

Magnetic Nanostructures Investigated by Small Angle Neutron Scattering

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VIII School of Neutron Scattering "Francesco Paolo Ricci" Structure and Dynamics of magnetic systems

http://www.hmi.de/bereiche/SF/SF3/methods/sans/index_en.html

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Scientific committee decides

http://www.hmi.de/bensc/

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Content



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Scattering law Nuclear and magnetic scattering Scattering of polarised neutrons (SANSPOL, POLARIS)

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- **1.3 Magnetic scattering**
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- **1.5 Instrument**



0.5-2 nm 0.5-300nm 0.1-20°

fluctuations of density, composition, magnetization



SAXS (X-rays): Interaction with electrons: Scattering length: b_x = Z*0.282 10⁻¹² cm

SANS (neutrons): Interaction with nucleousIntegralScattering length: b_n (1 H: -0.2, 2 D: +0.5)Between neutron spin and magnetic moment μ Scattering length: b_{mag} = 0.27 *10⁻¹² * μ

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Local

Scattering cross-section

h	m	:



phase 1 $\eta_1 = \sum c_i b_i / \Omega_i$ volume fraction f phase 2 $\eta_2 = \sum c_i b_i / \Omega_i$ volume fraction (1-f) average scattering length density $<\eta> = f \eta_1 + (1-f) \eta_2$ Deviation from $<\eta>$: $\delta_1 = \eta_1 - <\eta> = (1-f)(\eta_2 - \eta_1)$ $\delta_2 = \eta_2 - <\eta> = f(\eta_2 - \eta_1)$

Mean square fluctuation of scattering length densities

$$\langle \eta^2 \rangle = f \delta_1^2 + (1-f) \delta_2^2 = f(1-f) (\eta_1 - \eta_2)^2$$

Nuclear contrast-matching

Scattering from different parts of particle (Micelles in solvents)

a) Isotope mixtures of solvents



b) Isotope substitution (H by D) in molecules



 $I(Q) = <\eta^{2} > \int \gamma(r) d^{3}r \exp [i Qr]$ For isotropic, centro-symmetric scatterers $< \exp [i Qr] > = \sin(Qr)/Qr$ $I(Q) = <\eta^{2} > \int \gamma(r) 4\pi r^{2} dr \sin(Qr)/Qr$ Fourier transform $\gamma(r) = (2\pi^{2} < \eta^{2} >)^{-1} \int I(Q) Q^{2} dQ \sin(Qr)/Qr$

	Limits
r = ∞ :	$\gamma(\infty) = 0$
r = 0 :	γ (0) = 1

Pair distance distribution function p (r)



General results from SANS

"Invariant", independent of shape:

$\int \mathbf{I}(\mathbf{Q}) \mathbf{Q}^2 \mathbf{d}\mathbf{Q} = 2\pi^2 < \eta^2 >$
Extrapolation to Q = 0: $I(0) = \langle \eta^2 \rangle \int \gamma(\mathbf{r}) \ 4\pi \ \mathbf{r}^2 \ \mathbf{dr} = \langle \eta^2 \rangle \ \mathbf{V}_p$
Combination \Rightarrow Volume : I(0) / \int I(Q) Q ² dQ = V _p / 2 π^2
$\begin{array}{llllllllllllllllllllllllllllllllllll$
Porod-approximation (large Q) ⇒ Surface S: γ(r) ∞ 1-Sr/4 V+ ⇒ I(Q) ∞ S / Q ⁴

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Real space versus reciprocal space



Particles Scattering : Spheres





 $\begin{array}{ll} & \mbox{Guinier-approximation:} \\ (\mbox{for } QR_g < 1.5) \\ & \mbox{I}(Q) = \Delta \eta^2 \, V_p^{\ 2} \exp((R_g^{\ 2}Q^2/3)) \\ & \mbox{I}_0 & = \Delta \eta^2 \, V_p^{\ 2} \\ & \mbox{R}_g & = (3/5)^{1/2} \, R_0 \end{array}$

Porod approximation: (for QR_g > 2.5) $I(Q) = 2\pi\Delta\eta^2 S^*Q^{-4} = P Q^{-4}$ $S = 4\pi R_0^2$, $V_p = 4\pi/3R_0^3$

Combination

 $R_0^4 = 9 I_0/2 P$

 $I(Q) = \Delta \eta^2 V_p^2 [3(\sin(QR) - (QR)\cos(QR))/(QR)^3]^2$



Cylinders of length L and radius Rc $F^2(Q) = V^2 \Delta \eta^2 \exp(-Q^2 R_c^2/2) / (QL)$ Flat particles with area A and thickness T $F^2(Q) = AT^4 \Delta \eta^2 [\sin(-QT/2) / (-QT/2)]^2 / (QT)^2$ Random coil (Debye) $F^2(Q) = 2 / (QR_g)^4 [(QR_g)^2 - 1 + \exp(-Q^2 R_g^2)]$

Critical fluctuations (Ornstein-Zernicke) $F^{2}(Q) = 1/[1 + (Q^{2}\zeta^{2})]$ with $\zeta^{2} = R_{g}^{2}/3$

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Poly-disperse multi-phase systems



$I(Q) = \iint F_i^2(QR) N_i(R) S_i(QR) dR \otimes Res(Q)$ Form-factor: F(QR) = $\Delta \eta V(R) f(QR)$ contrast \downarrow volume \downarrow shape-factor, size-distribution \downarrow structure factor instrument resolution

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Inter-particle correlations

 $I(Q) = \langle F(Q) | ^{2} \langle 1 + | \langle F(Q) \rangle | ^{2} \langle F(Q) | ^{2} \langle S(Q) - 1 \rangle \rangle$







$S(Q) = 1 + N_p \int [g(r) - 1] exp(iQr) dr$

Pair correlation function g(r) :

 $g(r)=1 + (1/2\pi^2 N_p) \int [S(Q)-1] \sin(QR)/QR Q^2 dQ$

- S(Q) = 1 ideal gas
- S(Q) < 1 repulsive potential:
 - Excluded volume, electrostatic repulsion
- S(Q) > 1 attractive interaction

In practice

 $S(Q, \alpha) = I(Q, \alpha)_{measured} / I(Q, \alpha)_{non-interacting}$

for diluted samples

Magnetic scattering



Magnetic amplitude

p=($\gamma e^2/2mc^2$) g S f_m(Q)=0.27 10⁻¹² cm/ μ_B |M| f_m(Q)

