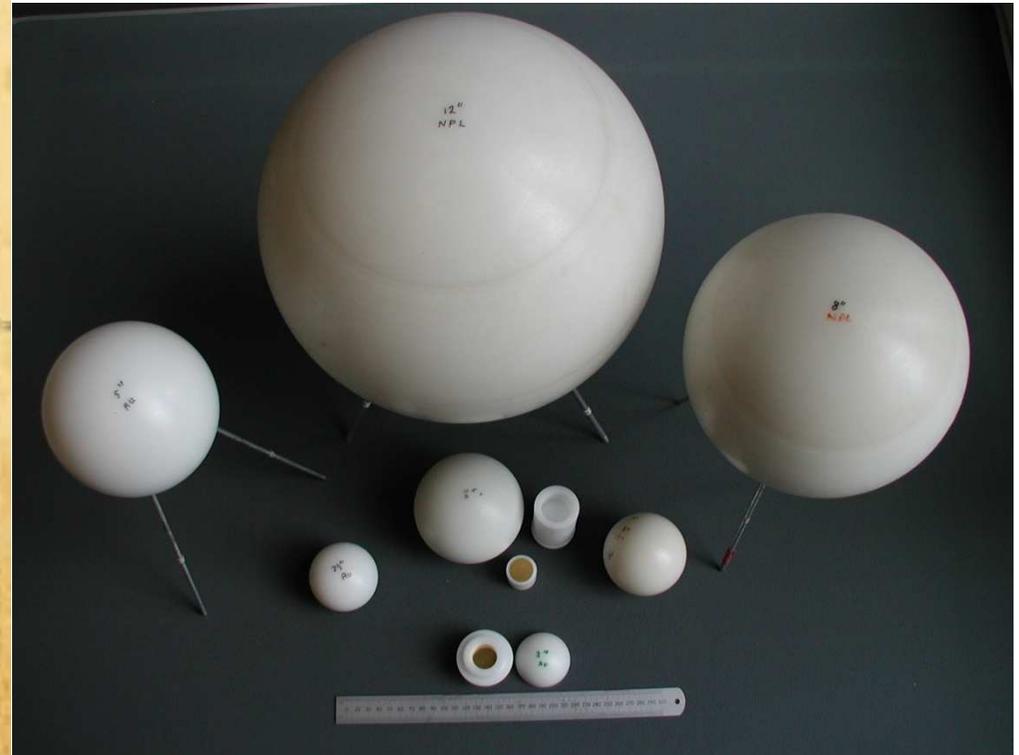
$$\Phi \approx 10^6 \text{ n/cm}^2/\text{s} \text{ above } 10 \text{ MeV}$$

One hour on CHIPIR will be equivalent to 114 years on a plane

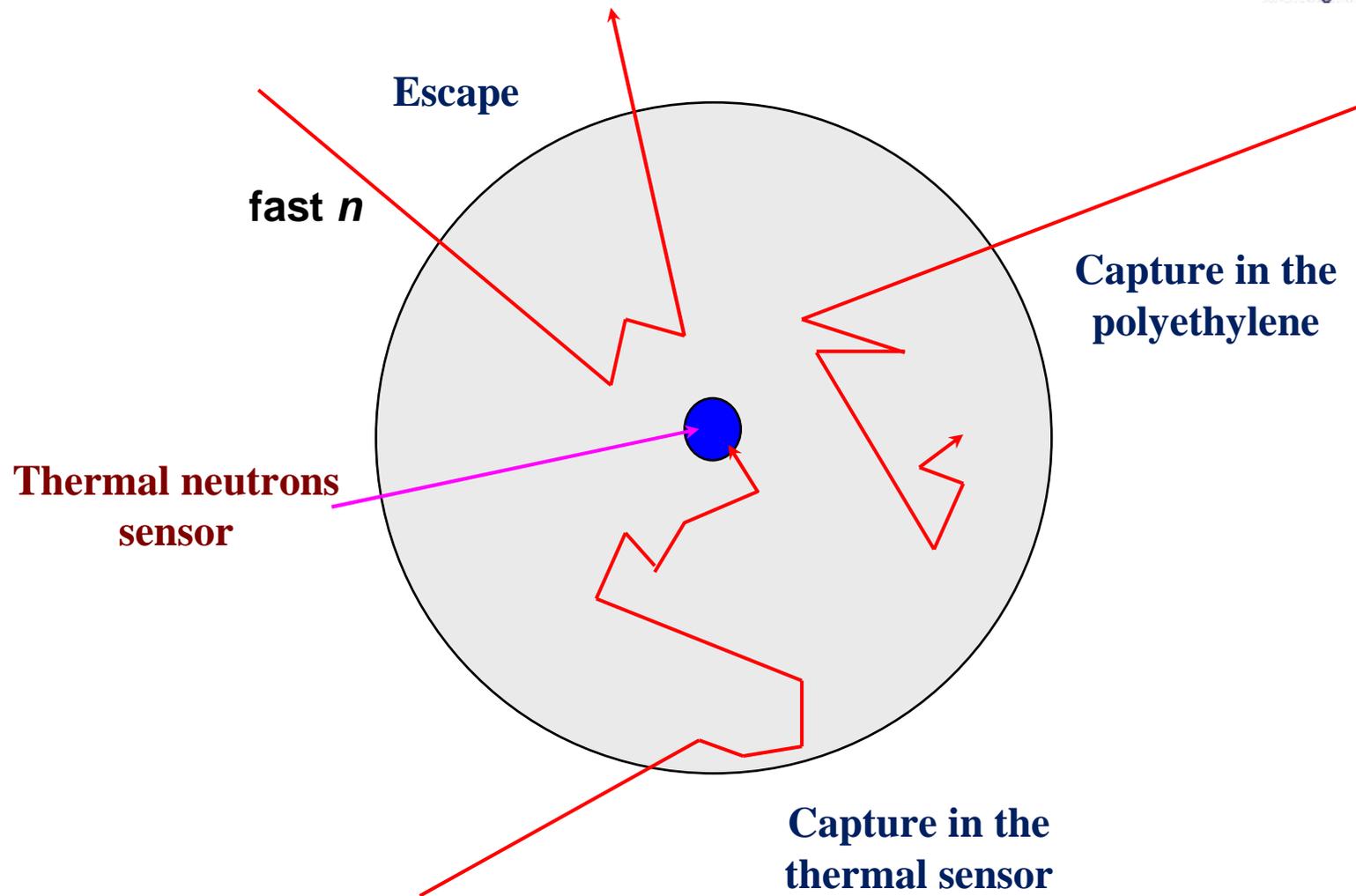
Need for beam monitor of the fast neutrons for the flux measurements