Magnetic scattering lengh density

 $η_{mag}$ = 0.27 10⁻¹² cm/μ_B $\sum M^{\perp}_{l} c_{i} / \Omega_{i}$

Magnetic contrast

 $\Delta \eta_{mag} = \eta_{mag}$ (particle) - η_{mag} (matrix)

Magnetic Scattering from different parts

a) Change of magnitude and /or direction of applied magnetic field



Magnetic particle scattering: *M*⊥

M^{\perp} : Projection of magnetization onto a plane perpendicular to the scattering vector Q



Magnetic particle scattering : $M \perp$

M^{\perp} : Projection of magnetization onto a plane perpendicular to the scattering vector Q



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Magnetic particle scattering : $M \perp$

M^{\perp} : Projection of magnetization onto a plane perpendicular to the scattering vector Q





 $I(Q) = A(Q) + B(Q) \sin^2 \alpha$



Nuclear scattering $A(Q) \propto F_N^2 \propto \Delta \eta_{nuc}^2$ Magnetic scattering $B(Q) \propto F_M^2 \propto \Delta \eta_{mag}^2$

I(Q) $\propto \Delta \eta_{nuc}^2$ + $\Delta \eta_{mag}^2 \sin^2 \alpha$

Transmission-Polariser



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Polarisation

 $P=(n^+-n^-)/(n^++n^-) \qquad n^{+:} \text{ Neutron spin opposite to } H//z$

Magnetic amplitude

 $p=(\gamma e^2/2mc^2) g S f_m(Q)=0.27 10^{-12} cm/\mu_B |M| f_m(Q)$

Atomic scattering amplitudes

 $a^{s,s'} = \langle s' | \mathbf{b}_i - \mathbf{p}_i \mathbf{M}_i \perp \sigma | s \rangle$

$M_{i\perp}$: Component of the magnetic moment perpendicular to Q

σ: neutron-spin operator

 Spin-non flip $a(++) = b - p M \perp^z$ $a(--) = b + p M \perp^z$

 Spin-flip $a(+-) = -p(M \perp^x + i M \perp^y)$ $a(-+) = -p(M \perp^x - i M \perp^y)$

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Amplitudes

Spin-non flip $a(++) = \Delta \eta_N - p \Delta M \perp^z$ $a(--) = \Delta \eta_N + p \Delta M \perp^z$

Spin-flip $a(+-) = -p(\Delta M \perp^x + i \Delta M \perp^y) \quad a(-+) = -p(\Delta M \perp^x - i \Delta M \perp^y)$

Form-factors

 $p=0.27 \ 10^{-12} \ cm/\mu_{B} \mid \Delta M \mid f_{m}(Q), \ F_{N}=\Delta\eta_{N}V_{p}f(Q) \quad , \ F_{M}=p \ V_{p}f(Q)$ $F(++)=F_{N} - F_{M} \ \Delta M \perp^{z} \qquad F(--)=F_{N} + F_{M} \ \Delta M \perp^{z}$ $F(+-)=-F_{M} \ (\Delta M \perp^{x} + i \ \Delta M \perp^{y}) \qquad F(-+)=-F_{M} \ (\Delta M \perp^{x} - i \ \Delta M \perp^{y})$

Amplitudes

Spin-non flip $a(++) = \Delta \eta_N - p \Delta M \perp^z$ $a(--) = \Delta \eta_N + p \Delta M \perp^z$

Spin-flip $a(+-) = -p(\Delta M \perp^x + i \Delta M \perp^y) \quad a(-+) = -p(\Delta M \perp^x - i \Delta M \perp^y)$

Form-factors

$$\begin{split} p=0.27 \ 10^{-12} \ cm/\mu_{B} \mid \Delta M \mid f_{m}(Q), \ F_{N}=\Delta\eta_{N}V_{p}f(Q) \quad , \ F_{M}=p \ V_{p}f(Q) \\ F(++)=F_{N}-F_{M} \ \Delta M \perp^{z} \qquad F(--)=F_{N}+F_{M} \ \Delta M \perp^{z} \\ F(+-)=-F_{M} \ (\Delta M \perp^{x}+i \ \Delta M \perp^{y}) \qquad F(-+)=-F_{M} \ (\Delta M \perp^{x}-i \ \Delta M \perp^{y}) \\ \hline Intensities \ (for \ M \perp^{y}=0) \\ I++(Q) = \ <|F++|^{2} > \qquad I--(Q) = \ <|F--|^{2} > \\ I+-(Q) = \ <|F++|^{2} > \qquad I-+(Q) = \ <|F-+|^{2} > \end{split}$$

M^{\perp} : Projection of magnetization onto a plane perpendicular to the scattering vector Q



P // M // H // z

 $\mathbf{M} \perp^{\mathbf{y}} = \mathbf{0}$

POLARIS: SANS with analysis of polarisation

Spin non-flip scattering (for M⊥^y=0)

$$I++(Q) = (F_N - F_M \sin^2 \alpha)^2$$
 $I--(Q) = (F_N + F_M \sin^2 \alpha)^2$

Spin -flip scattering (for M⊥^y=0)

 $I + -(Q) = (F_M \sin \alpha \cos \alpha)^2$

 $I-+(Q) = (F_{M} \sin \alpha \cos \alpha)^{2}$

SANSPOL and **SANS**

Without polarisation analysis of scattered neutrons

$$I^{+}(Q) = \langle F^{++} | ^{2} \rangle + \langle F^{+-} | ^{2} \rangle = F_{N}^{2} + \{F_{M}^{2} - 2F_{N}F_{M}\}\sin^{2}\alpha$$
$$I^{-}(Q) = \langle F^{--} | ^{2} \rangle + \langle F^{-+} | ^{2} \rangle = F_{N}^{2} + \{F_{M}^{2} + 2F_{N}F_{M}\}\sin^{2}\alpha$$

Anisotropic scattering profile

 $I^{\pm}(Q) = A(Q) + B^{\pm}(Q) \sin^2 \alpha$

Magnetic-nuclear cross term

 $I^{-}(Q) - I^{+}(Q) = 4 F_{N}F_{M} \sin^{2}\alpha$

Sum signal ≡ I (non-polarised SANS)

 $(I^+ + I^-) / 2 = F_N^2 + F_M^2 \sin^2 \alpha$

Nuclear scattering contrast



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Nuclear and magnetic scattering contrast



Polarised neutrons: SANSPOL



Polarised neutrons: SANSPOL


Polarised neutrons (SANSPOL)





$$I(Q) (-) \propto \Delta \eta_N^2 + \{\Delta \eta_M^2 + 2\Delta \eta_M \Delta \eta_N \} \sin^2 \alpha$$

Polarised neutrons (SANSPOL)





$$I(Q) (+) \propto \Delta \eta_N^2 + \{ \Delta \eta_M^2 - 2 \Delta \eta_M \Delta \eta_N \} \sin^2 \alpha$$

SANSPOL Instrument V4



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II. Applications





2. Amorphous magnetic alloys: Soft magnetic materials Bulk Amorphous Hardmagnets

Magnetic Liquids





Magnetic core

Nonmagnetic shell of surfactant Carrier liquid

Applications



Magnetic "drug-targeting"

Diagnostic



Immunoassays BSE- Therapy: Hyperthermie



Transport in strong H- field-gradient Change of relaxation time (SQUID)

Reverse of magnetisation in ac field:(400kHz):local increase of temperature to 50-55°C

Questions



Core: Size, Density, Shape, Distribution?

Shell: Composition, Density, Thickness, Shielding,?



Magnetic Nanostructure: Moment-value , orientation, single-domain, interfaces?

Solvent Penetration?

Aggregation: formation of chains, influence of magnetic field?

Materials

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Core	Со	Fe ₃ C	Fe ₃ O ₄	Ba-
Shell	C ₂₁ H ₃₉ NO ₃ +surfact.	C ₂₁ H ₃₉ NO ₃ +surfact.	Charge Dextrane L-M	Ferrite C ₂₁ H ₃₉ NO ₃
Magnet. Moment	1.7 μ _Β / at	1.6 μ _Β /at	1.24 μ _B / at	0.9 μ _в / at
Carrier liquid	C ₇ H ₈	C ₁₀ H ₂₂	H ₂ O	C ₁₂ H ₂₆
Concentra tion	0.5- 5 vol.%	0.5- 5 vol.%	1- 6 vol. %	2.5 -4 vol.%

2D-SANSPOL scattering patterns

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e.g. Co-Ferrofluid

 $I(Q) = A(Q)+B^{\pm}(Q) \sin^2 \alpha$

 $A(Q) = F_N^2$ $B^+(Q) = F_M^2 - F_N F_M$ $B^-(Q) = F_M^2 + F_N F_M$



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SANSPOL sector perpendicular to H

 $I(\perp)=A(Q)+B^{\pm}(Q)$

Flipping ratio I(-) / I(+)



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 $\Delta \eta$ ^(±) = η ^{nuc} ± η ^{mag} - $\eta_{solvent}$



Contrast variation in Co-Ferrofluids





Density profile in Co-Ferrofluid



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SANSPOL intensities

l(on)	$= \mathbf{F}_{\mathbf{N}}^{2} + \{\mathbf{F}_{\mathbf{M}}^{2} +$	$2\mathbf{F}_{\mathbf{N}}\mathbf{F}_{\mathbf{M}}$
l(off)	$= \mathbf{F}_{\mathbf{N}}^2 + \{\mathbf{F}_{\mathbf{M}}^2 -$	$2\mathbf{F}_{\mathbf{N}}\mathbf{F}_{\mathbf{M}}$
l(on)-l(off)	$= 4\mathbf{F}_{\mathbf{N}} \mathbf{F}_{\mathbf{M}}$	

SANS intensities

 $I(nuc) = F_N^2$ I(mag) = F_M^2

using contraints on parameters e.g.:

∆η ₁ (on)	= ∆η ₁ (nuc)+ η ₁ (mag)
∆η ₁ (off)	= ∆η ₁ (nuc)- η ₁ (mag)

Software		
SASFIT	(J. Kohlbrecher)	
MATHEmatica	(A. Heinemann)	
FISH	(R. Heenan)	



Magnetic single-domain particles

Field variation of magnetization in superparamagnetic particles



$$\sigma/\sigma_{\infty} = L(M_{cr}V_{p}\mu_{0}H_{eff}/kT)$$

Langevin function:

L(x)=coth(x)-1/x



Field variation of SANSPOL cross-term



Superparamagnetic behaviour of non-interacting single domains

AW, Physica B (2001) 226-233 A. Heinemann et al AOC(2004) hmi

Diluted poly-disperse systems





Heinemann et al JAC 2002

Magnetite-FF: Contrast Variation





Magnetic core-shell particles

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Magnetic aggregate

Nonmagnetic micelles

Interparticle correlations: Structure factor S(Q)



$S(Q) = 1 + N_p \int [g(r) - 1] exp(iQr) dr$

- g(r) : Pair correlation function
- S(Q) = 1 ideal gas
- S(Q) < 1 repulsive potential:
 - Excluded volume, electrostatic repulsion
- S(Q) > 1 attractive interaction

In practice

 $S(Q, \alpha) = I(Q, \alpha)_{measured} / I(Q, \alpha)_{non-interacting}$

for diluted samples



Field induced ordering

Non-polarised neutrons, H=1 T





For Coconcentrations above 1 vol. % :





Peaks ! disappear at H=0

A.Wiedenmann et. al. Phys Rev E 2003

3 vol.% Co-FF:SANSPOL



Polarised neutrons



Sectors of 2D-SANSPOL



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Anisotropic in-plane structure factors



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Neutrons parallel to H

Radial averaged SANSPOL



Model fits at high Q
from diluted samples
using form-factors
alone

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Neutrons parallel to H



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Nature of field induced ordering