Fast neutrons detection

- Spherical dosimeters
- Scintillators
- ^3He proportional counters
- Proton Recoil Telescope (spectroscopy)
- Activation threshold targets
- TFBC
- Bonner Sphere
- Diamond detectors
- nGEM detector

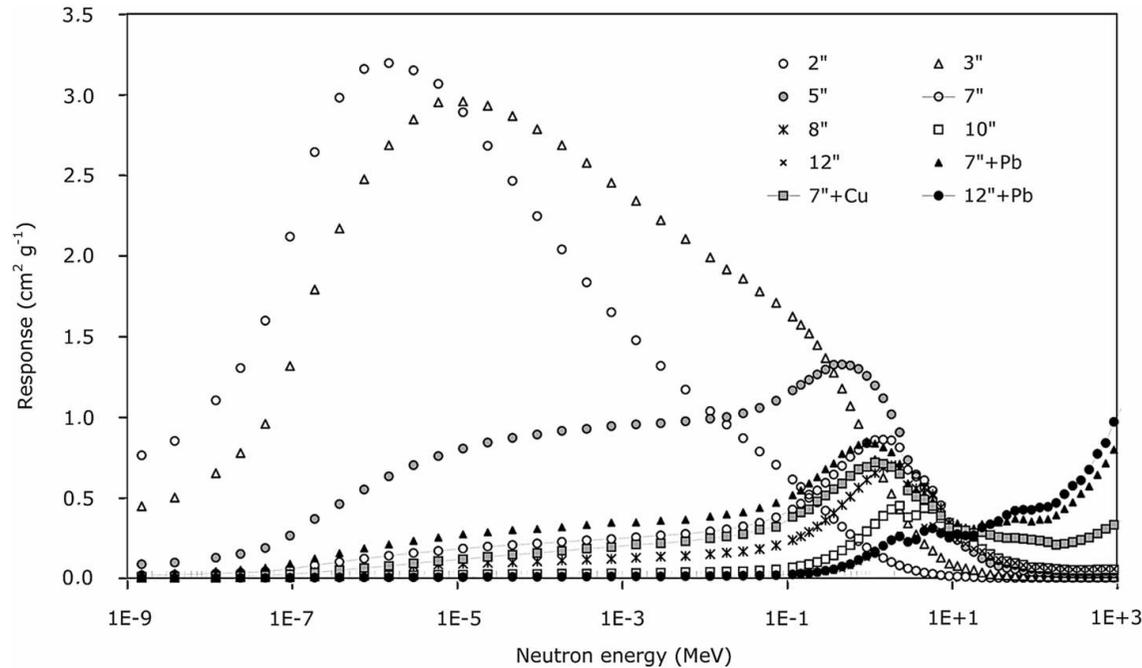


Principle of operation



From the measured activity on the sensor after a period of irradiation the neutron flux at the selected energy range is found

Response Functions of Bonner Spheres



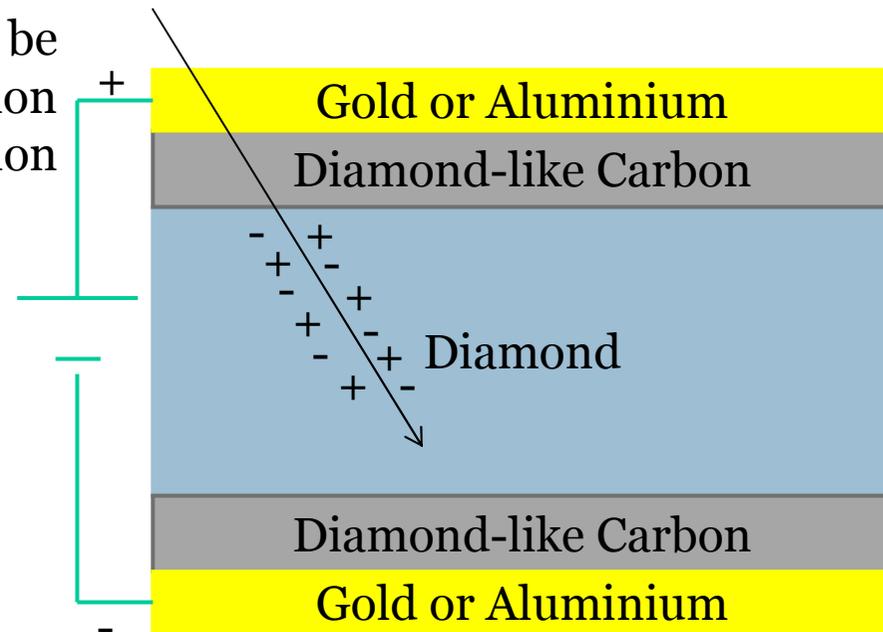
The very broad response function means that complicated deconvolution codes are needed to infer the incoming neutron spectrum from the measurement.

Diamond Detector

- Radiation hardness.
- High mobility of free charges (\rightarrow fast response, comparable to Si, Ge).
- Room temperature operation ($E_g=5.5$ eV) \rightarrow No Cooling.
- Compact volume solid state detector.

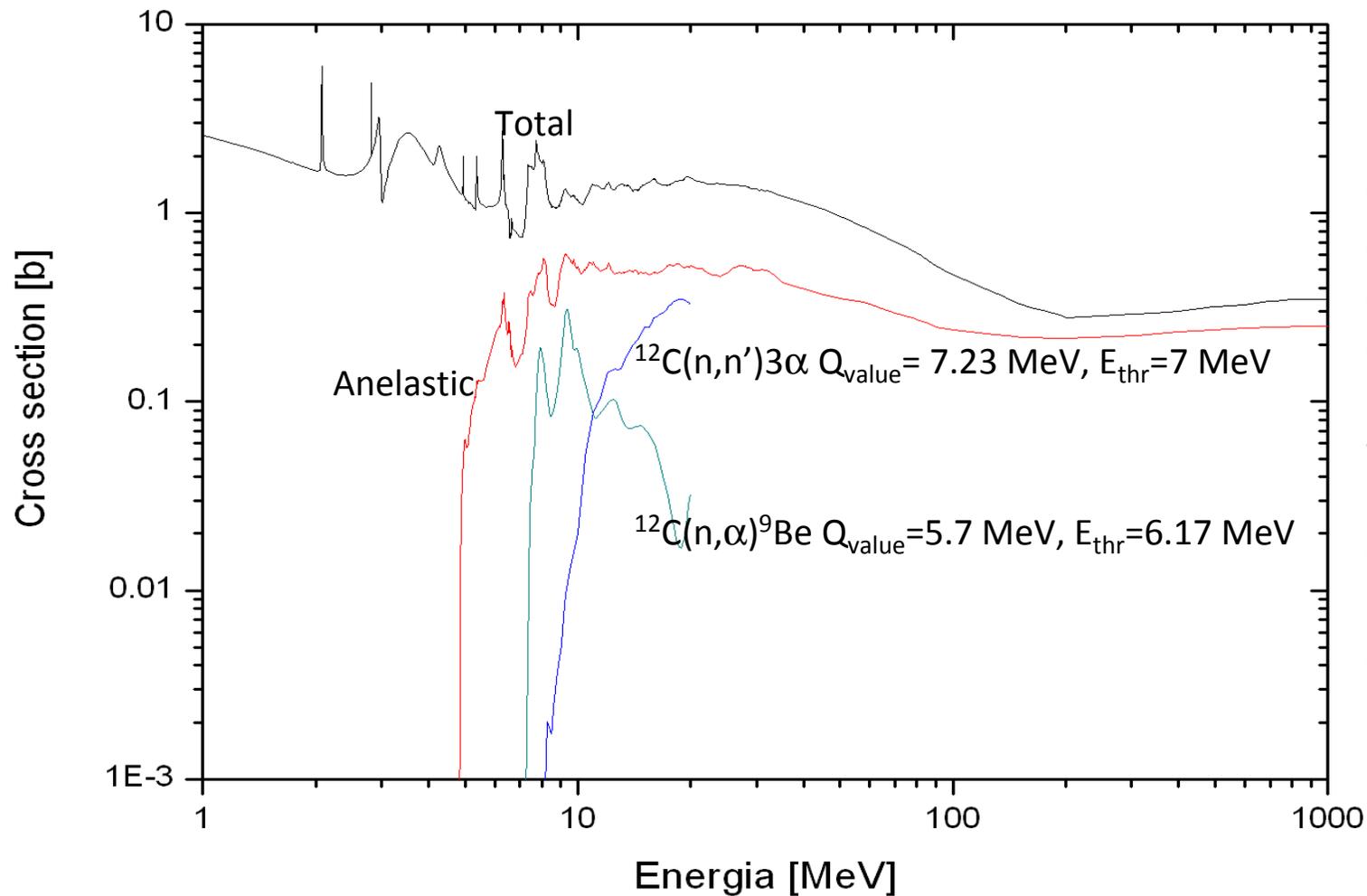
With the CVD technique diamonds can be produced with good energy resolution (<1%) and 100% charge collection efficiency.

A charged particle passes through the diamond and ionizes it, generating electron-hole pairs ($E_{e-h}=13$ eV)



Diamond Detectors Limited Technology

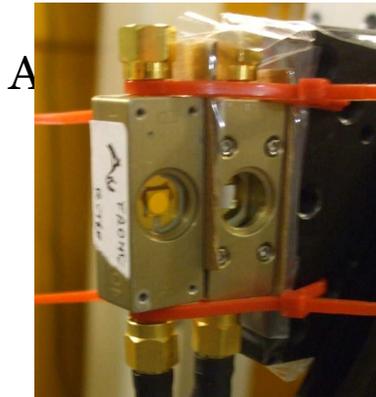
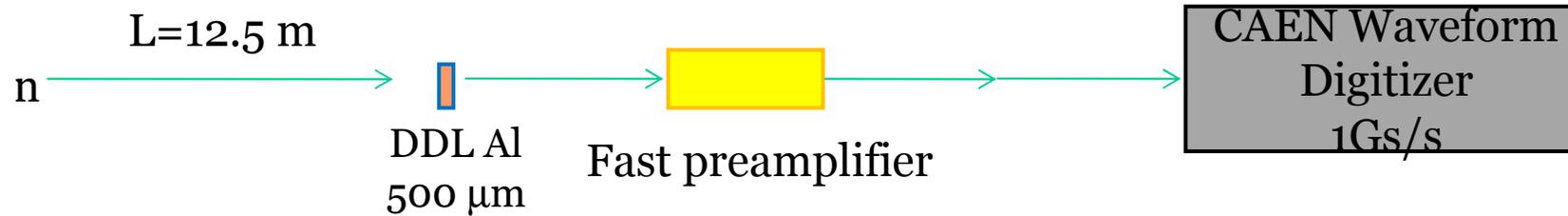
Carbon cross section