Hexagonal symmetry



Q₁=0.33 nm⁻¹

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 $Q_2 / Q_1 = \sqrt{3}$

Q₃= 0.24 nm⁻¹

(0 0 1), (0 0-1)

a_{hex} = 21.3 nm

c_{hex}= 26.1 nm

Q (hkl) = $2\pi / \{4(h^2+k^2+hk)/3a^2 + l^2/c^2\}^{1/2}$

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Texture type I





Turn around [110] direction of magnetic field by 90°

Particle moments M // [110] along H

Pseudo-crystalline hexagonal particle alignment

Texture type I







R_c = 3.8 nm, d=1.9 nm a_{hex}= 21 nm c_{hex}= 25...70 nm

A.Wiedenmann, A. Hoell, M. Kammel, Phys Rev E 2003

Field induced ordering: 1 vol.% Co-FF



Diffraction planes perpendicular to H



1 vol.% Co-FF: Field induced ordering

Diffraction planes perpendicular to H





Ordering- 1st step:chaining



SANSPOL Co-3 vol.%



Co-existence of uncorrelated chains with hexagonal ordered domains

Intensity I(Q \perp H) at low Q:



Ordering - 1st step:chaining

Particle size <R_c> = 3.8 nm, d=1.9 nm

Shortest possible distance 11.4 nm

OBSERVED:

Diffraction planes from quasi-1D chains

σ (1D) =2π /∆**Q= 21 nm**



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Second example: Co-L9



SANS H=1 T



Confirmation of hexagonal symmetry

Field dependence of ordering

6 vol.% Co in Oil L9

SANS H=1 T Saturation



 $I(Q,\alpha) = [F_{M}^{2}L^{2}(x) \sin^{2}\alpha + F_{N}^{2}] S(Q,\alpha) + [2L(x)/x - \sin^{2}\alpha(L^{2}(x) - 1 + 3L(x)/x)] F_{M}^{2}$

SANSPOL : Co-L9





2D Gaussian fits

At H=1 T:

Hexagonal ordering a_{hex} = 18.9 nm

c_{hex} = 21.9 nm

= 4 $F_N F_M L(x) S(Q,\alpha)$







At H=0



Isotropic scattering: no peaks

Cylinders L= 80 nm R=4.3 nm (5-6 core-shell particles)

Spontaneous formation of chains I(Q) ~Q⁻¹





At H=1 T



Anisotropic scattering peaks

Perpendicular to H $Q_3 = 0.28 \text{ nm}^{-1}$

> Sectors $\pm 30^{\circ}$ Q₁= 0.38 nm⁻¹

Co-existence with chain segments

I(Q) ∼Q⁻¹




SANS sector 90°: Difference I_H- I_{H=0}



Structure evolution in magnetic field



SANSPOL sector 30° I(on)-I(off)



Competition between magnetic dipol-interaction and thermal energy

$$\gamma = M_{sat}^2 V_c^2 \mu_o / 4\pi k_B^2 T \sigma^3$$
.

Present range 1.4 < γ < 8

Computer simulations (MSA): Chaining in head-to tail conformation (de Gennes)

Lateral attractions

Molecular dynamics predicts close-packed structures (Hess)





SANSPOL with Polarisation Analysis

A. Wiedenmann Physica B 2005

³He filter

Spin selective neutron absorption: (Passell, Schermer; 1966)

³He[↑] + n $\downarrow \rightarrow$ [⁴He^{*}] \rightarrow t + p + 0.76 MeV $\sigma^{\uparrow\downarrow} \approx 3000 \text{ barn} \cdot \lambda$ [Å]

 3 He \uparrow + n \uparrow \rightarrow 3 He \uparrow + n \uparrow $\sigma\uparrow\uparrow\approx$ 5 barn

Advantages:

- Unlimited angular acceptance
- Broad range of usable wavelength
- No deflection of neutrons
- Polarisation of white beams

Compression of polarised ³He

Piston compressor at University of Mainz



Neutron polarisation P_n P_n =tanh(x P_{He}) Transmission $T(+) = exp[-x (1 - P_{He})]$ $T(-) = exp[-x (1 + P_{He})]$ $T_0 = exp[-x]$

Neutron optique filter thickness : $x = N_{He} \sigma$

Performance of ³He filter



Outlook:Polarisation analysis at V4





Outlook:Polarisation analysis at V4





Outlook: POLARISation analysis at V4





Neutron Scattering cross-sections *



 $d\sigma/d\Omega(ij)(Q,\alpha,x) = C_{ij} S(Q,\alpha) + D_{ij} F_{M}^{2}(Q)$



* R.M.Moon, T.Riste, W.C. Koehler (1969) and R.Pynn, J.Hayter (1983)

$d\sigma/d\Omega$ (ij)(Q, α) = C_{ij} S(Q, α) + D_{ij} F_M²(Q)

a(x) = L(x) / x, $b(x) = L^2(x) - 1 + 3 a(x)$ $x = M_{cr}V_p \mu_0 H_{eff} / kT$

ij	C _{ij}	D _{ij}
POLARIS		
l(++)	$[F_M L(x) \sin^2 \alpha - F_N]^2$	a sin ² α – b sin ⁴ α
l()	$[F_M L(x) \sin^2 \alpha + F_N]^2$	a sin ² α – b sin ⁴ α
l(+ -),l(-+)	[F _M L(x) sin $\alpha * \cos \alpha$] ²	a(2- sin ² α) - b [sin α * cos α] ²
SANSPOL		
l(+)	$[F_{M}^{2}L^{2}(x) - 2F_{M}F_{N}L(x)]sin^{2}\alpha + F_{N}^{2}$	<mark>a - b sin</mark> ²α
l(-)	$[F_{M}^{2}L^{2}(x) + 2F_{M}F_{N}L(x)] \sin^{2}\alpha + F_{N}^{2}$	<mark>a - b sin</mark> ²α
l(-) - l(+)	4 F _M F _N L(x) sin²α	0
SANS		
I(Q)	$F_{M}^{2} L^{2}(x) \sin^{2}\alpha + F_{N}^{2}$	2a - b sin²α

$d\sigma/d\Omega$ (ij)(Q, α) = C_{ij} S(Q, α) + D_{ij} F_M²(Q)

a(x) = L(x) / x, $b(x) = L^2(x) - 1 + 3 a(x)$ $x = M_{cr}V_p \mu_0 H_{eff} / kT$

ij	C _{ij}	D _{ij}
POLARIS		
l(++)	$[F_M L(x) \sin^2 \alpha - F_N]^2$	a sin ² α – b sin ⁴ α
l()	$[F_M L(x) \sin^2 \alpha + F_N]^2$	a sin ² α – b sin ⁴ α
l(+ -),l(-+)	Purely magnetic	Purely magnetic
SANSPOL		
l(+)	$[F_M^2L^2(x) - 2F_MF_NL(x)]sin^2\alpha + F_N^2$	a - b sin²α
l(-)	$[F_{M}^{2}L^{2}(x) + 2F_{M}F_{N}L(x)] \sin^{2}\alpha + F_{N}^{2}$	a - b sin²α
l(-) - l(+)	4 F _M F _N L(x) sin²α	0
SANS		
l(Q)	$F_{M}^{2} L^{2}(x) \sin^{2}\alpha + F_{N}^{2}$	2a - b sin²α



Field induced ordering?- Non-perfect alignment of M?