→ Fast neutron detection is achieved by detecting charge particles produced via the reactions:

- $^{12}\text{C}(n,\alpha)^9\text{Be}$
- $^{12}\text{C}(n,n')3\alpha$
- $^{12}\text{C}(n, n')^{12}\text{C}^*$

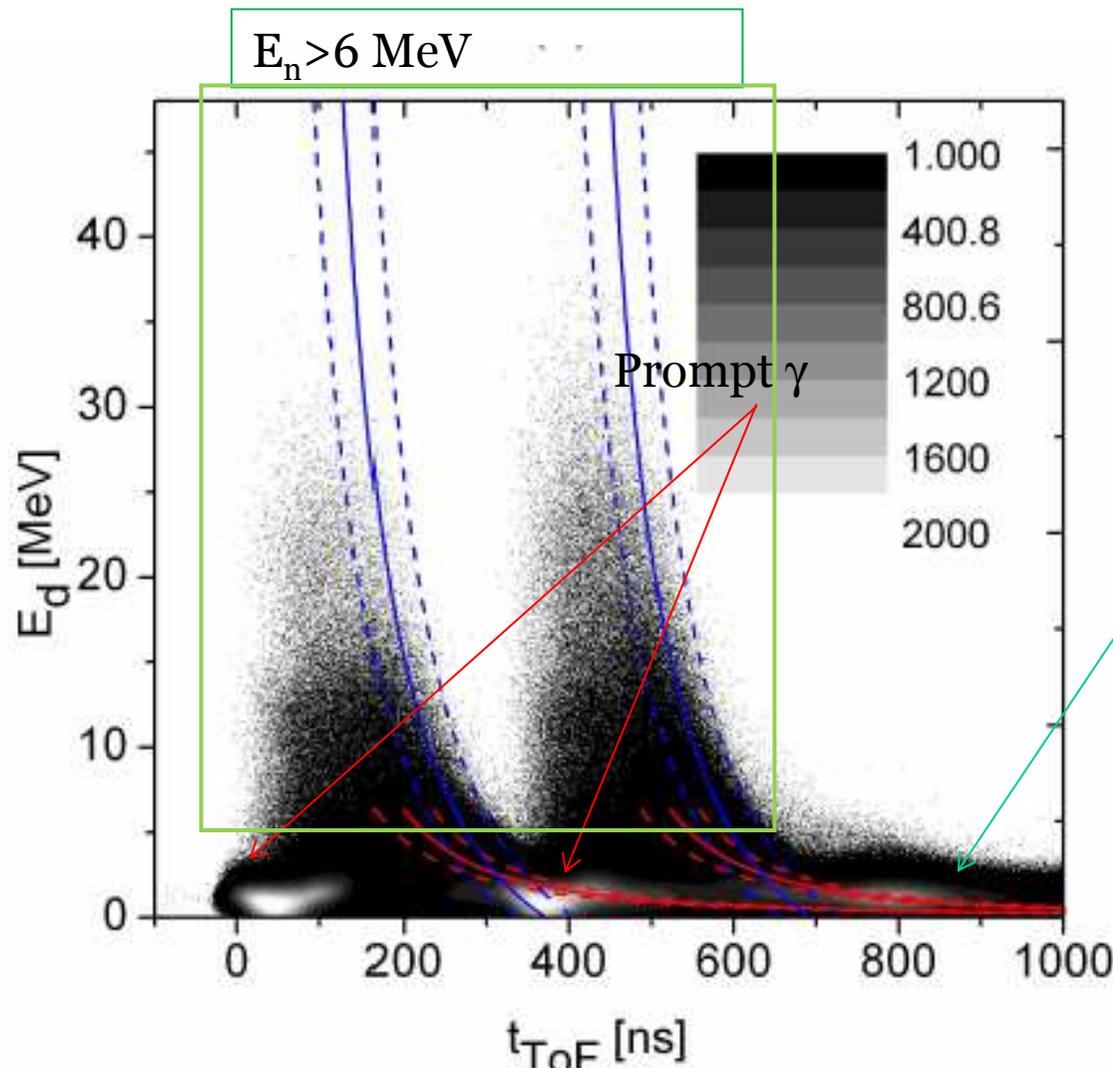
Single-crystal Diamond Detector



$4 \times 4 \times 0.5 \text{ mm}^3$ sCVD
 Al or Au contacts

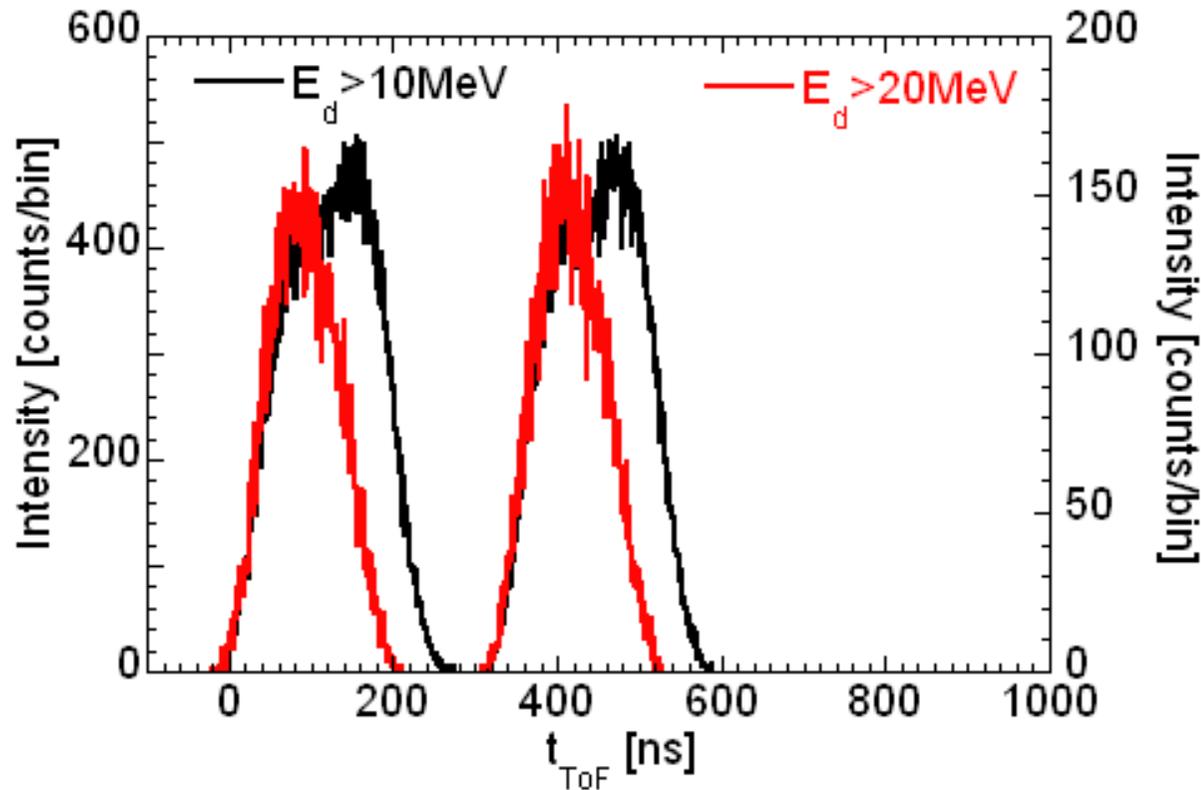
	Si	Ge	CVD-Diamond
Atomic Number Z	14	32	6
Density [g/cm ³]	2.33	5.33	5.47