5 vol % Co-FF in C_6D_8 H=0.05 T





Field induced ordering?- Non-perfect alignment of M? 5 vol % Co-FF in C₆D₈ H=0.05 T



POLARIS I(-): spin-flip

$I_{sf}(Q) = F_{M}^{2} \{ [L^{2}(x) S(Q) - L^{2}(x) + 1 - 3L(x)/x] [sin^{2}\alpha * cos^{2}\alpha] + (2 - sin^{2}\alpha)L(x)/x \}$



Uniaxial polarisation analysis in SANS: Additional contrast variation technique

Imperfect alignment of magnetic particle moments

Field induced correlations S(Q,α): Separation of purely magnetic contribution and purely nuclear contributions

Summary: Nanostructures in Ferrofluids





Competing interactions

Magnetic dipol-interaction versus thermal energy





1. Magnetic colloids: "Ferrofluids"

2. Amorphous magnetic alloys:

Soft magnetic materials

Bulk amorphous hardmagnets

Soft magnetic alloys

Fe_{73.5} Si _{15.5}B₇Nb₃ Cu₁: Amorphous ferromagnetic alloy

Improvement of soft magnetic properties by heat treatment



Magnetic contrast variation



Ferromagnetic single domain in ferromagnetic matrix Magnetisation ratio $\gamma(T) = M^{am} / M^{cr}$ Effective magnetic field $M_{eff}(T)$

Scattering cross-section [cm⁻¹]



SANSPOL :Fe-Si-B-Nb-Cu 2 % n-Fe₃Si



A.W. Physica B (2001) 226-233, Heinemann et al 1999

Diffusion zone: Enrichment of Nb

Contrasts for magnetic core, shell and matrix

$$\begin{split} \Delta n_1^{(\pm)} &= n_1^{nuc} \pm n_1^{mag} - n_m^{(\pm)} \\ \Delta \eta_2^{(\pm)} &= \eta_2^{nuc} \pm \eta_2^{mag} - \eta_m^{(\pm)} \\ \eta_m^{(\pm)} &= \eta_m^{nuc} \pm \eta_m^{mag} \end{split}$$



Weak magnetization in the interface between n- Fe₃Si matrix

11111

M^{\perp} : Projection of magnetization onto a plane perpendicular to the scattering vector Q



∠β (ΔΜ , z)

 $|M \perp | = |M| \sin \alpha$

$$\begin{array}{ll} \mathsf{M}\bot^z &= \mid \mathsf{M}\bot\mid \cos\left(90^\circ - (\alpha+\beta)\right) \\ &= \mid \mathsf{M}\bot\mid \sin(\alpha+\beta) \\ &= \mid M \mid \sin\alpha\,\sin(\alpha+\beta) \end{array}$$

 $\begin{array}{ll} \mathsf{M} \bot^{\mathsf{x}} &= \mid \mathsf{M} \bot \mid \sin(90^{\circ} \text{-}(\alpha + \beta)) \\ &\mid M \mid \sin \alpha \cos (\alpha + \beta) \end{array}$

 $\mathbf{M} \perp^{\mathbf{y}} = \mathbf{0}$

Results: Soft magnetic metallic glasses

Diffusion controlled enrichment of Nb around nanocrystals: magnetic dilution:



Weak magnetic interlayer reduces ferromagnetic coupling between matrix and nanocrystals

$Nd_{60} Fe_x Co_{30-x} Al_{10}$ alloys



Magnetic Properties

Non-magnetic

Soft-magnetic

hard-magnetic

versus

compositions and cooling conditions Mold-casting Melt-spinning

Magnetic and crystalline microstructure

Magnetization at H=1T





H= 1T

Almost isotropic patterns

Near Tc₂: Decrease of I(Q) at low Q

Weak magnetic contribution



SANSPOL in Nd₆₀Fe₂₀Co₁₀Al₁₀









Volume distributions

Scattering contrasts





Scattering length densities in Nd₆₀Fe₂₀Co₁₀Al₁₀ Imi



Microstructure model Nd₆₀Fe₂₀Co₁₀Al₁₀





Fe-rich crystalline phase

Nd-rich crystalline phases f₁ and f₂

Nd-rich amorphous phase f₃

Nanostructure $Nd_{60}Fe_{20}Co_{10}AI_{10}$ T > T_{c2} $\frac{hmi}{2}$



Nd-rich Amorphous f₃ Nd-rich crystalline f₁ and f₂

Fe-rich crystalline

$Nd_{60}Fe_{20}Co_{10}AI_{10}$ $Tc_1 < T < Tc_2$



Nd-rich
amorphousNd-rich
crystallineFe-rich
crystallineparamagneticparamagneticMagnetic domainsHard magnetic
$Nd_{60}Fe_{20}Co_{10}AI_{10}$ T < Tc₁



Nd-rich amorphous ferromagnetic

Soft-magnetic

Nd-rich crystalline ferromagnetic

Fe-rich crystalline ferromagnetic

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- Nanosized Nd-rich crystalline and amorphous particles embedded in Fe-rich ferromagnetic crystals.
- They are paramagnetic between T_{c1} and T_{c2} .
- They act as pinning center for magnetic domains (hard-magnetic behaviour).
- Below T_{c1} they are ferromagnetic with higher magnetization than Fe-rich crystals.