Band Gap [eV]	1.1	0.6	5.5
Electron Mobility [cm ² /V · s]	1350	3900	1800
Hole Mobility [cm ² /V · s]	480	1900	1200
Breakdown field [MV/cm]	0.3	0.1	10

Contour plot



Events due to the interaction of 3.5 MeV neutron which release energy via elastic recoil on the ^{12}C nucleus.

ToF spectra for high E_{dep} events



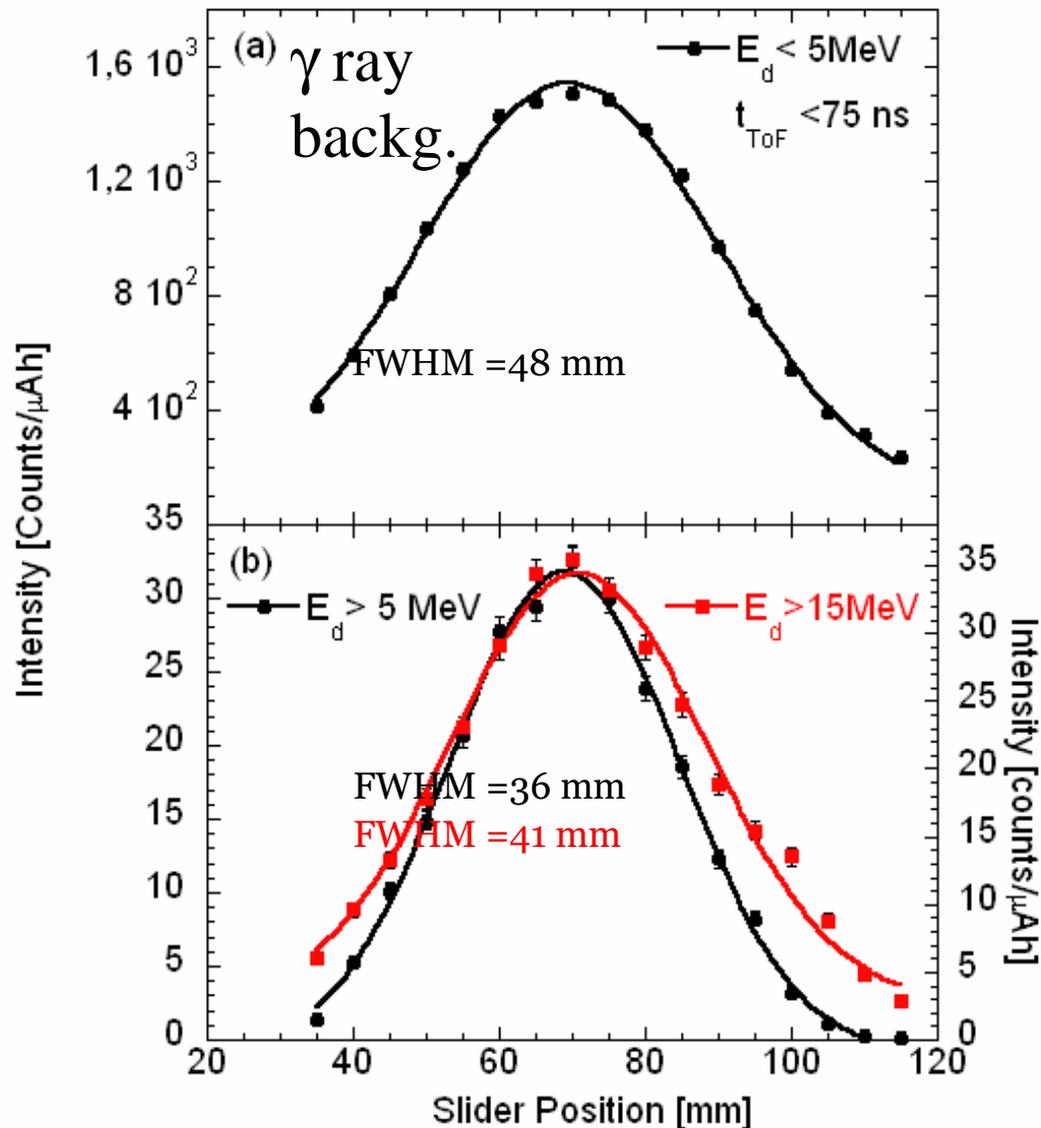
Again a clear correlation between t_{ToF} and E_{d} is observed: the maximum t_{ToF} is shorter for the higher energy.

A neutron that deposits 10 MeV in the SDD should have $E_{\text{n}} > 15.7$ MeV; that is $t_{\text{ToF}} < 230 \pm 35$ ns. A neutron that deposits 20 MeV should have $E_{\text{n}} > 25.7$ MeV, i.e. $t_{\text{ToF}} < 182 \pm 35$ ns.

Time of flight spectrum for neutrons which deposit > 10 MeV and > 20 MeV

→ For a quantitative analysis knowledge of the SDD response to monoenergetic neutrons is needed.

Beam profile measurement

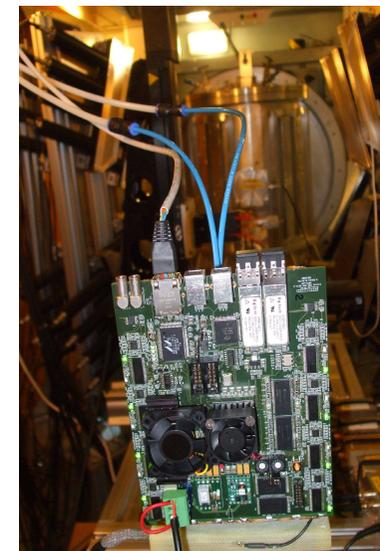
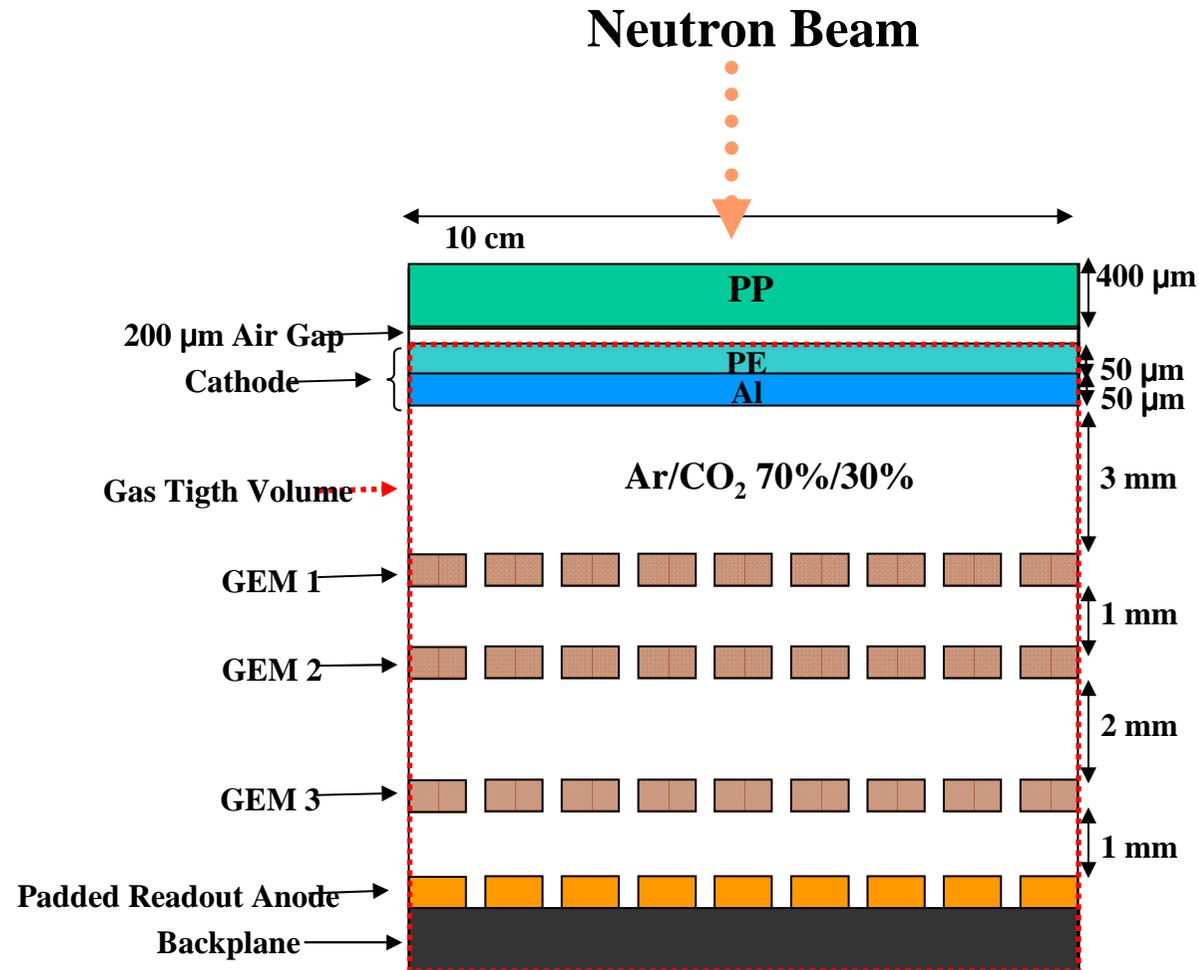


Horizontal beam profile obtained by selecting events with:

- (a) $E_d < 5 \text{ MeV}$ and $t_{\text{ToF}} < 75 \text{ ns}$,
- (b) $E_d > 5 \text{ MeV}$, $200 < t_{\text{ToF}} < 250 \text{ ns}$ (black), and $E_d > 15 \text{ MeV}$ (red).