Magnetic Nanostructures Investigated by Small Angle Neutron Scattering

Albrecht Wiedenmann

VIII School of Neutron Scattering "Francesco Paolo Ricci" Structure and Dynamics of magnetic systems

http://www.hmi.de/bereiche/SF/SF3/methods/sans/index_en.html

25.09-6.10.2006 Santa Margerita di Pula (Italy)

III. Practice of SANS

hmi

1. Instrument V4

2. SANS-Experiment, Data Reduction and analysis

3. Tutorial and problem class

SANSPOL Instrument (V4 at hmi)



SANSPOL Instrument V4



SANSPOL components

Transmission-Polariser $\lambda > 0.48$ nm

V-shaped l = 1.8m, $\alpha = 8.33$ mrad

CoFe/Si Supermirror m = 2

FeNdB magnets: 7*100mm H=1kGs

Magnetic guide field

Permanent magnets NdFeB on steel rods Steel plates ST37, H=10 Gs in centre

Spin-flipper

 $B_0=100Gs$, dB/dl=3Gs/cmRF-coil 1=5cm, 3 w/cm: $\omega_0 = \gamma B_0 = 300 kHz$ =Larmor frequency $B_1 = 20 Gs$

Collimator





Transmission - Polariser

Critical angle for total reflection: $\Theta_{c (up)}$ [mrad]= m * 1.73 λ $\Theta_{c (dn)}$ [mrad]= 0.7* λ



No divergence of incident neutrons

 $\Theta_{\rm eff} = \alpha$



Divergence of incident neutrons



Critical angle for total reflection:

 $\Theta_{c (up)}$ [mrad]= m * 1.73 λ $\Theta_{c (dn)}$ [mrad]= 0.7* λ Prototype Fe-Si SM m=3.8

$$\Theta_{\rm eff} = \alpha \pm \delta i v$$

For collimators: $\delta iv = \phi / L$



For neutron guides m=1

 $\delta iv = 1.73 \lambda$

Decrease of Polarisation

Divergence from collimators



Divergence from a neutron guide



hm

V4 polariser behind a neutron guide



V4-m=2 polariser behind a Ni-guide



Magnetic guide-field





Spin-flipper





Performance of SANSPOL



T. Keller, T. Krist, A. Danzig, U. Keiderling, F. Mezei, A. Wiedenmann J. Nuclear Instruments A451(2000), 474-479

III. Practice of SANS

1. Instrument V4

2. SANS-Experiment, Data Reduction and analysis

3. Practice and sensitivity of SANSPOL

Choice of Q-range





 $\begin{array}{ll} \mbox{Guinier-approximation:} \\ \mbox{(for QR}_g < 1.5) \\ \mbox{I(Q)} = \Delta \eta^2 \, V_p^{\ 2} \exp((R_g^{\ 2}Q^2/3)) \\ \mbox{I}_0 &= \Delta \eta^2 \, V_p^{\ 2} \\ \mbox{R}_g &= (3/5)^{1/2} \, R_0 \end{array}$

Porod approximation: (for QR_g > 2.5) $I(Q)=2\pi\Delta\eta^2 S^*Q^{-4} = PQ^{-4}$ $S = 4\pi R_0^2$, $V_p = 4\pi/3R_0^3$ $R_0^4 = 9 I_0/2 P$

 $I(Q) = 9 \Delta \eta^2 V_p^2 [(\sin(QR) - (QR)\cos(QR))/(QR)^3]^2$

Choice of Q-range





hm

Intensity versus resolution



Optimized Conditions: L = I d_E=2d_s

Practice: Intensity contributions



Differential scattering cross section of scuple $\frac{ds}{ds}_{s} = \frac{I_{0} - I_{0}}{C(\lambda) \cdot t_{c} \cdot T_{e}} \cdot I_{BC}$

1.	Scattering with sample + holder	•	ID
2.	Scattering of holder abue		I _{R6}
3.	Residual noise		Io
4.	Transmission of sample - holder	١	Ts
5.	Transmission of holder alone	:	TRE
6.	Absolute calibration	•	Ca

Must be measured for each spectrometer configuration L, L, de, dsp, 7,...

hmi

Calculated

$$T = \exp(\sigma_{tot} t_s / \rho_s)$$

- ts : Thickness of sample
- ps : Atomic density
- otot: Total cross section

 $\sigma_{tot} = \sigma_{abs}(\lambda) + \sigma_{coh} + \sigma_{inc}$ $\sigma_{abs}(\lambda)$: absorption cross section (tabulated for $\lambda=1.8A$) $\sigma_{abs}(\lambda_1) / \sigma_{abs}(\lambda_2) = \lambda_1 / \lambda_2$

Measured : Direct beau : $T = \mathbf{I}_{0}(0) / \mathbf{I}(0) = \frac{\sum_{i=1}^{3 \text{ choch}} (Gouts) \text{ with sam}}{\sum_{i=1}^{3 \text{ choch}} \sum_{i=1}^{3 \text{ choch}} (Gouts) \text{ with sam}}$

- Attenuation of primary beam $f_a = 10.200$
- Saturation corrrections
 - I (true) = I (measured) * $[1-t_1 * I(measured)]$
- · Good collimation (small profile- low contributions of scattering contained in primary beam)



Measurement of empty can with attenuator $f_a=270$



Calibration of absolute intensity

$$C(r) = \phi_0(r) E(r) \cdot \Delta P \cdot ds^2$$

$$d_s^2/L^2$$
Kut be known cell by cell.

[I] Using Aandards with known (ds/dr)

(u) • eg H20 & dinc. = 160 bom/at

• Vanadium dinc = 5 barn/at

Incolumnt scalling inclopendent of Q:

 $T(a) = C(r) \cdot (1 - T_s)/4\pi \cdot scalled$

in 4TT?

• H20: g(1) • ~0.8-1.3 inclose could be

• Vanadium : couly for large & clus

- to SANS from imparties

- the incolumnt B6.