The differences in profile width are well outside the uncertainties in the measurement.

nGEM Detector



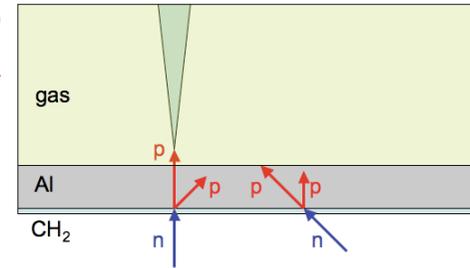
nGEM Detector Components

Optimized for 2.5 MeV neutrons detection

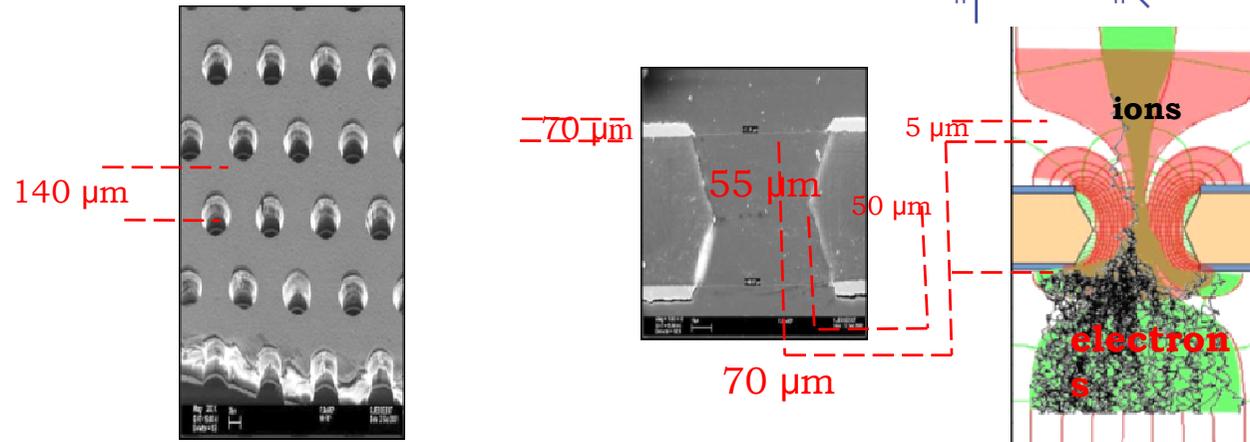
Converter Cathode

Composed by two layers of **CH₂ 50 μm thick** and of **Aluminium 50 μm thick**

Recoil protons produced by neutrons release energy in the gas



GEM Foils



Padded Readout Anode

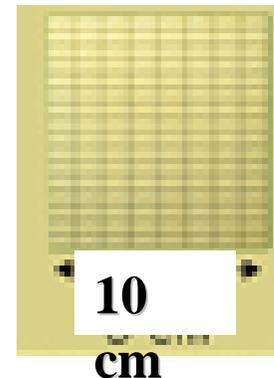
128 Pads

Pad Area

12 x 6

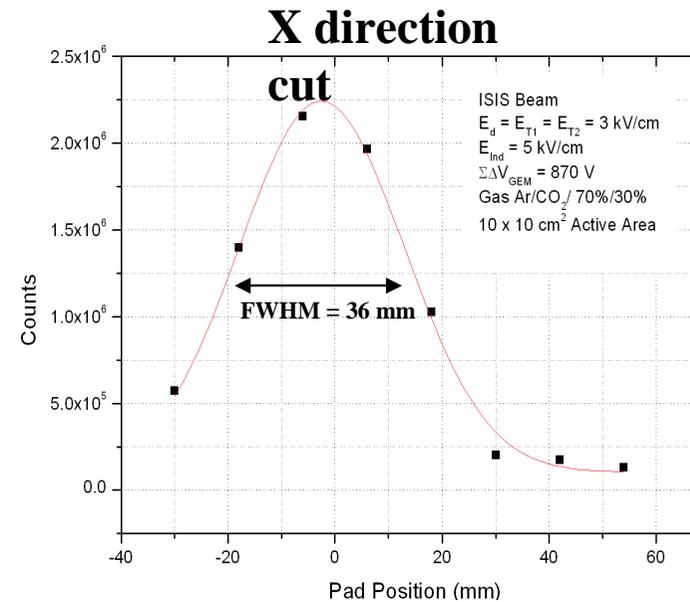
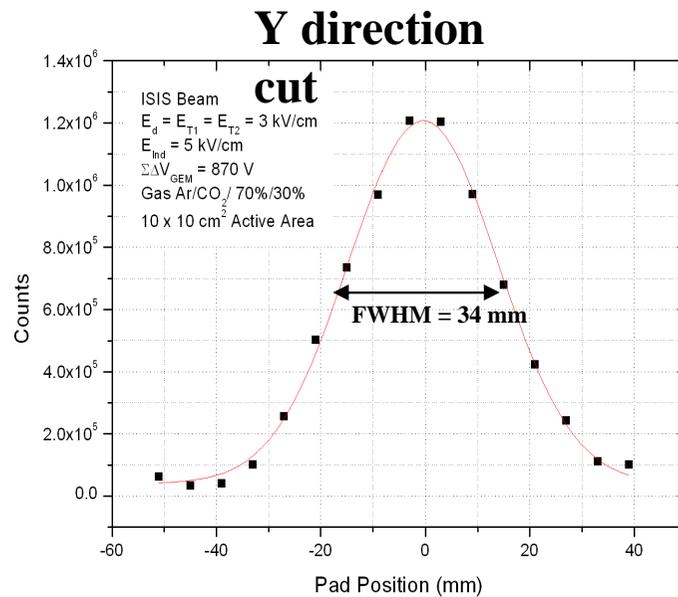
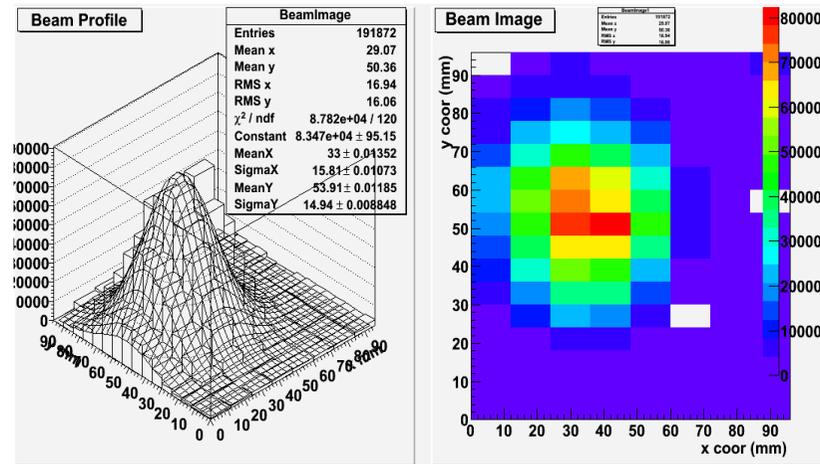
mm²