- · Compressibility of a liquid
- · Dorous syntees well grain site > tim
 - $\left(\frac{ds}{dx}\right) = 2\pi \left(0b\right)^2 \frac{s}{\sqrt{2}} \frac{1}{Q_y}$ Porod
- · Glany carbon: Well-know pre calibrated Q-dependencity
- · Ceramics .
- 2 <u>Direct beau measure meub</u> $I_3 = f_a(\lambda) \cdot \phi_0 \ \mathcal{E}(\lambda) \cdot \Delta \mathcal{R}$ with alternation factor $f_a(\lambda)$ lenown: $f_a(\lambda) = \sum (Caunts) \text{ of strong sections Absorbed}$ Un alternation. $\phi_0 \ \mathcal{E}(\lambda) \cdot \Delta \mathcal{R} = I_0 / f_a(\lambda)$ No calibration of individual cells: $E_i(\lambda)!$

ti mi

Data reduction : cell by cell corrections



Anisotropic raw data:

"SC 16", $\lambda = 1.2$ nm, SD = 16 m / 4 m (water)





sample 150 min, 26 n/s

sample holder 180 min, 16 n/s



water cell 120 min, 93 n/s





empty cell 120 min, 64 n/s

cadmium 180 min, 2.4 n/s

Anisotropic reduction procedure:

corrected intensity



error of corrected intensity

$$\Delta I = \sqrt{\left(\frac{\partial I}{\partial Cd} * \Delta Cd\right)^2 + \left(\frac{\partial I}{\partial W} * \Delta W\right)^2 + \left(\frac{\partial I}{\partial WB} * \Delta WB\right)^2 + \left(\frac{\partial I}{\partial SB} * \Delta SB\right)^2 + \left(\frac{\partial I}{\partial S} * \Delta SB\right)^2}$$

- masking
- result: corrected anisotropic data file







corrected data



mask

Azimutal averages



Cuts starting from beam-centre





270[°]dag

Data reduction software: BERSANS





Data analysis: Model fitting (SASFIT, FISH..)



Non-linear least-squares fit

$$\chi^2 = \sum \{ (I(\mathbf{Q}_i)_{exp} - I(\mathbf{Q}_i)_{calc}) / w_i \}^2$$

 $\eta_2 \dots$

$$\Rightarrow$$
 <**R**₁>, $\sigma_{1,}$,**f**_{1,} η_1 <**R**₂>, $\sigma_{2,}$,**f**₂

Data analysis: (Inverse Fourier transform)





p (r)= 4π r² γ (r) e.g: spheres of Radius R





III. Practice of SANS

Instrument V4

SANS-Experiment, Data Reduction and analysis

Examples

Suspension of latex (C₈H₈) in D₂O

	Latex	Heavy water
Molecular weight	104	20
Number of electrons	56	10
Mass-density/ gcm ⁻³	1.0	1.1
Scattering length den	sity	
(X-ray) 0.281e-12	0.281e-12*1.1*10N _I /20	

Neutrons $N_L/104$ (8b_c+8b_H) 1.1*N_L / 20(2b_D+b_O)

ContrastX-ray $(\eta_1 - \eta_2) = -0.195 e^{10} cm^{-2}$ neutrons $(\eta_1 - \eta_2) = -4.970 e^{10} cm^{-2}$



Example 3:Sintering of nano-ceramics





Example3:Sintering of nano-ceramics



Moments of distribution

1. scatt. contrib.:	calc: yes
LogNorm	Sphere
<r^1> = 4.3403</r^1>	<r^5> = 11935.4</r^5>
<r^2> = 23.1189</r^2>	<r^6> = 144147</r^6>
<r^3> = 151.127</r^3>	<r^7> = 2.13335e+006</r^7>
<r^4> = 1212.39</r^4>	<r^8> = 3.85619e+007</r^8>
R_lc = 8.0223	lc = 5.3482
R_li = 6.53696	li = 4.90272
R_Ac = 8.88684	Ac = 31.4235221923
R_VP = 9.8436	VP = 227.705365117
R_RG = 16.356	RG = 12.6693031221

 $R^2 = \langle R^8 \rangle / \langle R^6 \rangle = 16.3$ RG=(5/3)^{0.5}R=12.7

Number density: $n_p (1-f) = N \ 10^{42} = 1.14 \ 10^{21} T \ cm^{-3}$

Volume fraction: $f(1-f)=n_p 4/3\pi < R^3 > =1.15 4/3\pi 151.12 10^{21}10^{-21}= 0.07$

Example3 : Sintering of nano-ceramics





Guinier radius: 11.65 nm

Invariant =f(1-f)/ $2\pi^2 \Delta \eta^2$ $\Delta \eta$ =5.33 10¹⁰ cm⁻² f(1-f)=4.04 10²¹/ $2\pi^2 \Delta \eta^2$ =0.07
Example 4: 2 D-scattering pattern



background



empty beam

sample

water







Wask file; no file Data file: D0035964.00

> Correction for background, efficiency. Beam center Cuts and radial averages



Н

2 D-fit: I=A+B[±] sin² α or segments:

hmi

$I(Q//H) = I_{nuc}(Q)$ $I(Q\perp Q) = A+B^{\pm}$





Scattering length density profile in Co-Ferroflui



Scattering length densities [10¹⁰cm⁻²]

Scattering length densities

1.	Solvent: 43% D / Η mixture of Toluene : η (x)= -0.7(1-x) + 6.8 x	ղ (solv)	= 3.0 10 ¹⁰ cm ⁻²
2.	Non-deuterated shell (C ₂₄ -H ₂₀ -N-O ₂)	n(shell)	$= 0.3 \ 10^{10} \ \mathrm{cm}^{-2}$
3.	Nuclear sld (Co)		
	Ω(Co)=0. 0.01099 nm³/at ,		
	b= 10 ⁻¹² cm	ղ (Co)	= 2.56 10 ¹⁰ cm ⁻²
4.	Magnetic sld Co		
	$m_0 = 1.715 \ \mu_B / atom$		
	η(mag) = 0.27 10 ⁻¹² m ₀ /Ω	ղ (mag)	= 4.3 10 ¹⁰ cm ⁻²

Contrasts for SANSPOL

core - solvent :Δη₁ (on) = η(Co) +η(mag) - η(solv) = 2.56+4.3-3.0= 3.83 10¹⁰ Δη₁ (off) = η(Co) - η(mag) - η(solv) = 2.56-4.3-3.0= -4.70 10¹⁰ Shell-solvent Δη₂ = η(shell) - η(solv) = 0.3-3.0 = -2.73 10¹⁰





Fit of nuclear scattering (I(Q) // H)







I(Q) – on (I perp. H)





I(on) same parameters –except η_1 (on)





I(off) same parameters –except η_1 (off)





Optimisation by simultaneous fits l(on) l(off) and l(nuc) or l(mag) using contraints on parameters e.g.:

> $\Delta \eta_1 \text{ (on)= } \Delta \eta_1 \text{ (nuc)+ } \eta_1 \text{(mag)}$ $\Delta \eta_1 \text{ (off)= } \Delta \eta_1 \text{ (nuc)- } \eta_1 \text{(mag)}$

Software: SASFIT (J. Kohlbrecher , PSI) FISH (R. Heenan ISIS) MATHematica (A.Heinemann, hmi) Max Entropie: (Tatschev, hmi)



Number density: $n_p (1-f) = N \ 10^{42} = 6.59 \ 10^{15} T \ cm^{-3}$

Volume fraction: $f(1-f)=n_p 4/3\pi < R^3 > =6.59 4/3\pi 55.5 10^{15}10^{-21} = 0.0015$

Moments of distribution



Problem class



1.) Calculate the nuclear and magnetic scattering contrast of ferromagnetic Co-particles with fully aligned magnetic moments in a solvent of a mixture of $40\% D_2O$ and $60\% H_2O$.

- H_2O : density 1 g/cm³
- D_2O : density 1.1 g/cm³
- H: Scattering lenght b= -0.374 10⁻¹²cm
- D: Scattering lenght b= 0.667 10⁻¹²cm
- O: Scattering lenght b= 0.581 10-12cm
- Co: Atomic volume Ω (Co)=0. 0.01099 nm³/at , Scattering lenght b= 0.278 10⁻¹²cm, Magnetic Moment m0 = 1.715 μ B/ atom

V. New developments-Dynamical SANS





Spacial fluctuations of density, composition, magnetization Dynamical fluctuations-What time range Relaxation of field-induced order: Time-resolved SANSPOL

Ordering and re-ordering in oscillating external field:

Continuous stroboscopic SANS Pulsed stroboscopic "TISANE"

Time-resolved SANSPOL





Decay of nuclear and magnetic correlations measured in time slices of 0.1 s

Duty cycle 15-30 s



MFT3N1: 6% Co in oil L9 H=0.5T



SANS cross-section



SANSPOL Difference cross-section



Magnetic and Nuclear correlations + Misalignment of magnetic moment + Nomagnetic contributions

Magnetic-Nuclear correlations

SANSPOL Differences



Chain segments remain partly aligned along remanent field

11111

Time dependence of I(+)- I(-)



Sectors α = 90° (Q \perp H)



Time dependence of I(+)- I(-)

hmi

Sectors α = 30°



A. Wiedenmann, U. Keiderling, R. P. May, C. Dewhurst Physica B 2006



•Slow decay of field-induced ordering (few seconds).

•Fully reversible relaxation onto equilibrium

•Single exponential decay: Time constants depending on Q and B_{max} τ (intra-chain) > τ (in-plane) > τ (inter-plane)

Switch-on of B:

Reordering follows B-sweep rate: Process too fast!

Response from oscillating magnetic field



B(t)=B₀*sin ($2\pi\nu$ t + ϕ)





Frequency v = 50 - 2200 Hz

Duty cycle: 3-5 orders of magnitude shorter

Trigger for list-mode data acquisition in 2 D detector

Continuous stroboscopic SANS





2D- Detector: Time stamped recording of each scattered neutron

Sample S Oscillating magnetic field: oscillating magnetic contrast

Continuous monochromatic flux

 $t_{TOF}[ms] = \lambda[nm] * L_2[m] * 2.52778$

Continuous stroboscopic SANS at V4



 $\Delta t_R / T_s = \Delta \lambda * t_{TOF} * v_s$



Pulsed time-involved SANS (TISANE)



Pulsed time-involved SANS (TISANE)



All neutrons scattered at the sample in the same oscillation state are recorded in the same time channel hm

Pulsed time-involved SANS (TISANE)



TISANE@NEAT



 L_1 =13 m, L_2 =4 m, v_E (max)=666 Hz , v_s (max)=2800 Hz



TISANE@NEAT



 L_1 =13 m, L_2 =4 m, v_E (max)=666 Hz , v_s (max)=2800 Hz





11 111 7

Results: Static SANS





Results: Static SANS







Time*frequency









with increasing frequency

MFT3N1-13

0,0 0,2 0,4 0,6 0,8 1,0

Frequency * time

hm

Continuous SANS




Pulsed TISANE:





α = **90°**



Superparamagnetic behaviour



 σ/σ_{∞} = L(M_{cr}V_p μ_0 H_{eff} / kT)

Langevin function:

L(x)=coth(x)-1/x



Known from SANS: V_p Particle volume M_{cr} Magnetic particle moment hm

Magnetic single-domain particles

Static SANS cross-section



$$[F_{M}^{2} L^{2}(x) \sin^{2}\alpha + F_{N}^{2}] S(Q,\alpha)$$

R. Pynn et al (1983), J. Kohlbrecher , AW 1997

Magnetic single-domain particles

Static SANS cross-section



Magnetic and Nuclear correlations

+ Misalignment of magnetic moments

SANS cross-section in oscillating field



Stroboscopic versus static SANS





SANS cross-section in oscillating field

$B(t)=B_0 \sin(2\pi v t) + B_{st}$



Fraction of freely rotating moments



Fraction of freely rotating moments





Neel-relaxation of single particle moment



Néel-relaxation of particle moment



hm



Néel-relaxation of particle moment







Attractive interactions





Formation of ordered domains



Relaxation: Rotation of ordered domains





Relaxation: Rotation of ordered domains





Summary



Mechanisms and dynamics of field-induced ordering in Co-Ferrfluids determined by Brownian relaxation

Dynamical processes in nanoscaled inhomogeneities are observable by time-resolved SANS

Limitation of continuous techniques: $\Delta\lambda/\lambda$ Pulsed TISANE technique: Sub-millisecond range

Complementary to Photon correlation spectroscopy (PCS,XPCS), Forced Raleigh scattering, ac-χ

Closes the gap between inelastic neutron scattering/ Mössbauer (10⁻¹²- 10⁻⁶ s) and static measuremets

A. Wiedenmann, U. Keiderling, K. Habicht, M. Russina, R. Gähler PRL 97 057202 (2006)





Combining *"SANS, SANSPOL, POLARIS "* Contrast variation technique for magnetic materials

Weak magnetic versus strong nuclear contributions and vice-versa:

Density profiles, interfaces Sign and magnitude of contrast

Separation of magnetic and nuclear contributions

Dynamics in nanoscaled materials in sub-ms range



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