45

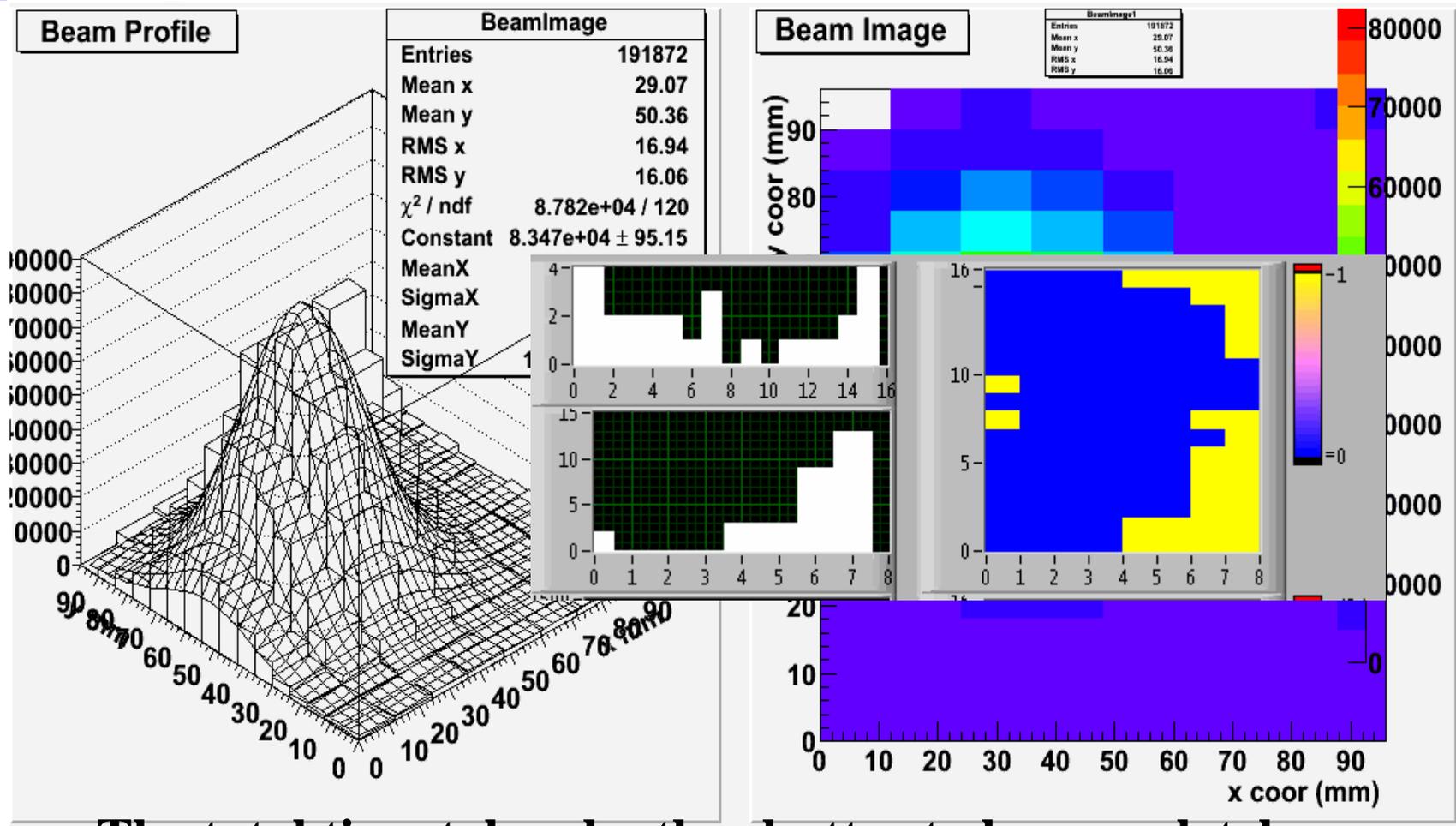


Vesuvio Beam Profile Measurement

Fast Neutron Intensity Map



Opening of Vesuvio shutter



The total time taken by the shutter to be completely open is about 11 seconds

END